

Search for Fermiophobic and Topcolor Higgs in the Diphoton Final States at DØ

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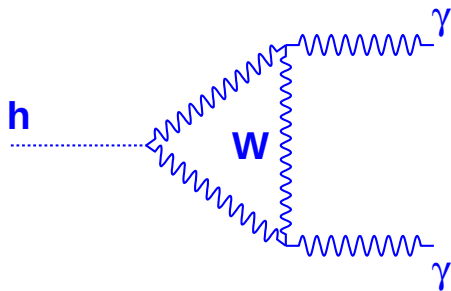


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Riverside, CA**



$h \rightarrow \gamma\gamma$ Decay Mode in the SM and Beyond. Why $\gamma\gamma$?

- no tree-level $H\gamma\gamma$ coupling (Higgs boson is neutral)
- Di-photon decays happen via W loop

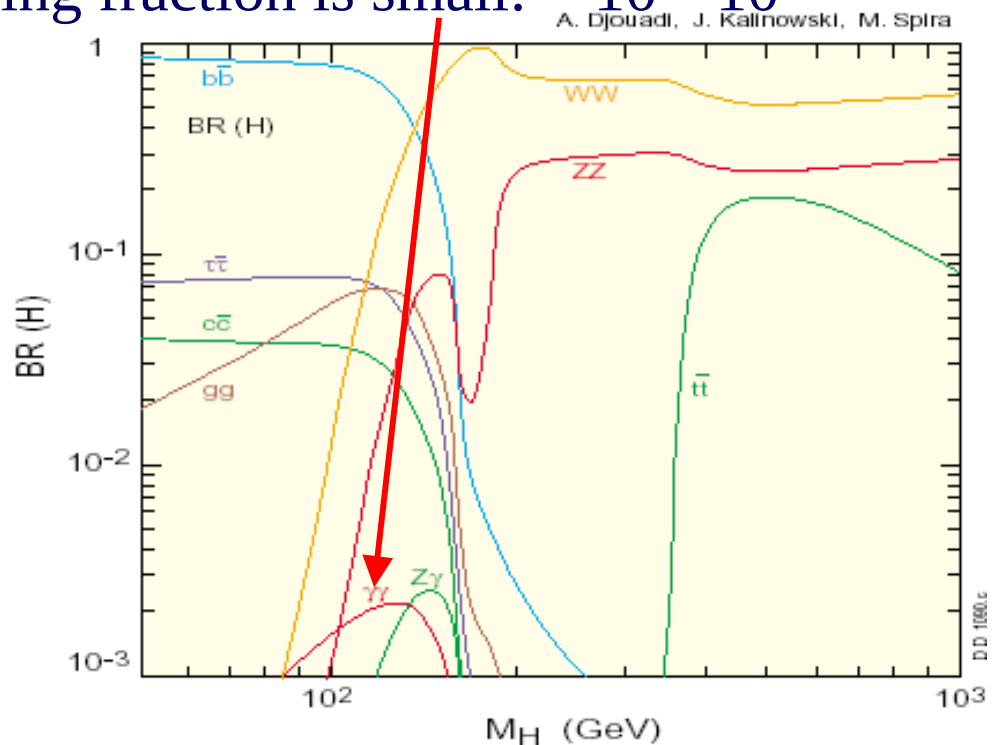


(top-quark loop reduces $H \rightarrow \gamma\gamma$ width by 20-30% due to destructive interference)



- $H \rightarrow \gamma\gamma$ branching fraction is small: $\sim 10^{-3}$ - 10^{-4}

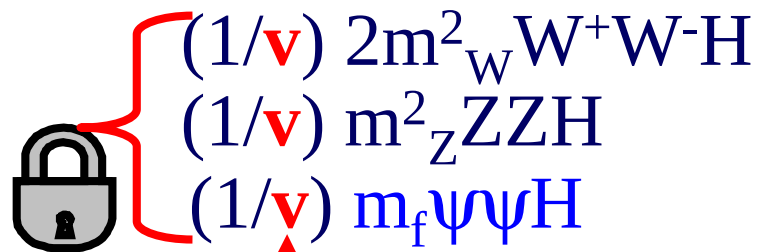
Higgs decays mostly to b-quark and W pairs



- However many extensions of the SM predict enhanced $\gamma\gamma$ decay rate of the Higgs boson

How Can $h \rightarrow \gamma\gamma$ Decays be Enhanced

- One possibility: suppress fermion decay rates (then the full decay width would only be shared with $h \rightarrow WW$ and $H \rightarrow ZZ$ decays)
- In the Standard Model Higgs is responsible for both **Electroweak symmetry breaking (EWSB)** and **fermion masses**
- Higgs-**W,Z/fermion** interaction terms in the SM Lagrangian are:

$$\begin{aligned} & \left\{ \begin{array}{l} (1/v) 2m_W^2 W^+ W^- H \\ (1/v) m_Z^2 Z Z H \\ (1/v) m_f \bar{\psi} \psi H \end{array} \right. \end{aligned}$$


- Relative strength of **Higgs-W,Z** and **Higgs-fermion** couplings is fixed by **v** (Higgs field vacuum expectation value)
- In a more general situation different sectors of theory (or different mechanisms) could be responsible for **EWSB** and (some of) **fermion masses** $\Rightarrow$ Higgs couplings to fermions and those to weak gauge bosons can vary independently

Some Examples of $h \rightarrow \gamma\gamma$ Enhancement

- Strong dynamics [1] :
EWSB = via technicolor condensation,
fermion masses = via “extended technicolor”
- Generic two Higgs doublet model [2]:
v1 – W,Z masses, v2 – fermion masses
- Top-condensation model + Higgs [3]
top and bottom quarks get their masses from the topcolor interaction while **Higgs gives mass to remaining fermions and gauge bosons**
- Topcolor Higgs=(top,vector-like quark)[4]
Top quark mass from Higgs, other fermion and **W/Z masses -- from other interactions present in the theory**
- Still another possibility of $h \rightarrow \gamma\gamma$ enhancement arises when the Higgs field extends into **large extra dimensions** (while fermion and gauge boson fields remain confined to 3D) [5]

[1] E.Farhi (CERN), L. Susskind Phys.,Rept.**74**, 277 (1981)

[2] H.E.Haber et al, Nucl. Phys. **B161**, 493 (1979)

[3] J.D. Wells, Phys. Rev. **D56**, 1504 (1997)

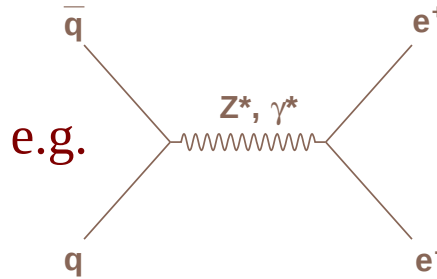
[4] B.Dobrescu, Phys. Rev. **D63**, 015004 (2001)

[5] L. Hall, C.Kolda, Phys. Lett. **B459**, 213 (1999)

Major Backgrounds :

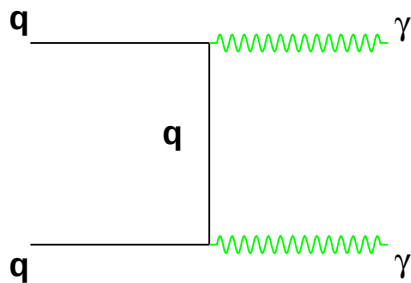
Drell-Yan and QCD

- $Z/\gamma^* \rightarrow e^+e^-$ with e^+e^- misidentified as photons (lost tracks)

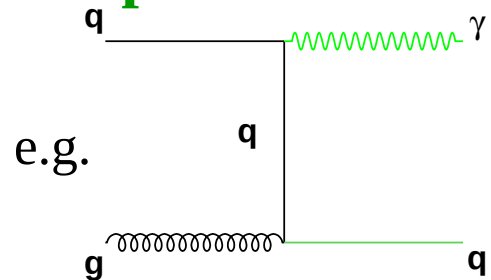


- QCD processes that in the final state contain :**

1. two photons

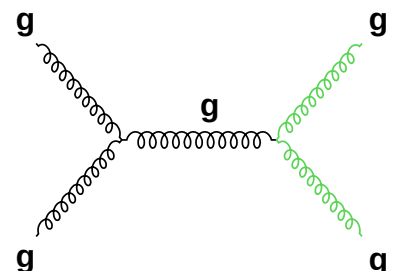
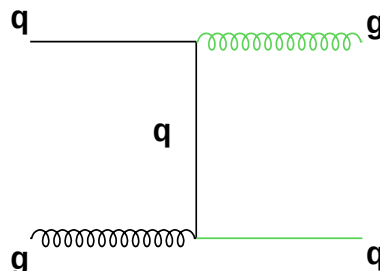
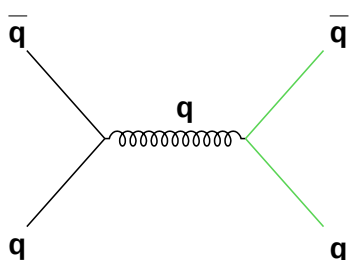


2. a photon and a hadronic jet misidentified as photon



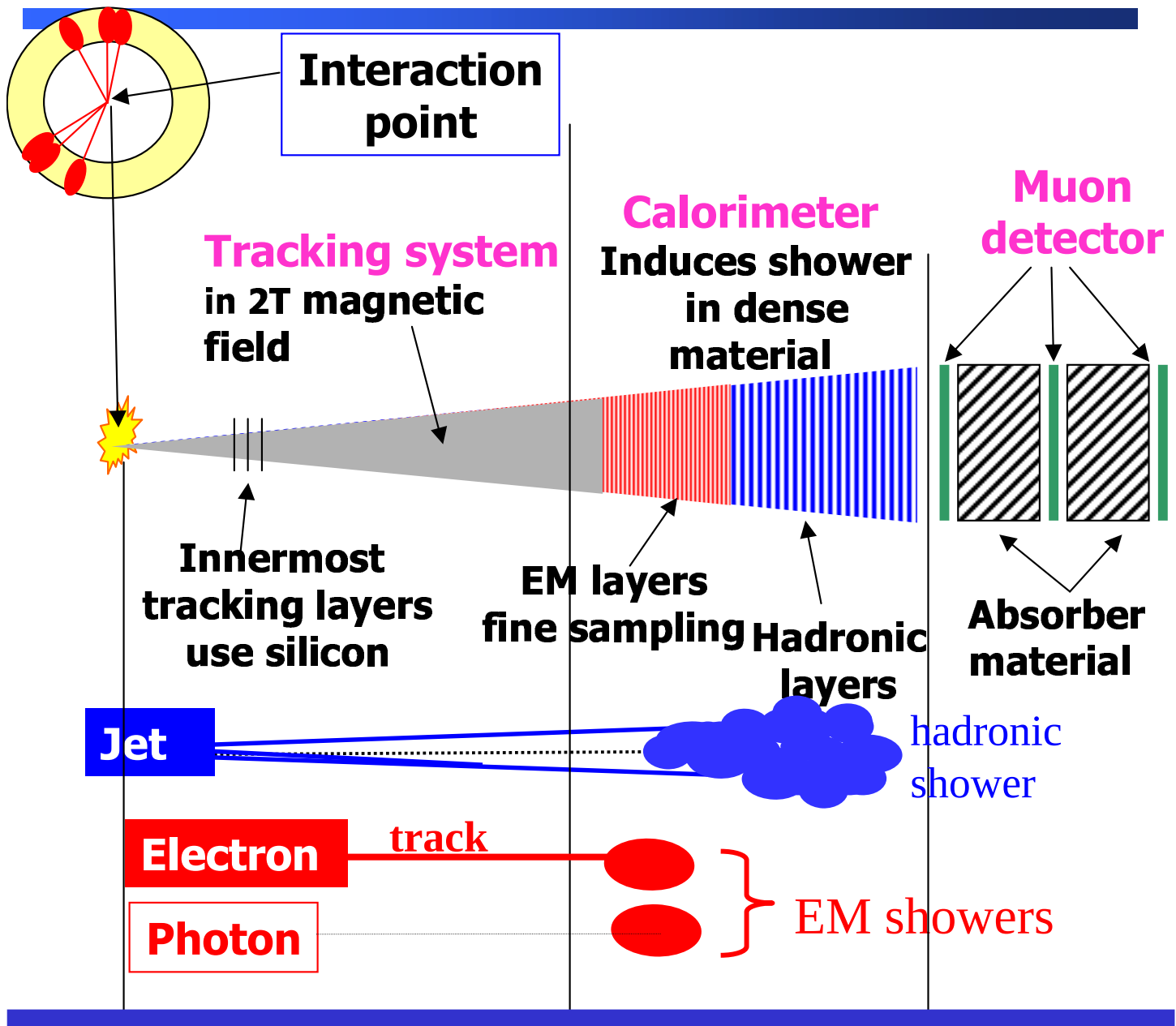
3. two hadronic jets misidentified as photons

e.g.



A Slice of the DØ Detector.

Experimental Signature of a Photon



Photon : EM-like shower in the calorimeter
WITHOUT associated track

Analysis Outline

- Select 2 Energetic Photon Objects
 - Estimate SM backgrounds:
 - QCD (jj and γj)
 - $Z/\gamma^* \rightarrow ee$
 - direct di-photons (PYTHIA MC with k factor of 1.4)
- from data with absolute normalization
- Look for a bump in the diphoton invariant mass spectrum

Selection of $\gamma\gamma$ Candidate Events

- **Trigger:**
di-EM* high p_T trigger
- **Offline:** (on both objects)
- **Kinematic cuts:** $p_T > 25\text{GeV}$
- **Acceptance cuts:** Central or End Cap
Calorimeter up to $|\eta|=2.4$
- **Photon ID:** -- shower shape consistent
with EM* shape (EMfraction,
Isolation, H-matrix χ^2)
-- sum of p_T of all tracks
in a hollow cone of $0.05 < R < 0.4$ around
photon candidate must be $< 2.0\text{ GeV}$
-- track veto
- *EM = *Electromagnetic Object (Photon or Electron)*

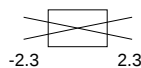
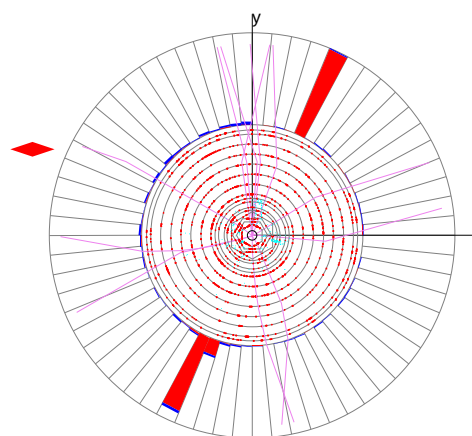
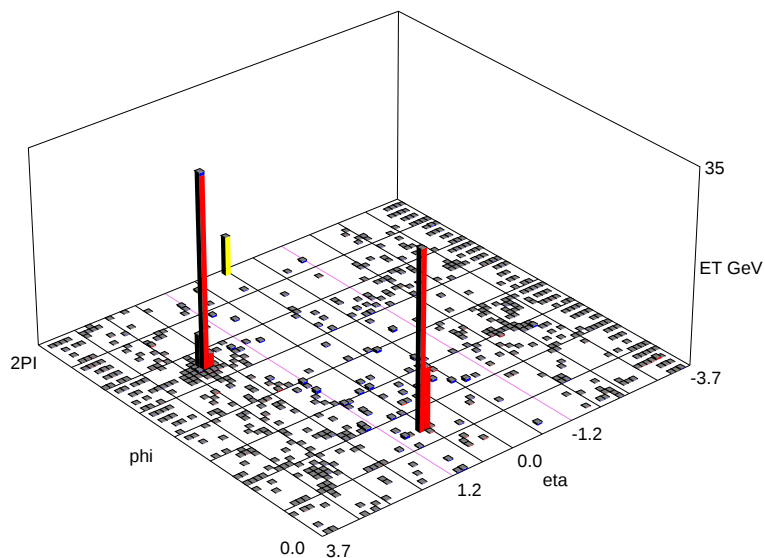
Event Displays of $\gamma\gamma$ Candidate

Run 148830 Event 3510187 Tue May 21 21:28:31 2002

14

Run 148830 Event 3510187 Tue May 21 21:28:43 2002

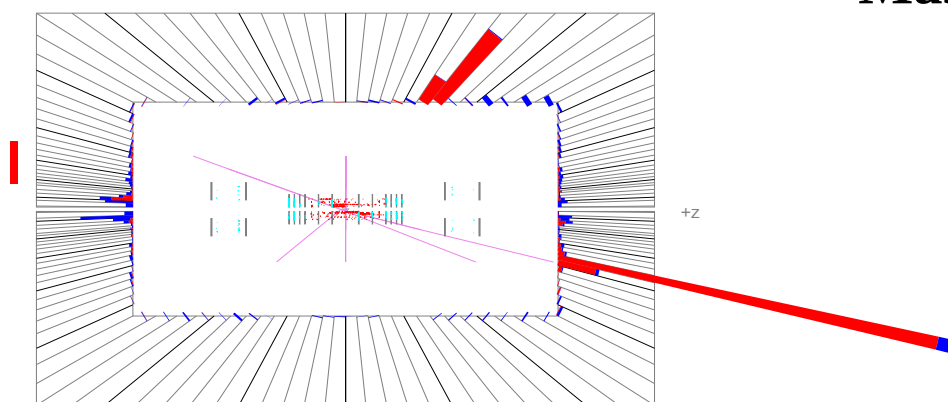
ET scale: 44 GeV



Run 148830 Event 3510187 Tue May 21 21:28:37 2002

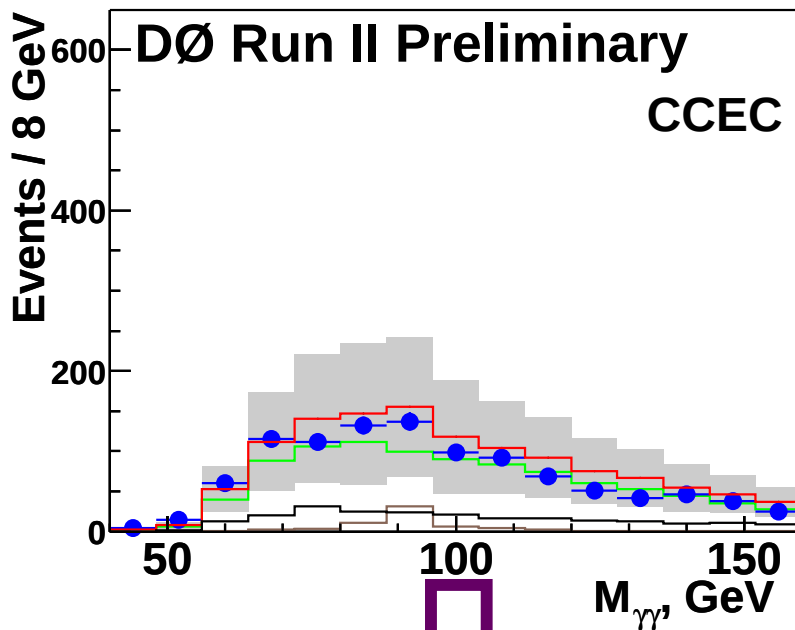
E scale: 42 GeV

Mass = 125.8 GeV



180 $\odot$ 0

Di-photon mass spectra, $\int L dt \approx 190 \text{ pb}^{-1}$ ($\approx$ half of the currently available data)



data = 1037.0

bkgd = 1211.5 $\pm$ 671.0

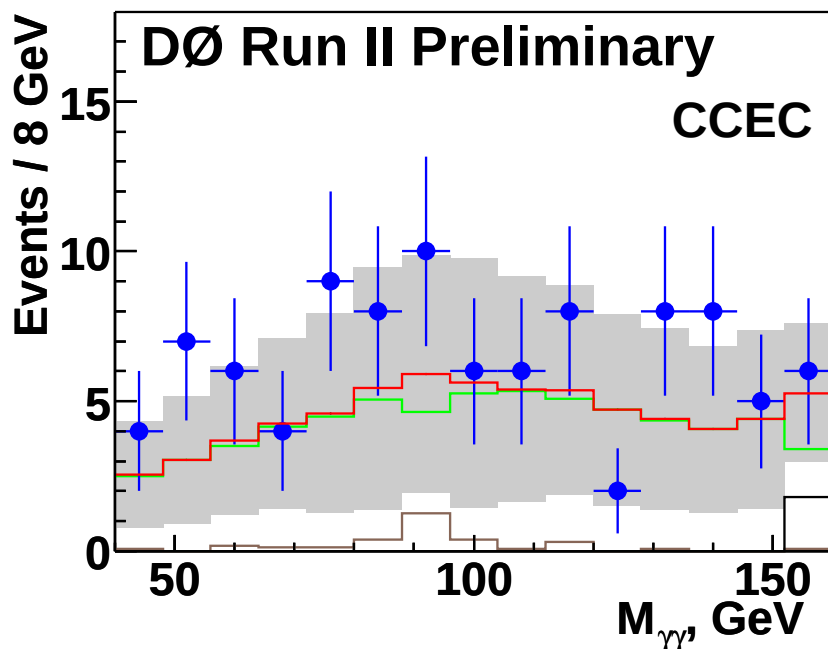
QCD = 922.1 $\pm$ 667.2

DY = 64.8 $\pm$ 69.4

$\gamma\gamma$ = 224.5 $\pm$ 16.6



Diphoton PT (35 GeV)cut



data = 97.0

bkgd = 68.8 $\pm$ 45.8

QCD = 64.0 $\pm$ 45.7

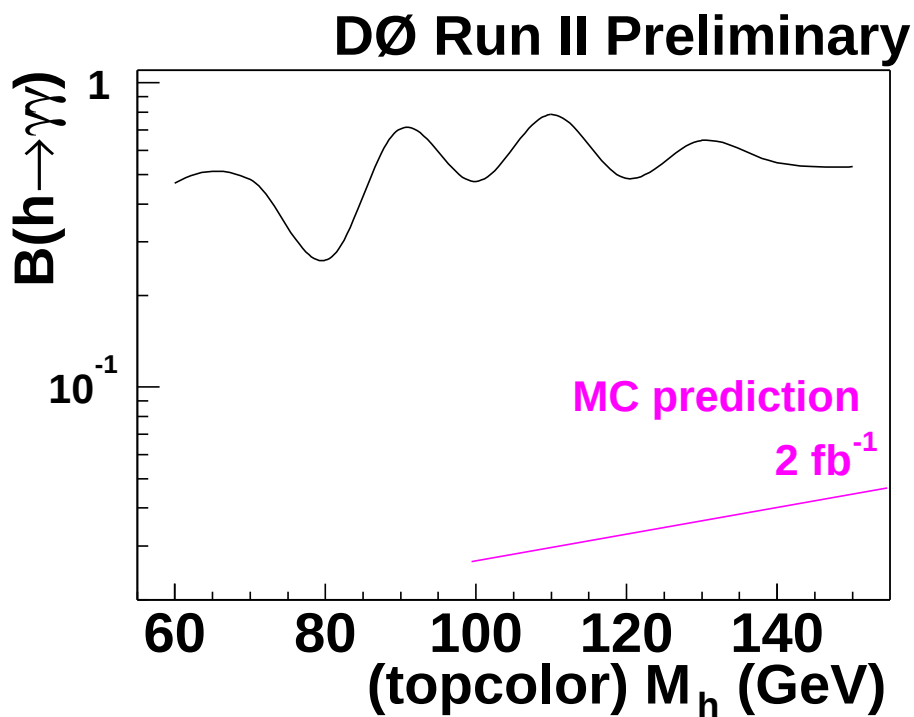
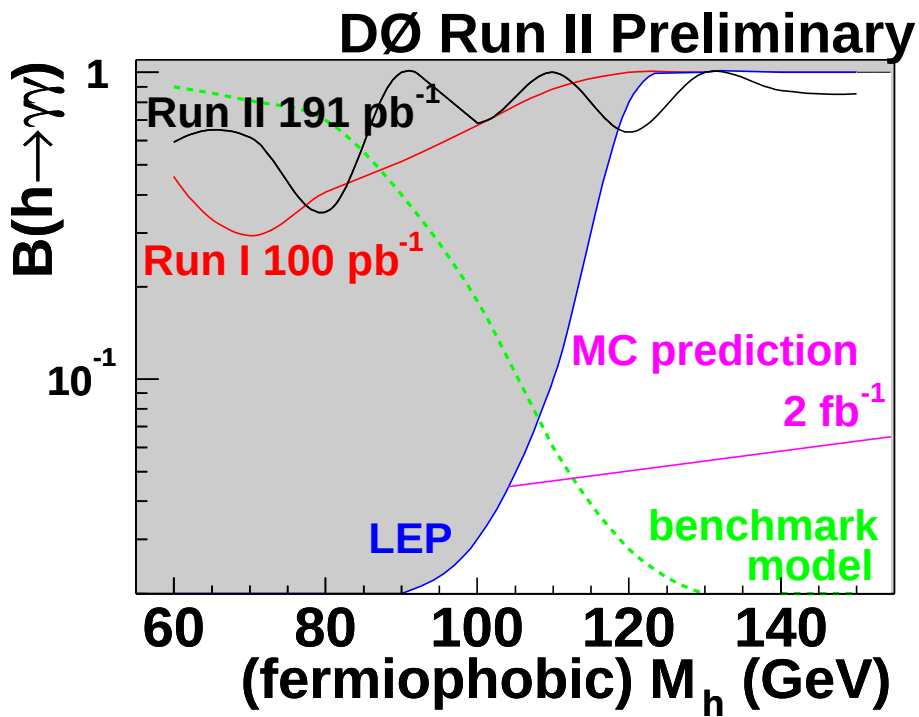
DY = 3.0 $\pm$ 3.0

$\gamma\gamma$ = 1.8 $\pm$ 0.1

Setting Limits on the Diphoton Branching Fraction

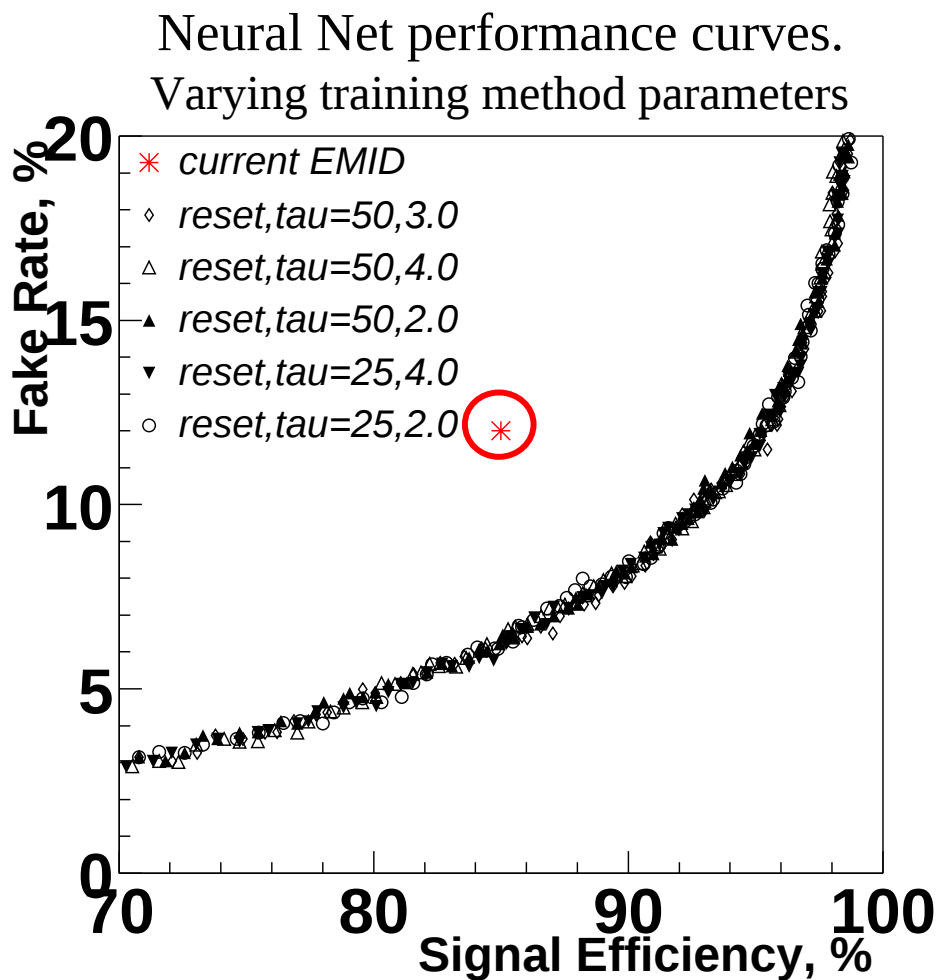
- Count di-photons in a mass window
*G.Landsberg , K.Matchev Phys. Rev. **D62**, 035004 (2000)*
- Derive Higgs cross section limits using Bayesian Approach
- Convert to limits on branching fractions
- Consider two scenarios:
 - ***Fermiophobic Higgs***
(does not couple to fermions)
Production mechanisms:
 1. W/Z associated production
 2. W/Z fusion
 - ***Topcolor Higgs***
(of all fermions couples only to top)
Production mechanisms:
 1. W/Z associated production
 2. W/Z fusion
 3. gluon fusion

$B(h \rightarrow \gamma\gamma)$ Limits



Neural Net Studies

- Try to use Neural Nets for photon identification to suppress background
- Preliminary studies indicate significant improvement over current EMID



- Plan to incorporate new ID in the analysis

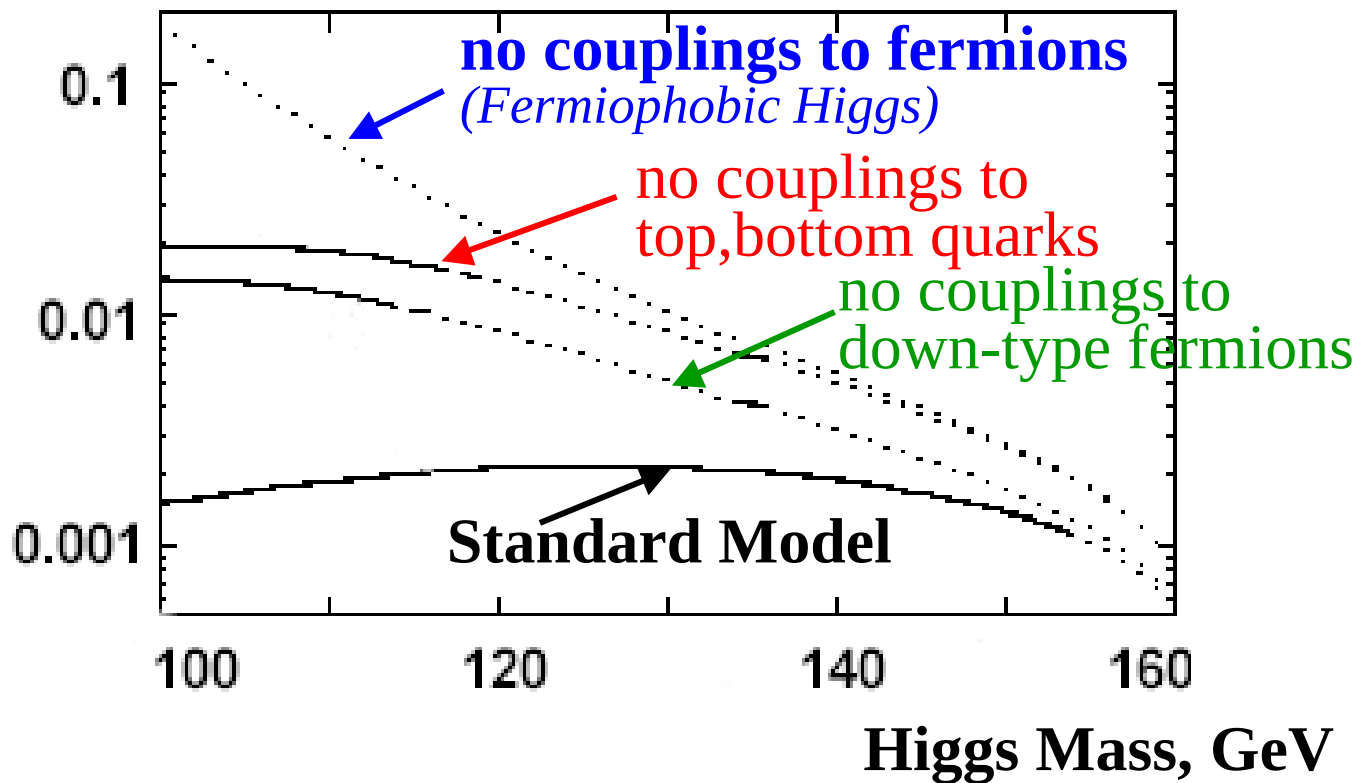
Conclusions

- A Search for the Fermiophobic and Topcolor Higgs boson has been performed in the inclusive $h \rightarrow \gamma\gamma + X$ channel with 190 pb^{-1} of DØ data collected between April 2002 and September 2003
- In the absence of excess we set limits on the branching fraction
- Limits are comparable with those of DØ Run I
- Work on improving sensitivity is in progress
- More data is being accumulated
- Stay tuned for new results

Back Up Slides Start Here

Examples of Enhancement of $h \rightarrow \gamma\gamma$ decays

$h \rightarrow \gamma\gamma$ Branching Fraction

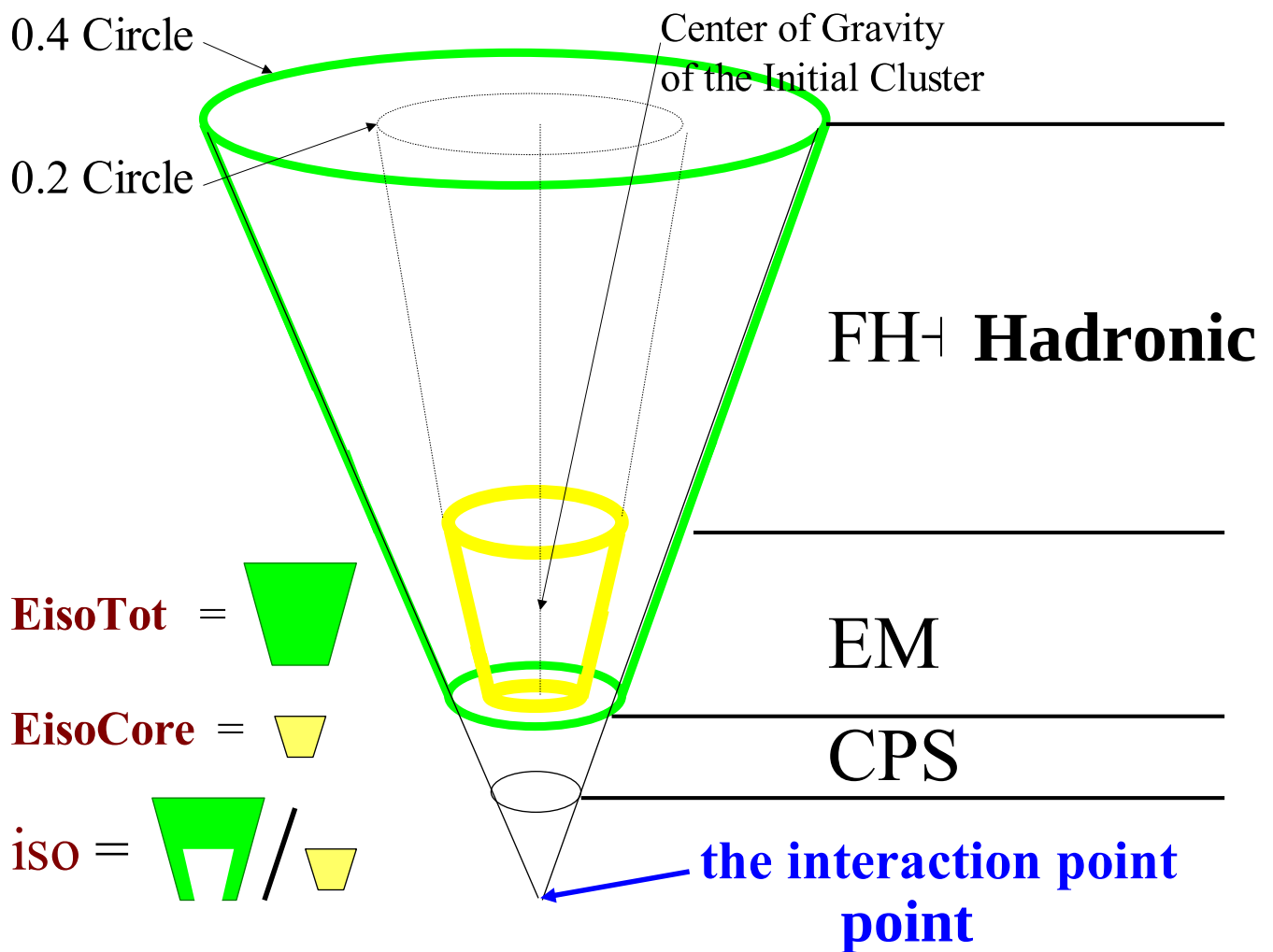


S.Mrenna, J.Wells, Phys. Rev. D63, 015006 (2001)

in general we should be prepared for any $h \rightarrow \gamma\gamma$ branching fraction (up to 1.0) due to new physics

Identification of a Photon Shower in the Calorimeter. Isolation

Photon-induced shower is smaller than quark/gluon shower both transversely and longitudinally

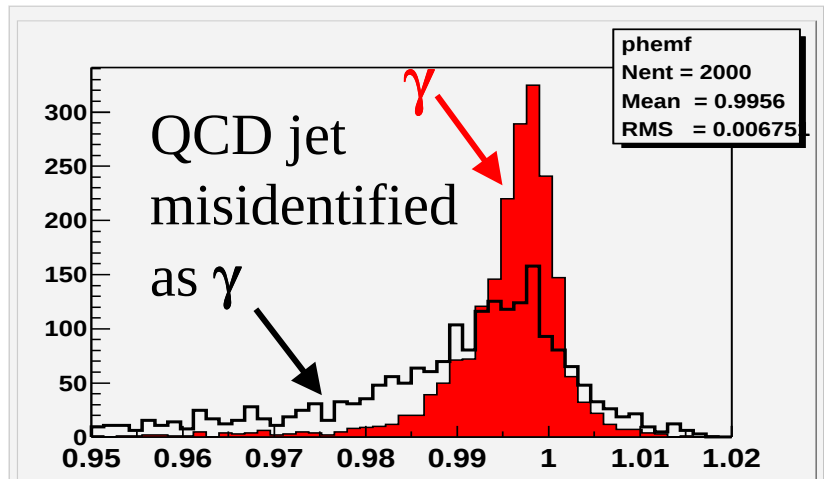


Photon ID Tools

(Monte Carlo Distributions)

EM fraction

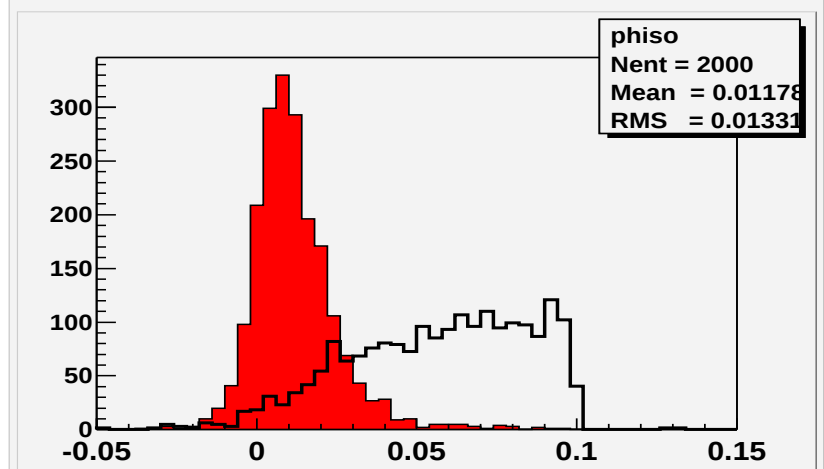
ratio of EM cluster energy deposited in EM calorimeter and total energy



Isolation

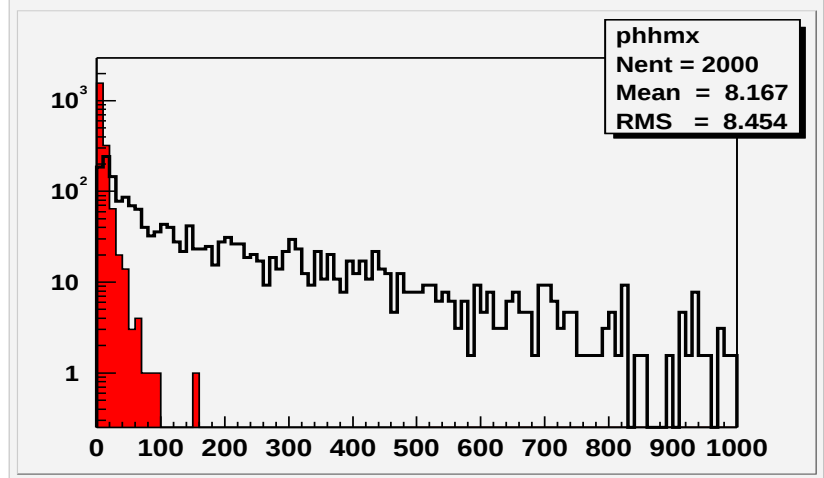
(previous slide)

measure of cluster narrowness

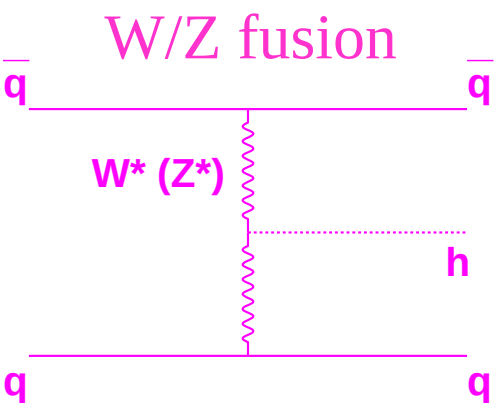
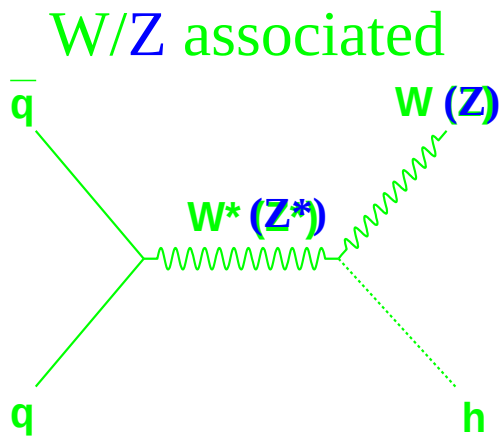
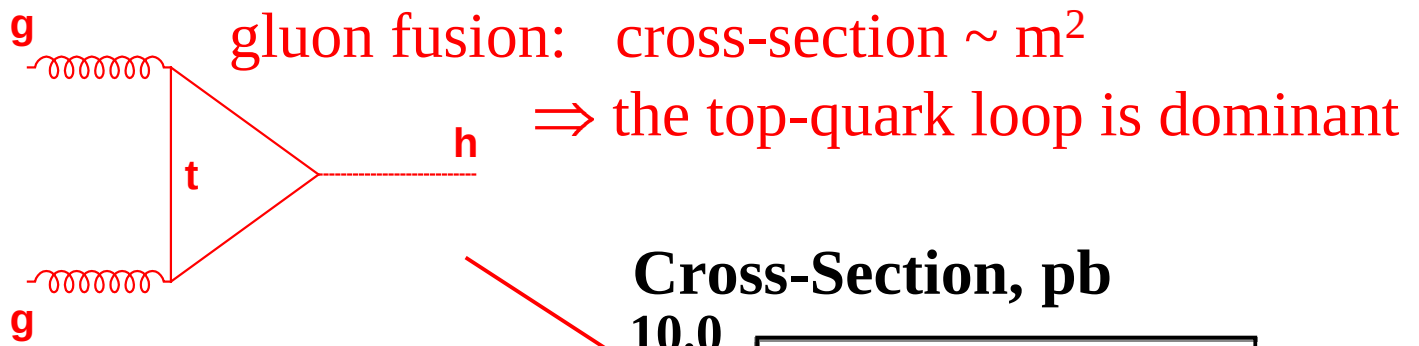


multi-variable shower shape tool

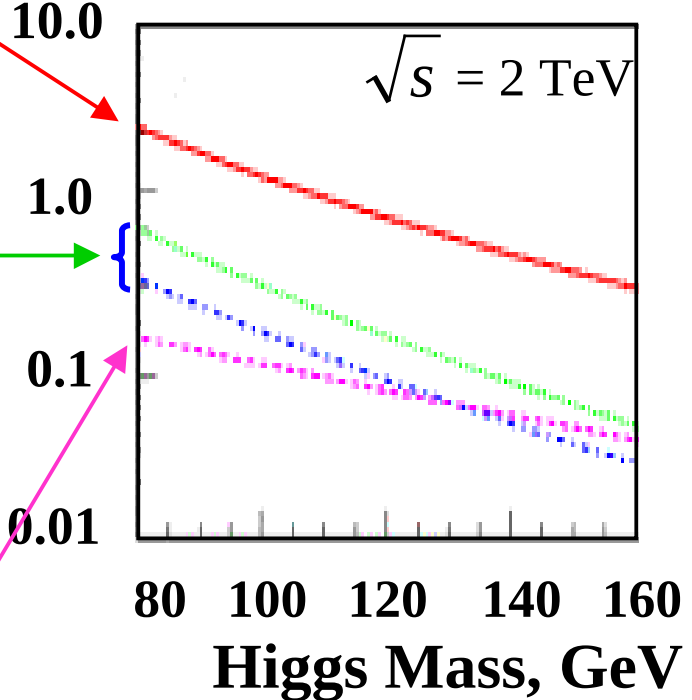
- layer energy fractions
- width at shower maximum



Leading SM Higgs Production Processes at Tevatron



Cross-Section, pb

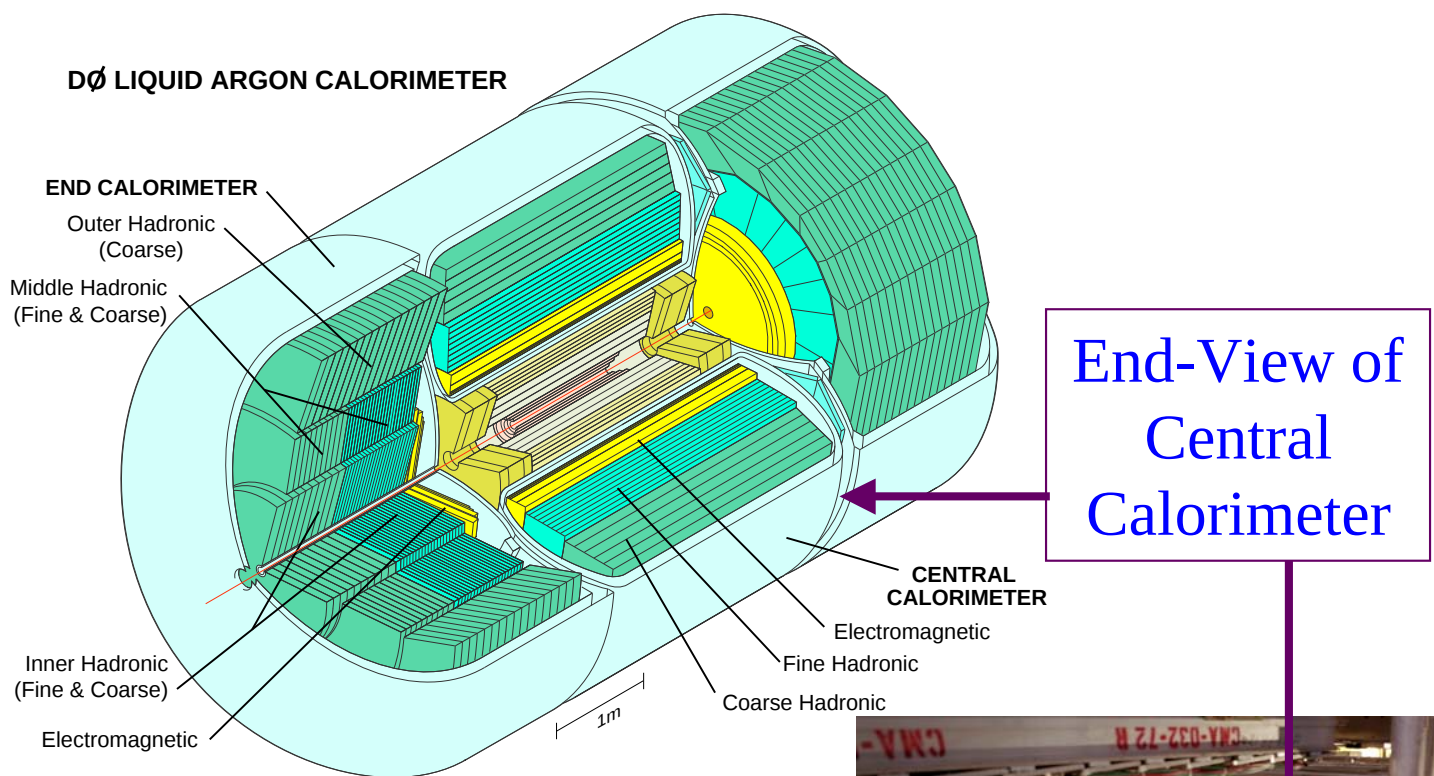


quark-antiquark fusion
 cross-section is small :

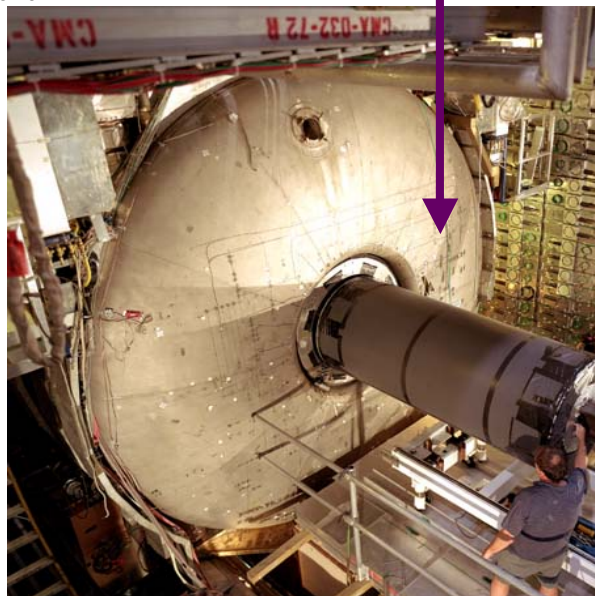
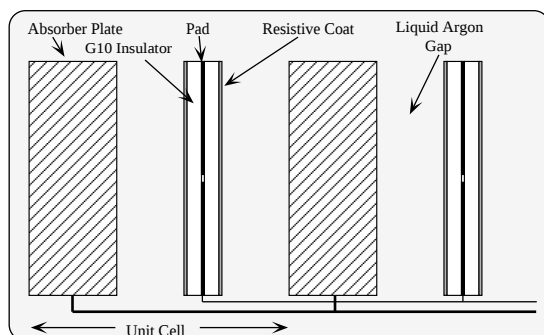
- Higgs-fermion coupling $\sim m_f$
- Masses of u,d quarks are small

DØ Calorimeter

- Uranium/Liquid Argon Sampling Calorimeter
- Three modules: -- Central Calorimeter (CC)
-- Two End Calorimeters (EC)

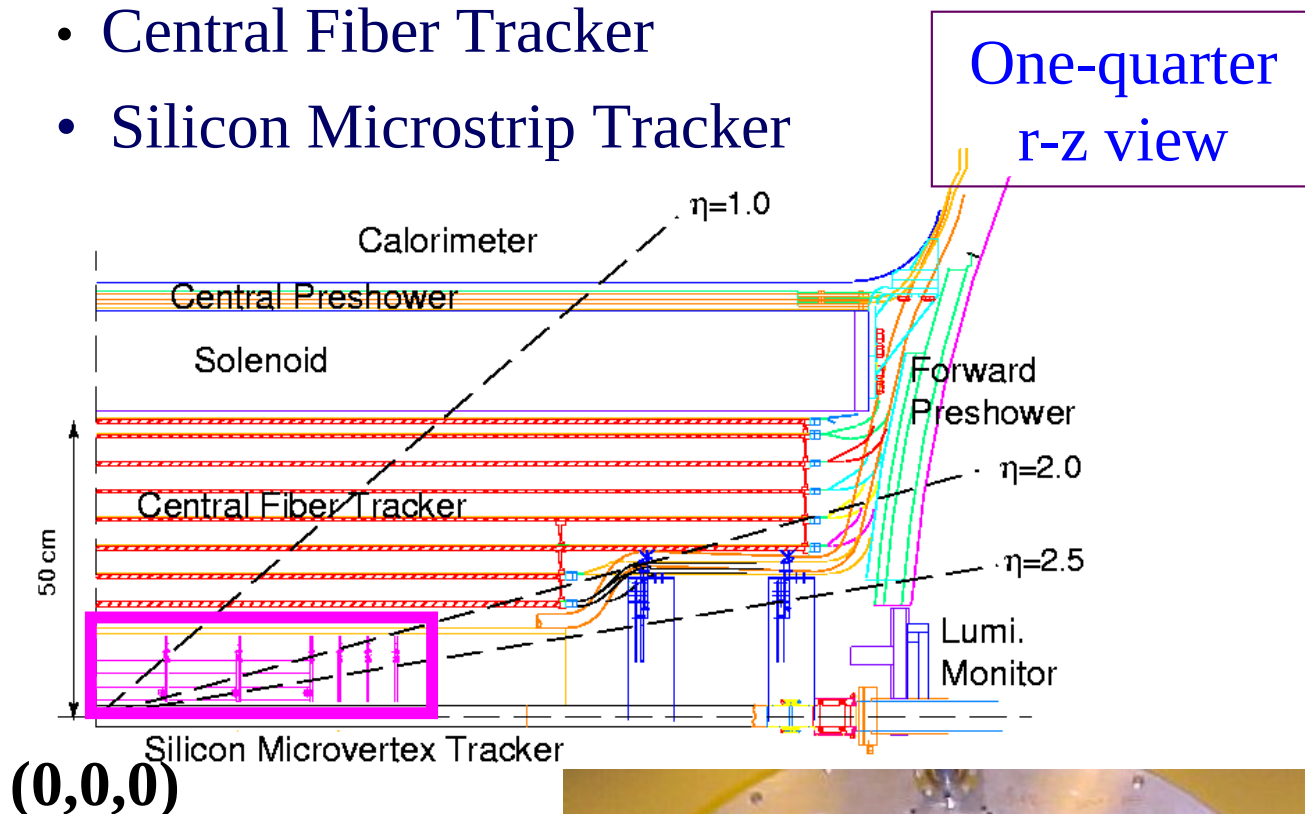


Unit cell

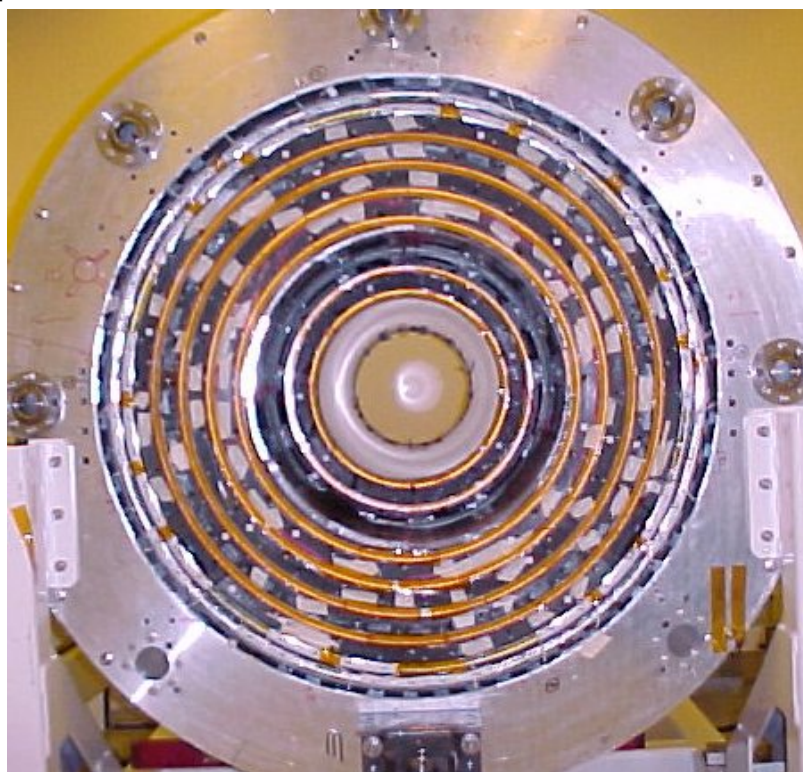


DØ Tracking System

- Central Fiber Tracker
- Silicon Microstrip Tracker

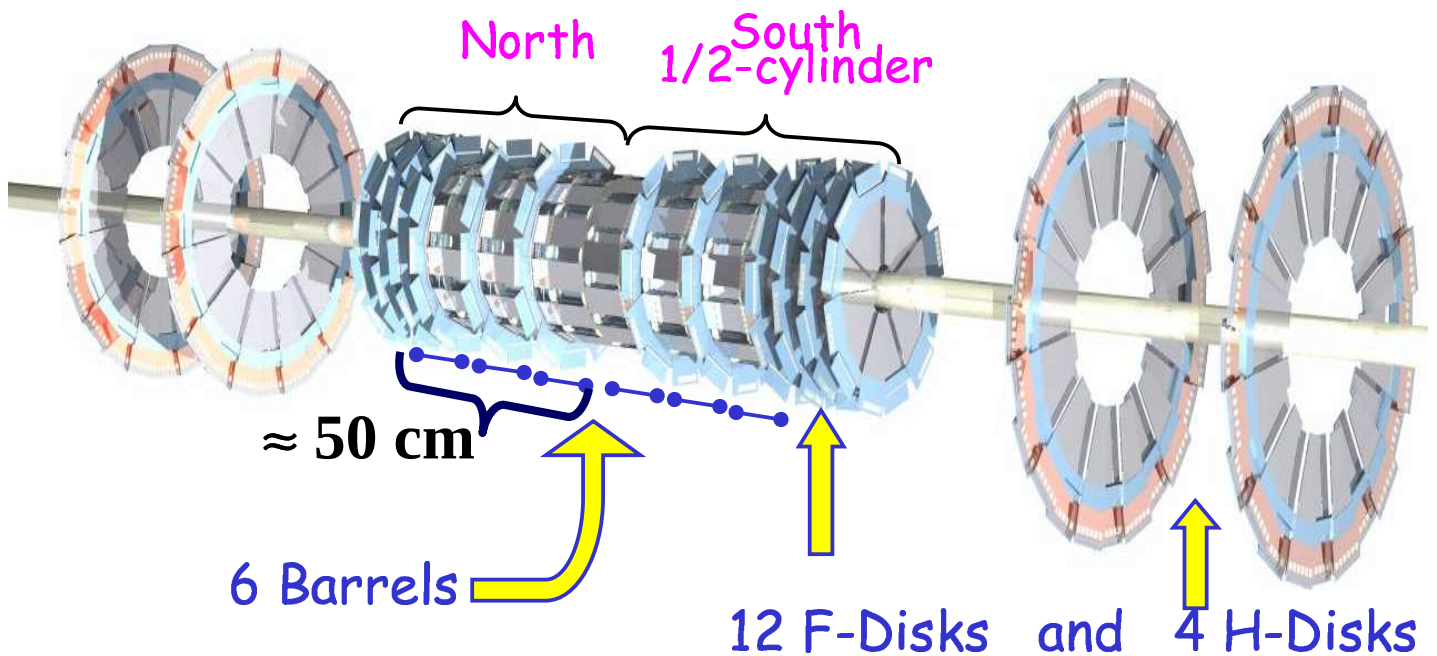
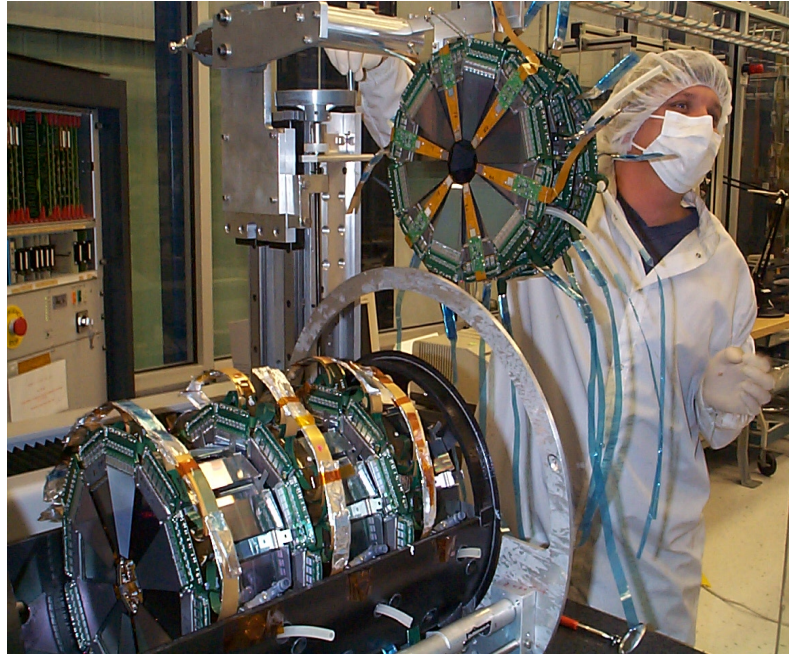


End-View
of the
Central Fiber
Tracker



Silicon Tracker. Longitudinal View

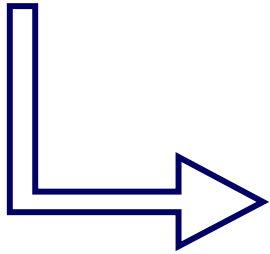
Barrel/Disk
Design:



Background Estimation

Example: $Z\gamma^* \rightarrow ee$

$$\left\{ \begin{aligned} N_{DY}(ee) &= \sigma(pp \rightarrow Z/\gamma^* + X) \times \text{Br}(ee) \times L_{\text{int}} \times A \times \epsilon(\text{trigger}) \\ &\quad \times \epsilon^2(\text{EM}) \times \epsilon^2(\text{track}) \\ N_{DY}(\gamma\gamma) &= \sigma(pp \rightarrow Z/\gamma^* + X) \times \text{Br}(ee) \times L_{\text{int}} \times A \times \epsilon(\text{trigger}) \\ &\quad \times \epsilon^2(\text{EM}) \times [1 - \epsilon(\text{track})]^2 \end{aligned} \right.$$



$$N_{DY}(\gamma\gamma) = N_{DY}(ee) \times \frac{[1 - \epsilon(\text{track})]^2}{\epsilon^2(\text{track})} =$$

$$= \{N(ee) - N(\text{misidentified } ee)\} \times \frac{[1 - \epsilon(\text{track})]^2}{\epsilon^2(\text{track})}$$

$$= \{N(ee) - N(\text{di-LooseEM}) \times f^2(e)\} \times \frac{[1 - \epsilon(\text{track})]^2}{\epsilon^2(\text{track})}$$

Electron misidentification rate

Tracking efficiency