

**A FIELD THEORETICAL MODEL
IN NONCOMMUTATIVE MINKOWSKI SUPERSPACE**

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Outline

- ★ **An Introduction**
- ★ **The Noncommutative SUSY Algebra**
- ★ **The Star Product**
- ★ **The Wess-Zumino Lagrangian**
 - ▶ **Chiral and Antichiral Superfields**
 - ▶ **Non-associativity and Weyl ordering**
 - ▶ **The Lagrangian**

I: Non(anti)commutativity from String Theory

★ Consider standard Nambu-Goto string action with extra term $B^{\mu\nu}$, which is constant.

$$S = \frac{1}{2\pi\alpha'} \int d\sigma d\tau \left(g^{ab} \eta_{\mu\nu} \partial_a x^\mu \partial_b x^\nu + \epsilon^{ab} B_{\mu\nu} \partial_a x^\mu \partial_b x^\nu \right)$$

► quantize string in the presence of a background field

$B^{\mu\nu} \Rightarrow$

$$[\hat{x}^\mu, \hat{x}^\nu] = 2\pi\alpha' \left((1 - B^2)^{-1} B \right)^{\mu\nu} \equiv i\Theta^{\mu\nu}.$$

I: Non(anti)commutativity from Superstring Theory

Ooguri, Vafa :hep-th/0302109, Seiberg :hep-th/0305248...

★ Consider superstrings in a self-dual graviphoton

background $F^{\alpha\beta}$, keep $F^{\dot{\alpha}\dot{\beta}} = 0 \Rightarrow$

work in flat Euclidean superspace \Rightarrow

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work in flat Euclidean superspace \Rightarrow

$$\{\hat{\theta}^\alpha, \hat{\theta}^\beta\} = (\alpha')^2 F^{\alpha\beta} \equiv C^{\alpha\beta},$$

$$\{\hat{\theta}^{\dot{\alpha}}, \hat{\theta}^{\dot{\beta}}\} = 0,$$

$$[\hat{y}^\mu, \hat{\theta}^\alpha] = [\hat{y}^\mu, \hat{\theta}^{\dot{\alpha}}] = [\hat{y}^\mu, \hat{y}^\nu] = 0$$

$$[\hat{x}^\mu, \hat{\theta}^\alpha] = iC^{\alpha\beta} \sigma_{\beta\dot{\alpha}}^\mu \bar{\theta}^{\dot{\alpha}}, \quad [\hat{x}^\mu, \hat{x}^\nu] = \bar{\theta}\bar{\theta}C^{\mu\nu}.$$

I: Non(anti)commutativity from Superstring Theory

★ Lorentz invariant Wess-Zumino lagrangian,
but nonhermitian

★ deformation of supersymmetry algebra is chirally
asymmetric (Berenstein, Rey: hep-th/0308049)

⇒ operators induced by deformation are nonhermitian,

⇒ assign non-canonical scaling dimensions for the Grass-
man coordinates, and superfields,

⇒ proof for renormalizability of 4-dimensional QFT's
with $\mathcal{N} = 1/2$ supersymmetry.

II: The Noncommutative SUSY Algebra

★ Consider 4D Minkowski space, $\bar{\theta}^{\dot{\alpha}} \equiv (\theta^{\alpha})^{\dagger}$,
 $\hat{y}^{\mu} \equiv \hat{x}^{\mu} + i\hat{\theta}\sigma^{\mu}\hat{\theta}$, $\hat{\bar{y}}^{\mu} \equiv \hat{x}^{\mu} - i\hat{\theta}\sigma^{\mu}\hat{\theta}$, $\bar{C}^{\dot{\alpha}\beta} = (C^{\beta\alpha})^{\dagger}$,

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Require $\{\hat{\theta}^{\alpha}, \hat{\theta}^{\beta}\} = C^{\alpha\beta} \quad \Rightarrow \quad \{\hat{\theta}^{\dot{\alpha}}, \hat{\theta}^{\dot{\beta}}\} = \bar{C}^{\dot{\alpha}\dot{\beta}}$

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Require $\{\hat{\theta}^{\dot{\alpha}}, \hat{\theta}^\alpha\} = 0 \quad \Rightarrow$

$$[\hat{y}^\mu, \hat{\theta}^\alpha] = 2iC^{\alpha\beta}\sigma^\mu_{\beta\dot{\beta}}\hat{\theta}^{\dot{\beta}},$$

$$[\hat{\bar{y}}^\mu, \hat{\theta}^{\dot{\alpha}}] = 2i\bar{C}^{\dot{\alpha}\dot{\beta}}\theta^\beta\sigma^\mu_{\beta\dot{\beta}}.$$

II: The Noncommutative SUSY Algebra

★
$$[\hat{y}^\mu, \hat{y}^\nu] - [\hat{\bar{y}}^\mu, \hat{\bar{y}}^\nu] = 4(\bar{C}^{\dot{\alpha}\dot{\beta}} \hat{\theta}^\alpha \hat{\theta}^\beta - C^{\alpha\beta} \hat{\bar{\theta}}^{\dot{\alpha}} \hat{\bar{\theta}}^{\dot{\beta}}) \sigma_{\alpha\dot{\alpha}}^\mu \sigma_{\beta\dot{\beta}}^\nu.$$

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$$[\hat{y}^\mu, \hat{y}^\nu] - [\hat{\bar{y}}^\mu, \hat{\bar{y}}^\nu] = 4(\bar{C}^{\dot{\alpha}\dot{\beta}}\hat{\theta}^\alpha\hat{\theta}^\beta - C^{\alpha\beta}\hat{\bar{\theta}}^{\dot{\alpha}}\hat{\bar{\theta}}^{\dot{\beta}})\sigma_{\alpha\dot{\alpha}}^\mu\sigma_{\beta\dot{\beta}}^\nu.$$

If
$$[\hat{y}^\mu, \hat{y}^\nu] = [\hat{\bar{y}}^\mu, \hat{\bar{y}}^\nu] = 0,$$

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If $[\hat{y}^\mu, \hat{y}^\nu] = [\hat{\bar{y}}^\mu, \hat{\bar{y}}^\nu] = 0,$

$\Rightarrow \bar{C}^{\dot{\alpha}\dot{\beta}}\hat{\theta}^\alpha\sigma_{\alpha\dot{\alpha}}^\mu\hat{\theta}^\beta\sigma_{\beta\dot{\beta}}^\nu = C^{\alpha\beta}\sigma_{\alpha\dot{\alpha}}^\mu\hat{\theta}^{\dot{\alpha}}\sigma_{\beta\dot{\beta}}^\nu\hat{\theta}^{\dot{\beta}}.$

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If $[\hat{y}^\mu, \hat{y}^\nu] = [\hat{\bar{y}}^\mu, \hat{\bar{y}}^\nu] = 0,$

$\Rightarrow \bar{C}^{\dot{\alpha}\dot{\beta}}\hat{\theta}^\alpha\sigma_{\alpha\dot{\alpha}}^\mu\hat{\theta}^\beta\sigma_{\beta\dot{\beta}}^\nu = C^{\alpha\beta}\sigma_{\alpha\dot{\alpha}}^\mu\hat{\bar{\theta}}^{\dot{\alpha}}\sigma_{\beta\dot{\beta}}^\nu\hat{\bar{\theta}}^{\dot{\beta}}.$

Thus, for consistency we choose,

$[\hat{y}^\mu, \hat{y}^\nu] = (4\bar{C}^{\dot{\alpha}\dot{\beta}}\hat{\theta}^\alpha\hat{\theta}^\beta - 2C^{\alpha\beta}\bar{C}^{\dot{\alpha}\dot{\beta}})\sigma_{\alpha\dot{\alpha}}^\mu\sigma_{\beta\dot{\beta}}^\nu,$

$[\hat{\bar{y}}^\mu, \hat{\bar{y}}^\nu] = (4C^{\alpha\beta}\hat{\bar{\theta}}^{\dot{\alpha}}\hat{\bar{\theta}}^{\dot{\beta}} - 2C^{\alpha\beta}\bar{C}^{\dot{\alpha}\dot{\beta}})\sigma_{\alpha\dot{\alpha}}^\mu\sigma_{\beta\dot{\beta}}^\nu$

$\Rightarrow [\hat{\bar{y}}^\mu, \hat{\bar{y}}^\nu] = 2C^{\alpha\beta}\bar{C}^{\dot{\alpha}\dot{\beta}}\sigma_{\alpha\dot{\alpha}}^\mu\sigma_{\beta\dot{\beta}}^\nu.$

II: The Noncommutative SUSY Algebra

$$\{\hat{\theta}^\alpha, \hat{\theta}^\beta\} = C^{\alpha\beta},$$

$$\{\hat{\theta}^{\dot{\alpha}}, \hat{\theta}^{\dot{\beta}}\} = \bar{C}^{\dot{\alpha}\dot{\beta}},$$

$$\{\hat{\theta}^{\dot{\alpha}}, \hat{\theta}^\alpha\} = 0,$$

$$[\hat{x}^\mu, \hat{\theta}^{\dot{\alpha}}] = i\bar{C}^{\dot{\alpha}\dot{\beta}}\hat{\theta}^{\dot{\beta}}\sigma_{\beta\dot{\beta}}^\mu,$$

$$[\hat{x}^\mu, \hat{\theta}^\alpha] = iC^{\alpha\beta}\sigma_{\beta\dot{\beta}}^\mu\hat{\theta}^{\dot{\beta}},$$

$$[\hat{x}^\mu, \hat{x}^\nu] = (C^{\alpha\beta}\hat{\theta}^{\dot{\alpha}}\hat{\theta}^{\dot{\beta}} - \bar{C}^{\dot{\alpha}\dot{\beta}}\hat{\theta}^{\dot{\beta}}\hat{\theta}^\alpha)\sigma_{\alpha\dot{\alpha}}^\mu\sigma_{\beta\dot{\beta}}^\nu.$$

II: The Star Product

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$$\begin{aligned}\theta^\alpha \theta^\beta &\Rightarrow \mathbf{W}(\hat{\theta}^\alpha \hat{\theta}^\beta), \\ &= -\hat{\theta}^\beta \hat{\theta}^\alpha + \frac{1}{2} C^{\alpha\beta}, \\ &\Rightarrow -\theta^\beta * \theta^\alpha + \frac{1}{2} C^{\alpha\beta}, \\ &= \theta^\alpha \theta^\beta,\end{aligned}$$

$$\begin{aligned}\hat{\theta}^\alpha \hat{\theta}^\beta &\Rightarrow \theta^\alpha * \theta^\beta, \\ &= \theta^\beta \theta^\alpha + \frac{1}{2} C^{\alpha\beta}, \\ &\Rightarrow -\mathbf{W}(\hat{\theta}^\beta \hat{\theta}^\alpha) + \frac{1}{2} C^{\alpha\beta}, \\ &= \hat{\theta}^\alpha \hat{\theta}^\beta.\end{aligned}$$

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★ If $\{\hat{\theta}^\alpha, \hat{\theta}^\beta\} = C^{\alpha\beta}$, and $\{\theta^\alpha, \theta^\beta\} = 0$, and choose Weyl ordering in noncommutative space, then

$$\begin{aligned}\theta^\alpha \theta^\beta &\Rightarrow \mathbf{W}(\hat{\theta}^\alpha \hat{\theta}^\beta), & \hat{\theta}^\alpha \hat{\theta}^\beta &\Rightarrow \theta^\alpha * \theta^\beta, \\ &= -\hat{\theta}^\beta \hat{\theta}^\alpha + \frac{1}{2}C^{\alpha\beta}, & &= \theta^\beta \theta^\alpha + \frac{1}{2}C^{\alpha\beta}, \\ &\Rightarrow -\theta^\beta * \theta^\alpha + \frac{1}{2}C^{\alpha\beta}, & &\Rightarrow -\mathbf{W}(\hat{\theta}^\beta \hat{\theta}^\alpha) + \frac{1}{2}C^{\alpha\beta}, \\ &= \theta^\alpha \theta^\beta, & &= \hat{\theta}^\alpha \hat{\theta}^\beta.\end{aligned}$$

★ An important requirement: reality condition

$$(f_1 * f_2)^\dagger = f_2^\dagger * f_1^\dagger$$

II: The Star Product

$$\begin{aligned}
 \hat{f}\hat{g} \Rightarrow f(x, \theta, \bar{\theta}) * g(x, \theta, \bar{\theta}) = fg + \\
 f \left[-\frac{C^{\alpha\beta} \overleftarrow{Q}_\alpha \overrightarrow{Q}_\beta - \bar{C}^{\dot{\alpha}\dot{\beta}} \overleftarrow{Q}_{\dot{\alpha}} \overrightarrow{Q}_{\dot{\beta}}}{2} \right. \\
 + \frac{C^{\alpha\beta} C^{\gamma\delta} \overleftarrow{Q}_\alpha \overleftarrow{Q}_\gamma \overrightarrow{Q}_\delta \overrightarrow{Q}_\beta}{8} + \frac{\bar{C}^{\dot{\alpha}\dot{\beta}} \bar{C}^{\dot{\gamma}\dot{\delta}} \overleftarrow{Q}_{\dot{\alpha}} \overleftarrow{Q}_{\dot{\gamma}} \overrightarrow{Q}_{\dot{\delta}} \overrightarrow{Q}_{\dot{\beta}}}{8} \\
 \left. + \frac{C^{\alpha\beta} \bar{C}^{\dot{\alpha}\dot{\beta}}}{4} \left(\overleftarrow{Q}_{\dot{\alpha}} \overleftarrow{Q}_\alpha \overrightarrow{Q}_{\dot{\beta}} \overrightarrow{Q}_\beta - \overleftarrow{Q}_\alpha \overleftarrow{Q}_{\dot{\alpha}} \overrightarrow{Q}_\beta \overrightarrow{Q}_{\dot{\beta}} \right) \right] g
 \end{aligned}$$

II: The Star Product

$$\begin{aligned} f(\theta) * g(\theta) &= f(\theta) \exp \left(-\frac{C^{\alpha\beta} \overleftarrow{\partial} \overrightarrow{\partial}}{2 \partial\theta^\alpha \partial\theta^\beta} \right) g(\theta) \\ &= f(\theta) \left(1 - \frac{C^{\alpha\beta} \overleftarrow{\partial} \overrightarrow{\partial}}{2 \partial\theta^\alpha \partial\theta^\beta} - \mathbf{det}C \frac{\overleftarrow{\partial} \overrightarrow{\partial}}{\partial\theta\theta \partial\theta\theta} \right) g(\theta), \end{aligned}$$

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and

$$\begin{aligned} f(\bar{\theta}) * g(\bar{\theta}) &= f(\bar{\theta}) \exp \left(-\frac{\bar{C}^{\dot{\alpha}\dot{\beta}} \overleftarrow{\partial} \overrightarrow{\partial}}{2 \partial\bar{\theta}^{\dot{\alpha}} \partial\bar{\theta}^{\dot{\beta}}} \right) g(\bar{\theta}) \\ &= f(\bar{\theta}) \left(1 - \frac{\bar{C}^{\dot{\alpha}\dot{\beta}} \overleftarrow{\partial} \overrightarrow{\partial}}{2 \partial\bar{\theta}^{\dot{\alpha}} \partial\bar{\theta}^{\dot{\beta}}} - \mathbf{det}\bar{C} \frac{\overleftarrow{\partial} \overrightarrow{\partial}}{\partial\bar{\theta}\bar{\theta} \partial\bar{\theta}\bar{\theta}} \right) g(\bar{\theta}), \end{aligned}$$

II: The Star Product

$$\begin{aligned}
 \{\theta^\alpha, \theta^\beta\}_* &= C^{\alpha\beta}, & [x^\mu, \theta^\alpha]_* &= iC^{\alpha\beta} \sigma_{\beta\dot{\beta}}^\mu \bar{\theta}^{\dot{\beta}}, \\
 \{\bar{\theta}^{\dot{\alpha}}, \bar{\theta}^{\dot{\beta}}\}_* &= \bar{C}^{\dot{\alpha}\dot{\beta}}, & [x^\mu, \bar{\theta}^{\dot{\alpha}}]_* &= i\bar{C}^{\dot{\alpha}\dot{\beta}} \theta^{\beta} \sigma_{\beta\dot{\beta}}^\mu, \\
 \{\bar{\theta}^{\dot{\alpha}}, \theta^\alpha\}_* &= 0, & [x^\mu, x^\nu]_* &= \bar{\theta}\bar{\theta}C^{\mu\nu} + \theta\theta\bar{C}^{\mu\nu}.
 \end{aligned}$$

where,

$$\begin{aligned}
 C^{\mu\nu} &\equiv C^{\alpha\beta} \epsilon_{\beta\gamma} (\sigma^{\mu\nu})_\alpha{}^\gamma, \\
 \bar{C}^{\mu\nu} &\equiv \bar{C}^{\dot{\alpha}\dot{\beta}} \epsilon_{\dot{\beta}\dot{\gamma}} (\bar{\sigma}^{\mu\nu})^{\dot{\gamma}}{}_{\dot{\alpha}}.
 \end{aligned}$$

III: Supercharges and Covariant Derivatives

★ Q 's and D 's have their canonical forms,

$$\{Q_\alpha, Q_\beta\} = -4\bar{C}^{\dot{\alpha}\dot{\beta}}\sigma_{\alpha\dot{\alpha}}^\mu\sigma_{\beta\dot{\beta}}^\nu\frac{\partial^2}{\partial\bar{y}^\mu\partial\bar{y}^\nu},$$

$$\{\bar{Q}_{\dot{\alpha}}, \bar{Q}_{\dot{\beta}}\} = -4C^{\alpha\beta}\sigma_{\alpha\dot{\alpha}}^\mu\sigma_{\beta\dot{\beta}}^\nu\frac{\partial^2}{\partial y^\mu\partial y^\nu},$$

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$$\{\vec{Q}_\alpha, \vec{Q}_{\dot{\alpha}}\} = 2i\sigma_{\alpha\dot{\alpha}}^\mu\frac{\partial}{\partial y^\mu}.$$

★ Supersymmetry is broken by the star product,

$$\vec{Q}(\Psi_1 * \Psi_2) \neq \vec{Q}(\Psi_1) * \Psi_2 + (-1)^{p(\Psi_1)}\Psi_1 * \vec{Q}(\Psi_2),$$

$$\vec{\bar{Q}}(\Psi_1 * \Psi_2) \neq \vec{\bar{Q}}(\Psi_1) * \Psi_2 + (-1)^{p(\Psi_1)}\Psi_1 * \vec{\bar{Q}}(\Psi_2).$$

III: Supercharges and Covariant Derivatives

$$\{D_\alpha, D_\beta\} = 0,$$

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$$\star \quad \{D_\alpha, Q_\beta\} = \{\bar{D}_{\dot{\alpha}}, Q_\beta\} = \{D_\alpha, \bar{Q}_{\dot{\beta}}\} = \{\bar{D}_{\dot{\alpha}}, \bar{Q}_{\dot{\beta}}\} = 0.$$

Can keep the same definitions for $\Phi(y, \theta)$, and $\bar{\Phi}(\bar{y}, \bar{\theta})$,

$$\bar{D}_{\dot{\alpha}}\Phi(y, \theta) = 0,$$

$$D_\alpha\bar{\Phi}(\bar{y}, \bar{\theta}) = 0.$$

IV: Chiral and Antichiral Superfields

$$\begin{aligned}\Phi(\hat{y}, \hat{\theta}) &= A(\hat{y}) + \sqrt{2}\hat{\theta}\psi(\hat{\theta}) + \hat{\theta}\hat{\theta}F(\hat{y}), \\ \bar{\Phi}(\hat{y}, \hat{\theta}) &= A(\hat{y}) + \sqrt{2}\hat{\theta}\bar{\psi}(\hat{y}) + \hat{\theta}\hat{\theta}\bar{F}(\hat{y}).\end{aligned}$$

For the star product of two chiral fields we find,

$$\begin{aligned}\Phi_1(y, \theta) * \Phi_2(y, \theta) &= \Phi_1(y, \theta)\Phi_2(y, \theta) - C^{\alpha\beta}\psi_{1\alpha}\psi_{2\beta} - \underline{\det C F_1 F_2} \\ &\quad + \sqrt{2}\theta^\gamma C^{\alpha\beta} [\epsilon_{\beta\gamma}(\psi_{1\alpha}F_2 - \psi_{2\alpha}F_1) \\ &\quad + \bar{C}^{\dot{\alpha}\dot{\beta}}\sigma_{\alpha\dot{\alpha}}^\mu\sigma_{\gamma\dot{\beta}}^\nu(\partial_\mu A_1\partial_\nu\psi_{2\beta} - \partial_\mu A_2\partial_\nu\psi_{1\beta})] \\ &\quad + \theta\theta [2\bar{C}^{\mu\nu}\partial_\mu A_1\partial_\nu A_2 \\ &\quad + C^{\alpha\beta}\bar{C}^{\dot{\alpha}\dot{\beta}}\sigma_{\alpha\dot{\alpha}}^\mu\sigma_{\beta\dot{\beta}}^\nu(\partial_\mu A_1\partial_\nu F_2 - \partial_\mu A_2\partial_\nu F_1)],\end{aligned}$$

V: Non-associativity and Weyl ordering

$$\Phi_1 * \Phi_2 \neq \Phi_2 * \Phi_1$$

$$\bar{\Phi}_1 * \bar{\Phi}_2 \neq \bar{\Phi}_2 * \bar{\Phi}_1$$

V: Non-associativity and Weyl ordering

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★ limit discussion to symmetrized via Weyl ordering products of superfields. (Seiberg: hep-th/0305248)

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$$\Phi_1 * (\Phi_2 * \Phi_3) \neq (\Phi_1 * \Phi_2) * \Phi_3$$

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★ limit discussion to symmetrized via Weyl ordering products of superfields. (Seiberg: hep-th/0305248)

$$\Phi_1 * (\Phi_2 * \Phi_3) \neq (\Phi_1 * \Phi_2) * \Phi_3$$

★ Double Weyl ordering

VI: The Wess-Zumino Lagrangian

★ We find the following result for the Wess-Zumino lagrangian

$$\begin{aligned}\mathcal{L} &= \mathbf{w} \left[\int d^2\theta d^2\bar{\theta} \bar{\Phi} * \Phi + \int d^2\theta \left(\frac{1}{2}m\Phi * \Phi + \frac{1}{3}g\Phi * \Phi * \Phi \right) \right. \\ &\quad \left. + \int d^2\bar{\theta} \left(\frac{1}{2}\bar{m}\bar{\Phi} * \bar{\Phi} + \frac{1}{3}\bar{g}\bar{\Phi} * \bar{\Phi} * \bar{\Phi} \right) \right] \\ &= \mathcal{L}(C = 0) - \underline{\underline{\frac{1}{3}\det C F^3 - \frac{1}{3}\det \bar{C} \bar{F}^3}} + \text{total derivatives.}\end{aligned}$$

VII: Summery

- ★ We worked out a consistent algebra in Minkowski space,
- ★ Presented a star product that reproduces the algebra of deformed parameter space in its entirety, satisfies the reality condition, and preserves the chirality of a product of superfields,
- ★ Obtained Wess-Zumino lagrangian that is both Lorentz invariant, and Hermitian.