

Precise Measurement of the Differential Cross Section for $\nu - Fe$ and $\bar{\nu} - Fe$ Scattering

Martin Tzanov

Introduction

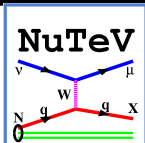
- NuTeV experiment
- $\nu(\bar{\nu}) - N$ Differential cross section and structure functions

Method

- Differential cross section extraction
- F_2 and xF_3 measurement

Results

Conclusions



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NuTeV Collaboration

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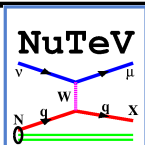
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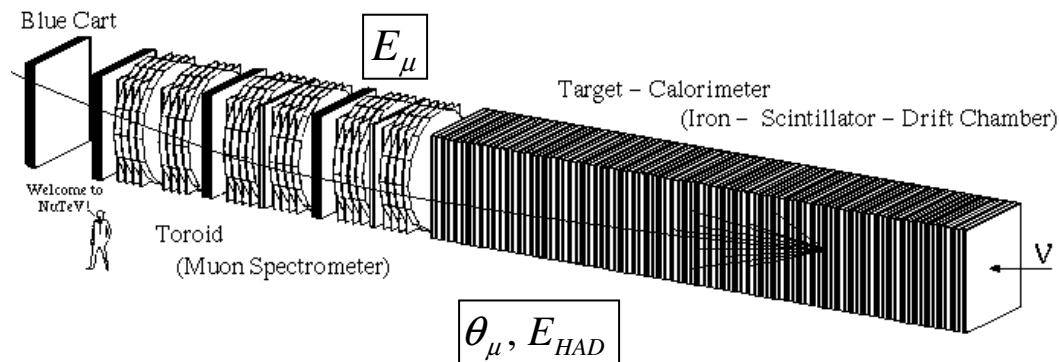
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NuTeV Detector



Target Calorimeter:

- Steel-Scintillator Sandwich (10 cm)

$$\frac{\delta E}{E} \approx \frac{0.86}{\sqrt{E}} \text{ -resolution}$$

- Tracking chambers for muon track and vertex

New feature: Continuous Calibration Beam

Hadron energy scale $\frac{\Delta E_{HAD}}{E_{HAD}} = 0.43\%$

Muon Spectrometer:

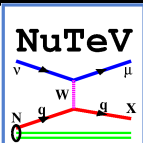
- Three toroidal iron magnets with five sets of drift chambers

$$\langle B_\phi \rangle \approx 1.7T, p_t \approx 2.4 GeV/c$$

$$\delta(1/p)/(1/p) \sim 11\% \text{ MCS dominated}$$

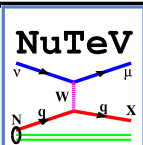
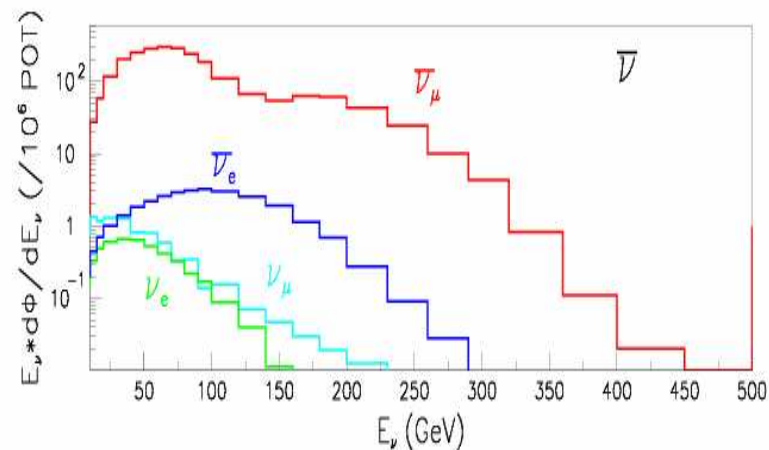
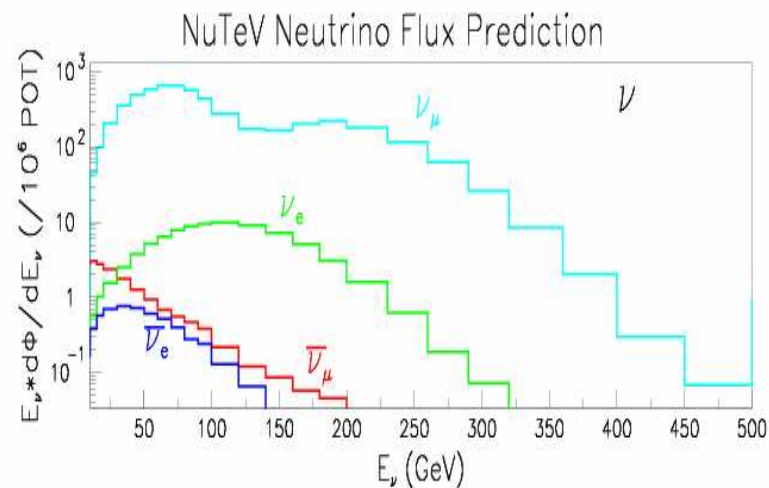
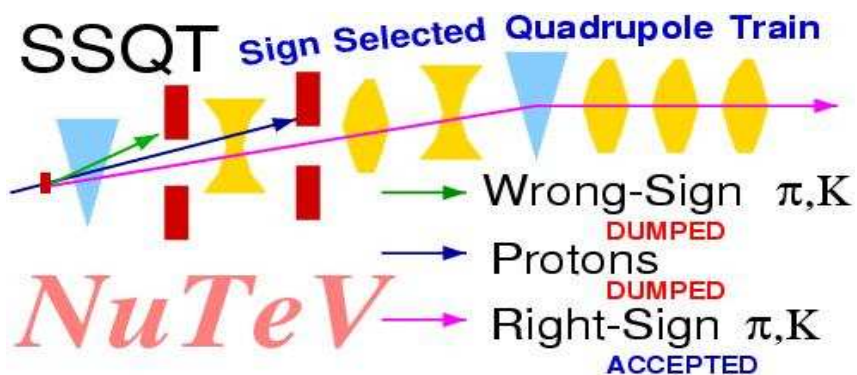
- Always focusing for lead muon

Muon energy scale $\frac{\Delta E_\mu}{E_\mu} = 0.7\%$

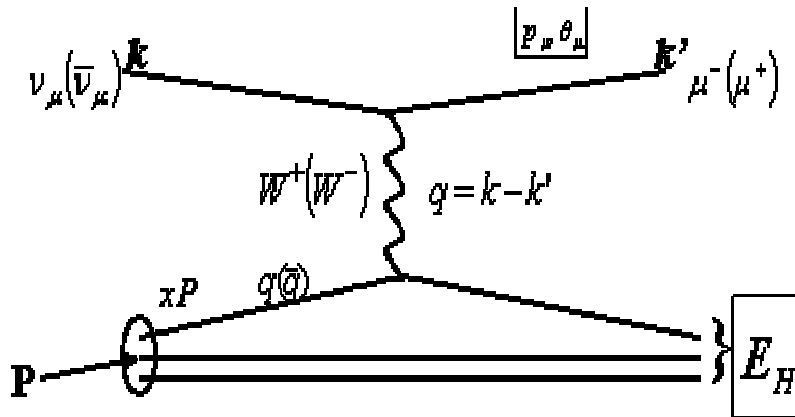


Neutrino Beamline

- NuTeV accumulated over 3 million neutrino/antineutrino events in 1996-1997
 - $20 \leq E_\nu \leq 400 \text{ GeV}$
- Wrong-sign π^\pm/K^\pm are dumped
 - NuTeV selects neutrinos or antineutrinos
- High purity separate neutrino and antineutrinos
 - high y CC sample
 - tag leading muon



Neutrino-Nucleon Scattering



$$Q^2 = 4E_\nu E_\mu \sin^2 \frac{\theta}{2}, \quad \text{Squared 4-momentum transferred to hadronic system}$$

$$x = \frac{Q^2}{2ME_{HAD}}, \quad \text{Fraction of momentum carried by the struck quark}$$

$$y = \frac{\nu}{E_\nu} = \frac{E_{HAD}}{E_\nu}, \quad \text{Inelasticity}$$

Differential cross section in terms of structure functions:

$$\frac{1}{E_\nu} \frac{d^2\sigma^{\nu(\bar{\nu})}}{dx dy} = \frac{G_F^2 M}{\pi(1+Q^2/M_W^2)^2} \left[\left(1 - y - \frac{Mxy}{2E_\nu} + \frac{y^2}{2} \frac{1+4M^2x^2/Q^2}{1+R(x,Q^2)} \right) F_2^{\nu(\bar{\nu})} \pm \left(y - \frac{y^2}{2} \right) xF_3^{\nu(\bar{\nu})} \right]$$

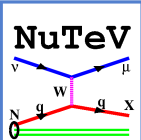
Structure Functions in terms of parton distributions

$$F_2^{\nu(\bar{\nu})N} = \sum [xq^{\nu(\bar{\nu})N}(x) + x\bar{q}^{\nu(\bar{\nu})N}(x) + 2xk^{\nu(\bar{\nu})N}(x)]$$

$$xF_3^{\nu(\bar{\nu})N} = x(d_\nu(x) + u_\nu(x)) \pm 2x(s(x) - c(x)), \quad \text{v-scattering only}$$

$$\Delta xF_3 = x(F_3^\nu - F_3^{\bar{\nu}}) = 4x(s(x) - c(x))$$

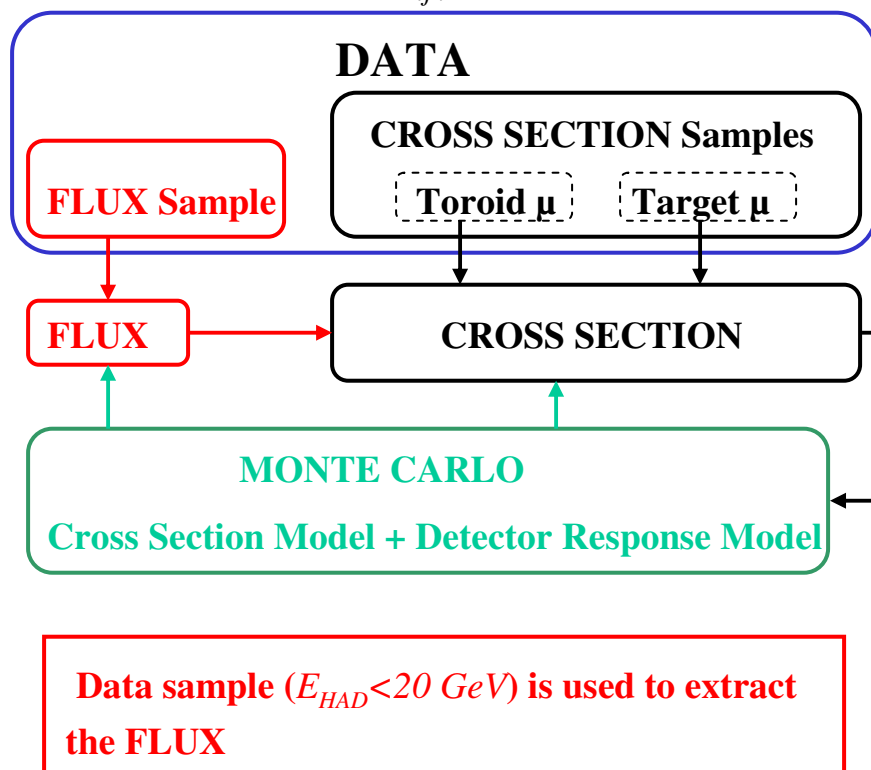
$$R = \frac{\sigma_L}{\sigma_T}$$



Extraction of the Differential Cross Section

Differential Cross Section in terms of number of events and the flux:

$$\left(\frac{1}{E} \frac{d^2\sigma}{dx dy} \right)_{ijk}^{\nu(\bar{\nu})} = \frac{1}{\Phi(E_k)} \frac{N_{ijk}^{\nu(\bar{\nu})}}{\Delta x_i \Delta y_j}$$



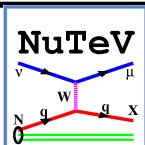
Events selection criteria:

- **Toroid muon** : Good muon track, containment, $E_\mu > 15 \text{ GeV}$, $E_{HAD} > 10 \text{ GeV}$, $30 < E_\nu < 300 \text{ GeV}$ $Q^2 > 1 \text{ GeV}^2/c^2$
- **Target muon**: $4 < E_\mu < 12 \text{ GeV}$

Monte Carlo event generator is used for acceptance:

- **Cross section model**:
 - QCD inspired fit using Buras – Gaemers¹ param.
 - Q^2 evolution from GRV for $Q^2 < 1.35 \text{ GeV}^2/c^2$. Includes fit for $x > 0.4$ to SLAC, NMC and BCDMS to account for non-perturbative behavior at low Q^2 and high x
- **Detector model**:
 - E_μ and E_{HAD} resolution functions are parameterized using Test Beam (TB) muons and hadrons
 - θ_μ is parameterized as a function of E_{HAD} and event position using GEANT hit level Monte Carlo

¹(A. Buras, K. Gaemers, Nucl. Phys. **B132**, 249 (1978))



Flux Extraction

Relative flux extraction using “fixed ν_0 method”

- We can rewrite the differential cross section in terms of $\nu = E_{HAD}$ integrate over x .

$$\frac{d^2\sigma^{\nu(\bar{\nu})}}{dx d\nu} = \frac{G^2 M}{\pi} \left[F_2 - \frac{\nu}{E} (F_2 \mp x F_3) + \frac{\nu^2}{2E^2} (F_2 \mp x F_3 + R_{TERM} F_2) \right],$$

$$\text{where } R_{TERM}(x, Q^2) = \frac{1 + 2Mx/\nu}{1 + R_L(x, Q^2)} - \frac{Mx}{\nu} - 1$$

$$\frac{dN_{\nu, \bar{\nu}}}{d\nu} = \Phi(E_\nu) \left(A_{\nu, \bar{\nu}} + B_{\nu, \bar{\nu}} (\nu/E_\nu) + \frac{C_{\nu, \bar{\nu}}}{2} (\nu/E_\nu)^2 \right),$$

$$A = \frac{G_F^2 M}{\pi} \int_0^1 F_2(x, Q^2) dx, \quad B = -\frac{G_F^2 M}{\pi} \int_0^1 [F_2(x, Q^2) \mp x F_3(x, Q^2)] dx,$$

$$C = B - \frac{G_F^2 M}{\pi} \left(\int_0^1 F_2(x, Q^2) dx \right) R_{TERM}(x, Q^2) = B - A R_{TERM}$$

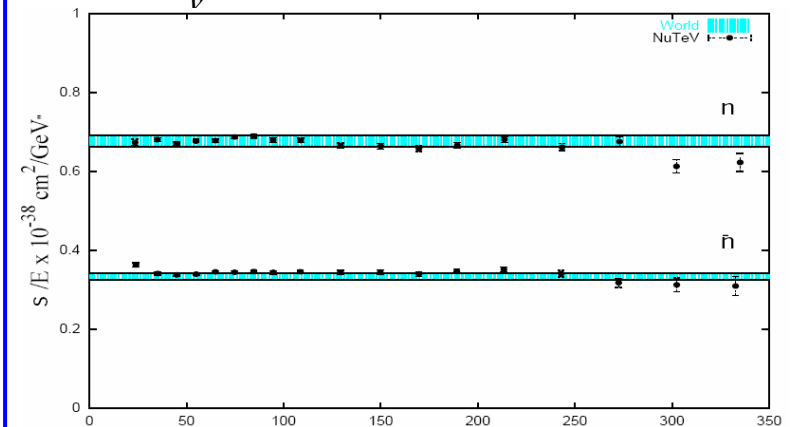
$$\frac{dN_{\nu, \bar{\nu}}}{d\nu} \xrightarrow{\nu \rightarrow 0} \Phi(E_\nu) A_{\nu, \bar{\nu}}$$

$$\Phi(E_\nu) = \int_0^{\nu_0/E_\nu} \left(\frac{dN(E_\nu)}{dy} / \left(1 + \frac{B}{A} y - \frac{C}{A} \frac{y^2}{2} \right) \right) dy$$

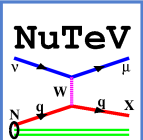
Normalizing the relative flux

- using world average neutrino cross section for $30\text{GeV} < E_\nu < 200\text{GeV}$

$$\frac{\sigma}{E_\nu} = 0.677 \pm 0.014 \times 10^{-38} \frac{\text{cm}^2}{\text{GeV}}$$



- the neutrino cross section $\frac{\sigma}{E_\nu}$ is flat as a function of E_ν within less than 2%
- the relative antineutrino cross section agrees with the world average



Systematic Uncertainties

There are **23 systematic uncertainties** which are considered:

- E_μ and E_{HAD} energy scales affect both the flux and the diff. cross section extraction

$$\frac{\Delta E_\mu}{E_\mu} = 0.7\%$$

$$\frac{\Delta E_{HAD}}{E_{HAD}} = 0.43\%$$

- m_c and B/A are important for the relative flux extraction

$$m_c = 1.4 \pm 0.13$$

$$B/A^v = -0.27 \pm 0.03, \quad B/A^{\bar{v}} = -1.64 \pm 0.03$$

- neutrino world average cross section has **2.1%** uncertainty – used to normalize the flux. (included as overall normalization)
- the uncertainty in the E_μ and E_{HAD} energy smearing models.
- 16** model fit parameters.

The χ^2 including all systematic uncertainties is given by

$$\chi^2 = \sum_{\alpha\beta} (D_\alpha - Theory) M_{\alpha\beta}^{-1} (D_\beta - Theory)$$

$$M_{\alpha\beta} = \sum_{ij} \rho_{ij} \partial_{\alpha i} \partial_{\beta j} \sigma_{\alpha i} \sigma_{\beta j}$$

$M_{\alpha\beta}$ - point to point covariance matrix

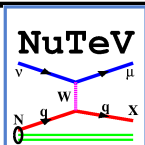
ρ_{ij} - 22x22 correlation matrix of all uncertainties

$\sigma_{\alpha i}$ - the size of systematic uncertainty i at data point α

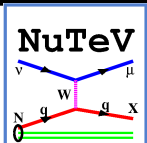
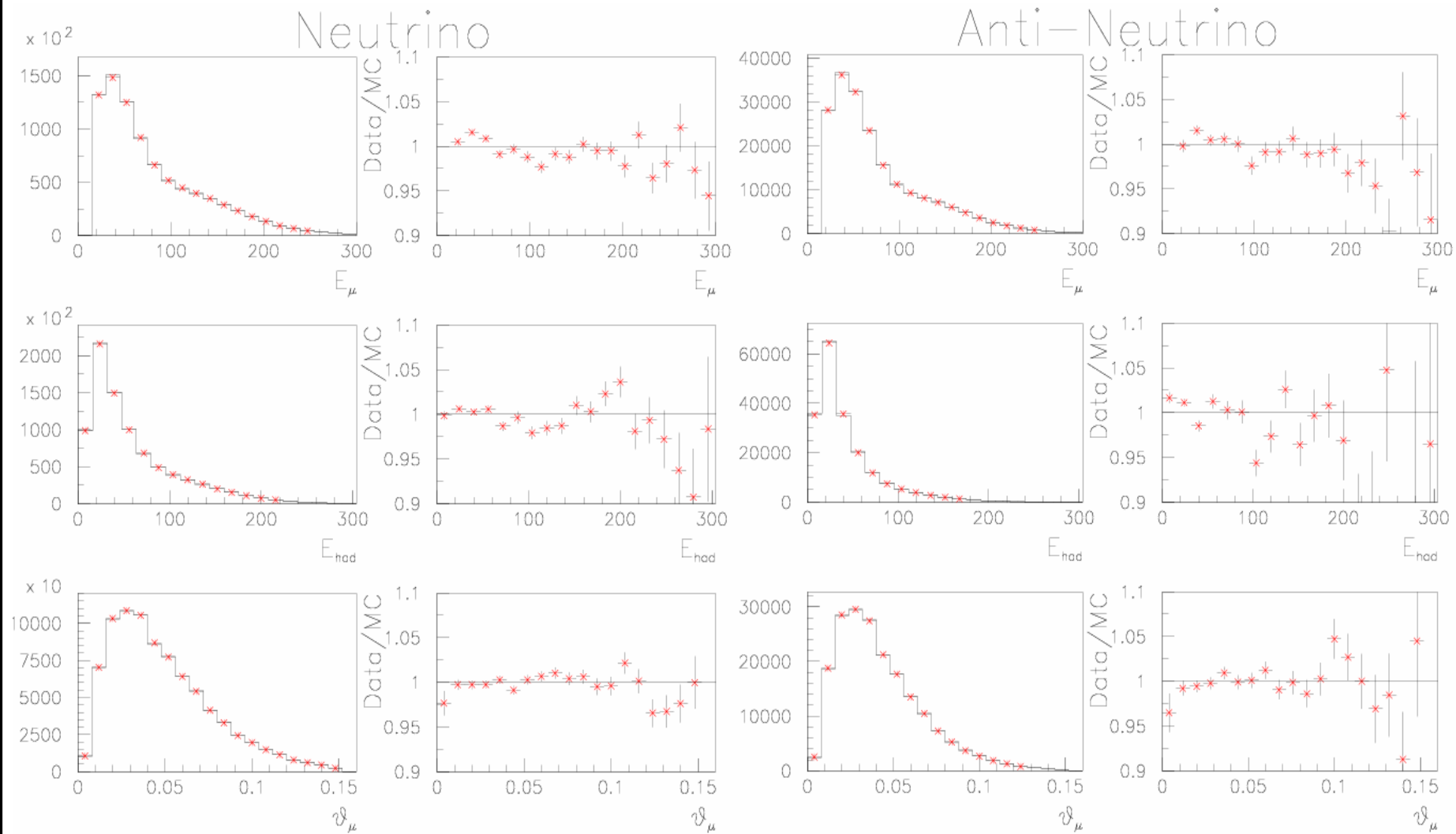
Each derivative is calculated by varying the corresponding systematic S_i by a small ε_i using:

$$\partial_{i\alpha} = \frac{\frac{d^2\sigma}{dx dy}(x, y, E)(S_i + \varepsilon_i) - \frac{d^2\sigma}{dx dy}(x, y, E)(S_i - \varepsilon_i)}{2\varepsilon_i}$$

- statistical uncertainty is added in quadrature to the diagonal element of the matrix



Monte Carlo Modeling of Data



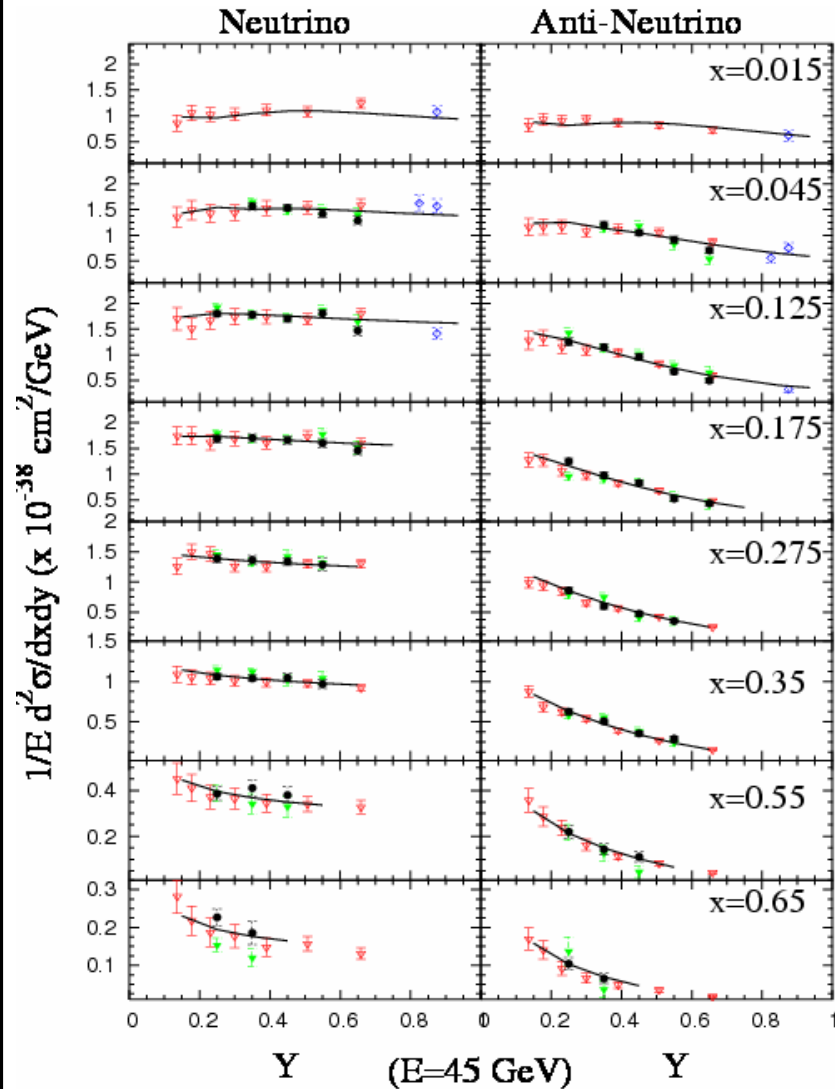
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NuTeV CC Cross Section



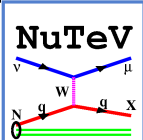
	E_μ scale	E_{HAD} scale	E_ν range (GeV)
CDHSW	2%	2.5%	20-200
CCFR	1%	1%	30-350
NuTeV	0.43%	0.7%	30-300

- all major systematics are included
- new high- y cross section at low x and low E_ν (different systematics – not included)
- NuTeV has comparable statistics to the other ν -Fe experiments.
- Reduction in the largest systematic uncertainties : E_μ and E_{HAD} scales

Other ν -Fe data shown on the plot:

CDHSW (Z. Phys. C49 187, 1991)

U. K. Yang CCFR Ph.D. Thesis



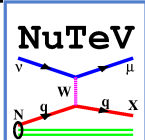
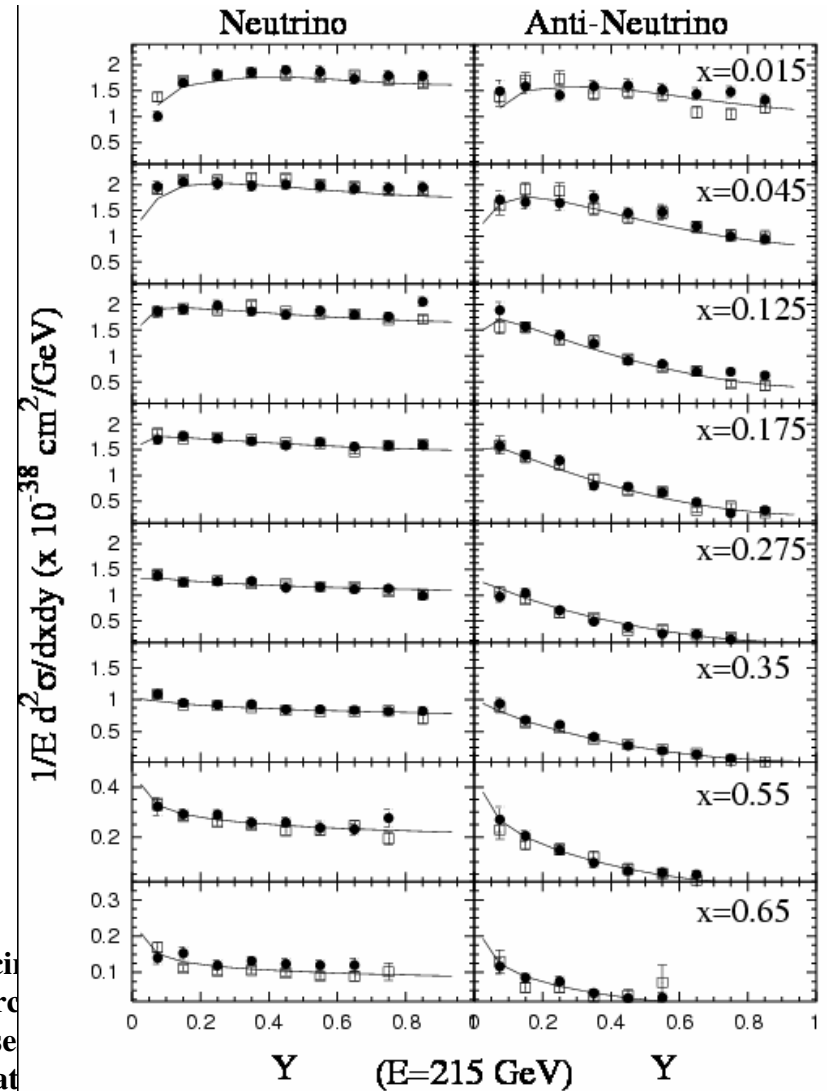
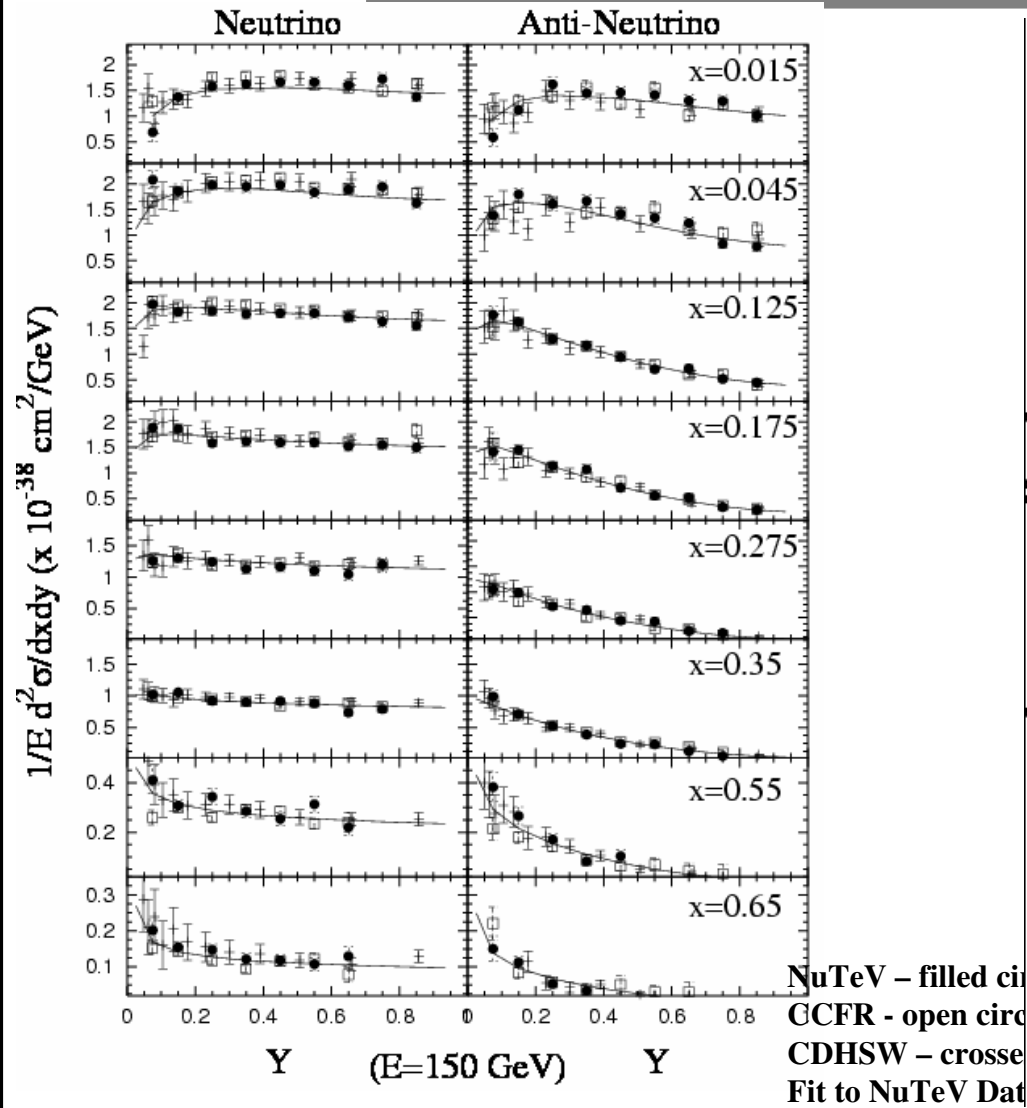
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Preliminary NuTeV's CC Cross Section



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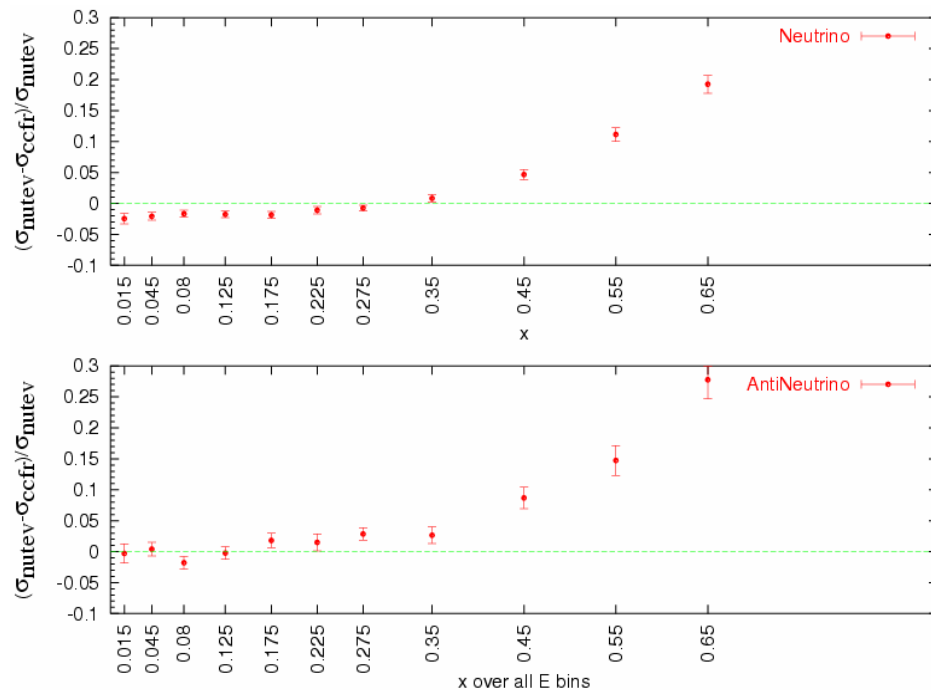
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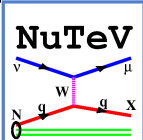
NuTeV CC Cross Section Result

- it agrees well with other ν -Fe experiments for $x < 0.4$
- for $x > 0.4$ it agrees in level, but not in shape with CDFSW
- for $x > 0.4$ NuTeV result is increasingly higher than CCFR
 - $x=0.45$ - 5%
 - $x=0.55$ - 10%
 - $x=0.65$ - 20%



Investigating the discrepancy with CCFR

- similar detectors and techniques
- NuTeV has improved calibration
 - 3% effect at $x=0.65$
- Different model fit parameters
 - 3% effect at $x=0.65$
- Other differences are:
 - NuTeV had separate neutrino and antineutrino beams
 - NuTeV - always focusing for the “right-sign” muon
 - CCFR – simultaneous neutrino and antineutrino runs
 - CCFR toroid polarity was half of the time focusing for μ^+ and half of the time μ^-



F₂ and xF₃ Measurement

F₂

$$\left[\frac{d^2\sigma^v}{dx dy} + \frac{d^2\sigma^{\bar{v}}}{dx dy} \right] \frac{\pi}{G_F^2 ME} =$$

$$= 2\bar{F}_2 \left(1 - y - \frac{Mxy}{2E} + \frac{y^2}{2} \frac{1+4M^2x^2/Q^2}{1+R} \right) + y \left(1 - \frac{y}{2} \right) \Delta x F_3$$

xF₃

$$\left[\frac{d^2\sigma^v}{dx dy} - \frac{d^2\sigma^{\bar{v}}}{dx dy} \right] \frac{\pi}{G_F^2 ME} =$$

$$= \Delta F_2 \left(1 - y - \frac{Mxy}{2E} + \frac{y^2}{2} \frac{1+4M^2x^2/Q^2}{1+R} \right) + 2y \left(1 - \frac{y}{2} \right) x\bar{F}_3$$

- Perform 1-parameter fit for F₂
- ²R_w and model for ΔxF₃ are used
- ¹TR-VFS ΔxF₃ model:

¹(R. S. Thorne and R. G. Roberts, *Phys. Lett. B* 421, 303 (1998)).

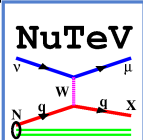
²(L.W. Whitlow et. al. *Phys. Lett. B* 250, 193 (1990)).

- Perform 1-parameter fit for xF₃
- ΔF₂ is very small and is neglected

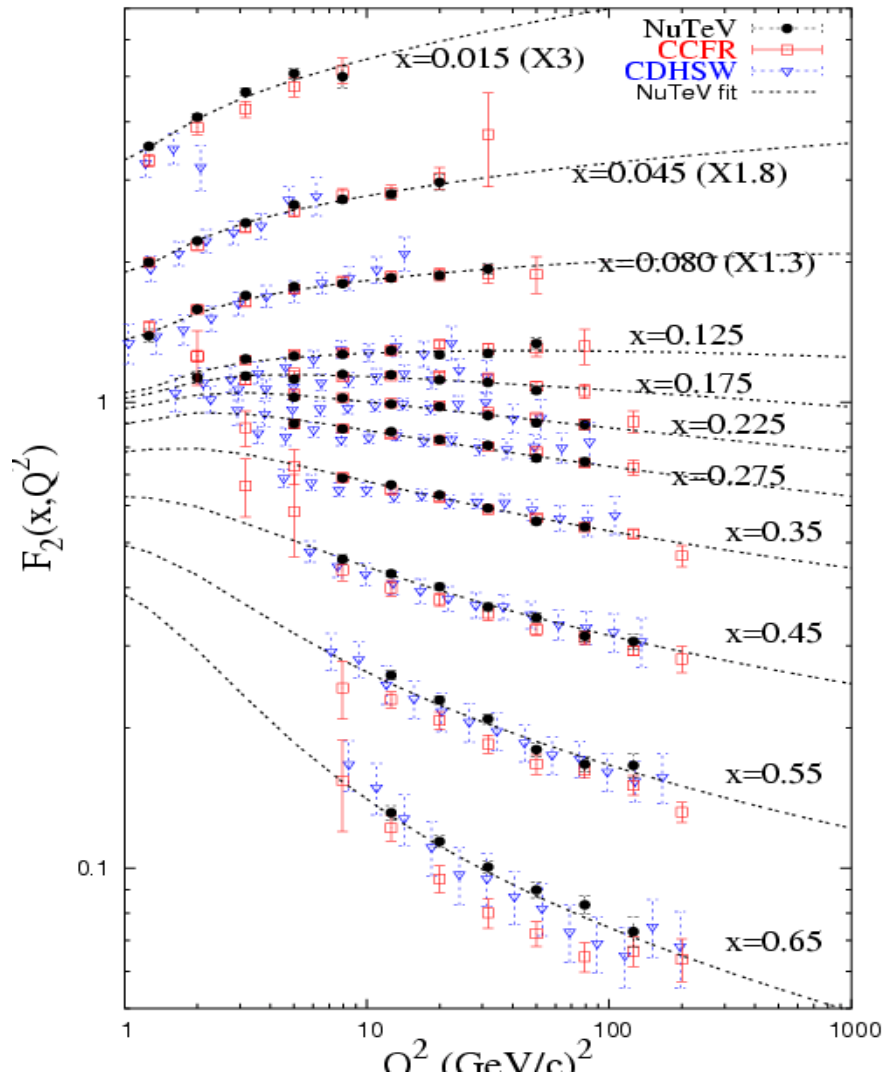
- Radiative corrections applied

(D. Y. Bardin and Dokuchaeva, *JINR-E2-86-260* (1986))

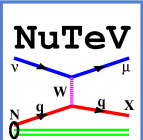
- Isoscalar correction applied
- Bin centering correction for Q²



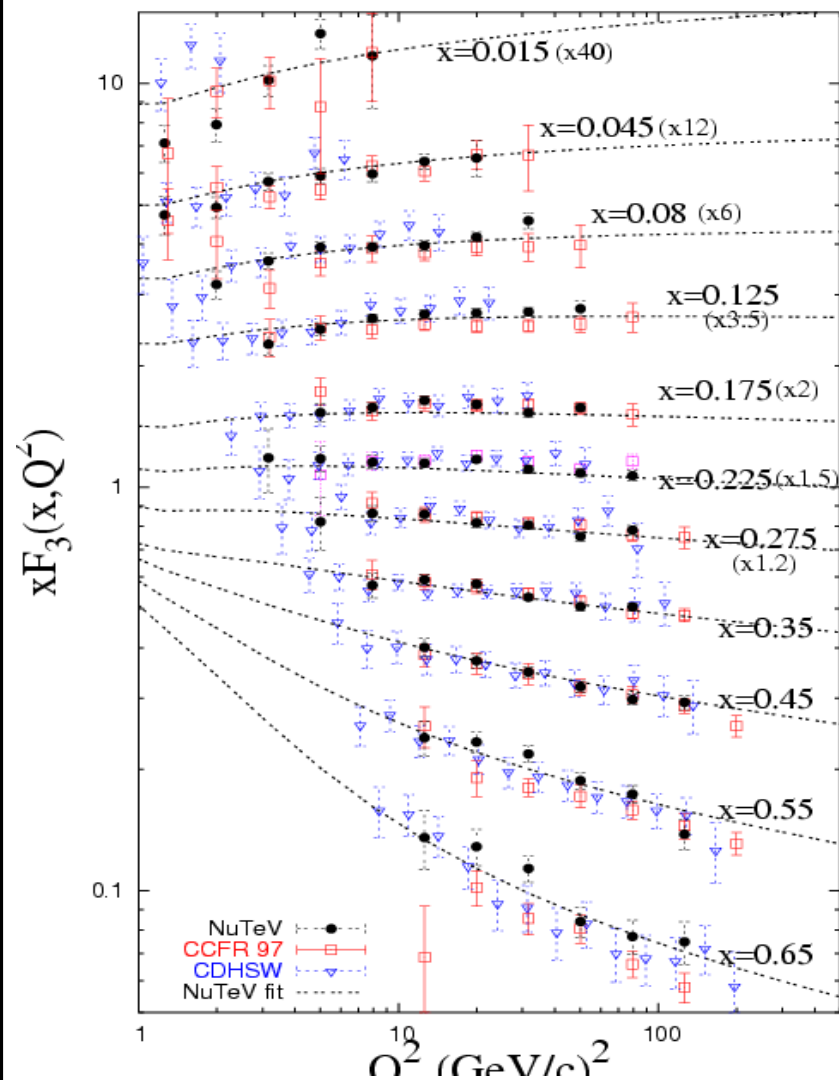
F_2 Measurement



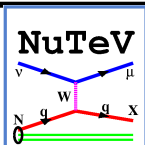
- Isoscalar ν -Fe F_2
- NuTeV F_2 is compared with **CCFR** and **CDHSW** results
 - the line is a fit to NuTeV data
- All systematic uncertainties are included
- All data sets agree for $x < 0.4$.
- At $x > 0.4$ NuTeV agrees in level with **CDHSW**, but not in shape
- At $x > 0.4$ NuTeV is systematically above **CCFR**



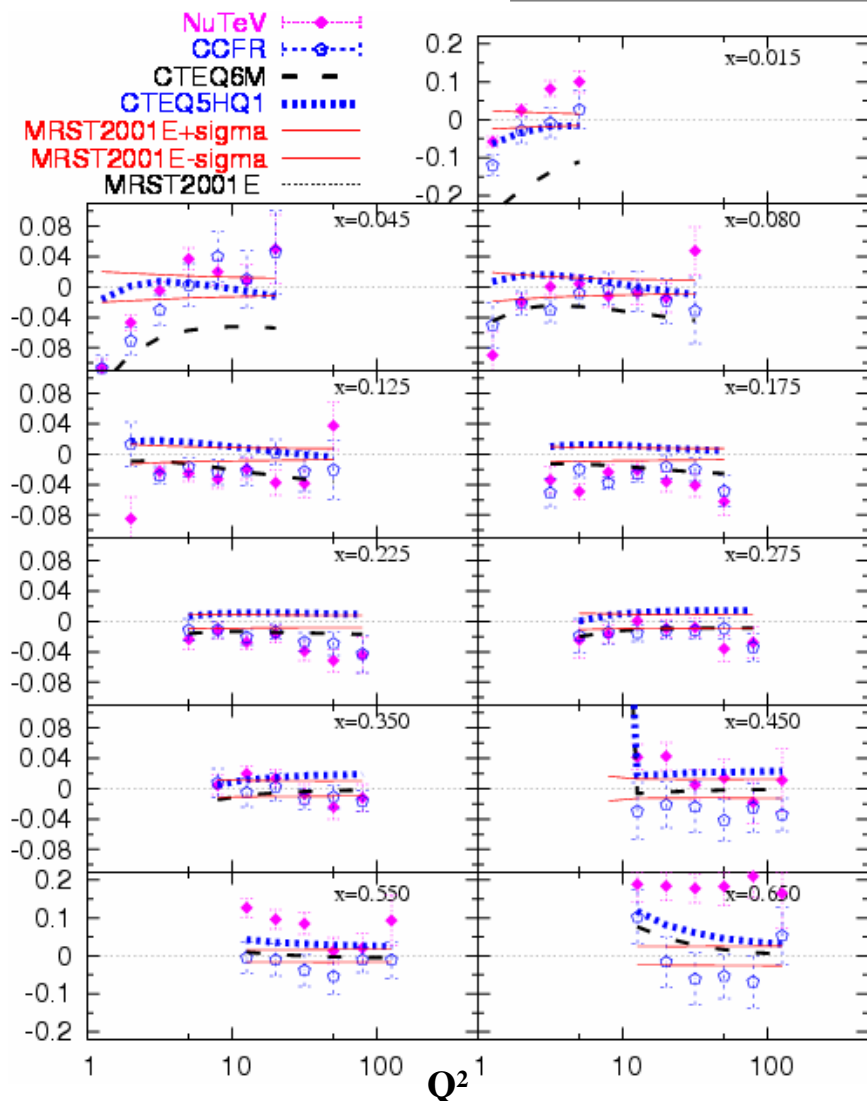
xF_3 Measurement



- Isoscalar ν -Fe xF_3
- NuTeV xF_3 is compared with **CCFR** and **CDHSW** results
 - the line is a fit to NuTeV data
- All systematic uncertainties are included
- All data sets agree for $x < 0.4$.
- At $x > 0.4$ NuTeV agrees in level with **CDHSW**, but not in shape
- At $x > 0.4$ NuTeV is systematically above **CCFR**



Comparison with Theory



- base line is TRVFS(MRST2001E)

- NuTeV and CCFR F_2 are compared to

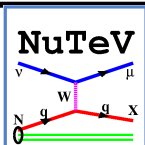
$$\text{TRVFS(MRST99)} \quad \frac{F_2^{\text{NuTeV}} - F_2^{\text{TRVFS}}}{F_2^{\text{TRVFS}}}$$

- Other theoretical models shown are

- ACOT(CTEQ6M)
- ACOT(CTEQ5HQ1)
- **TRVFS (MRST2001E)**

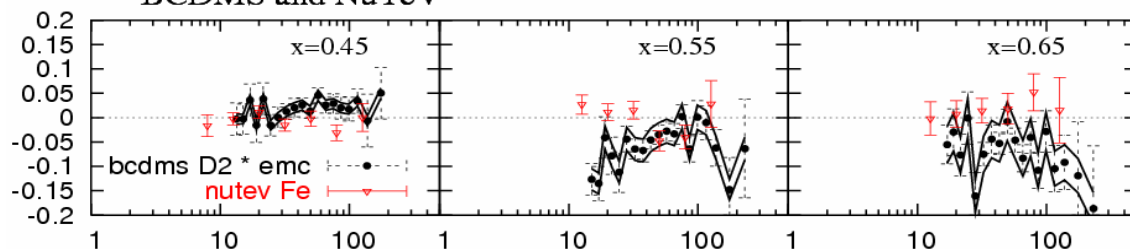
- theory curves are corrected for:

- target mass
- (*H. Georgi and H. D. Politzer, Phys. Rev. D14, 1829*)
- nuclear effects – parameterization from charge lepton data, assumed to be the same for neutrino scattering (no Q^2 dependence added)
- nuclear effects parameterization is dominated by SLAC (lower Q^2 in this region) data at high-x

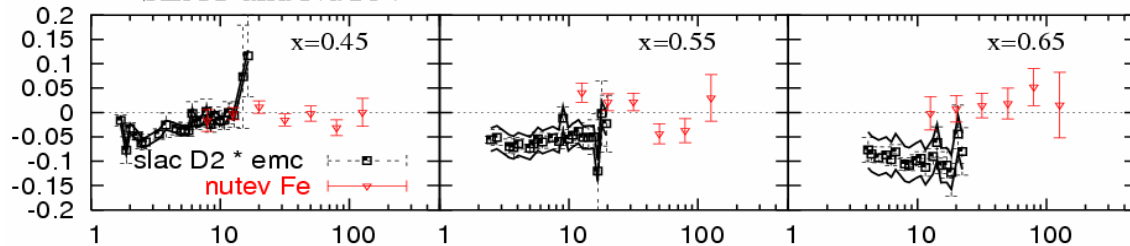


Comparison with Charge Lepton Data for $x > 0.4$

BCDMS and NuTeV



SLAC and NuTeV

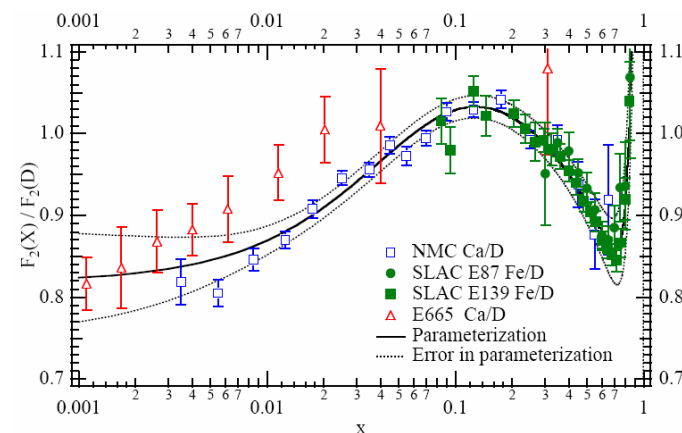


- base line is NuTeV model fit
- data points are
$$\frac{F_2^{DATA} - F_2^{BG}}{F_2^{BG}}$$
- **NuTeV** is higher than BCDMS(D_2) (top 3 plots)
 - 7% at $x=0.55$, and $x=0.65$
- **NuTeV** is higher than SLAC(D_2) (bottom 3 plots)
 - 5% at $x=0.55$, and 7% at $x=0.65$

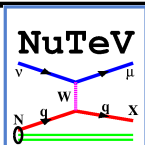
NuTeV data favors smaller nuclear correction.

- charge lepton data is corrected for:

- $\frac{F_2^V}{F_2^I}$ using CTEQ4D
- heavy target $\frac{F_2^N}{F_2^D}$



- **the nuclear correction is dominated by SLAC data, which is at lower Q^2 from NuTeV in this region**

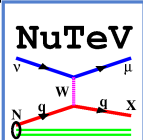


Conclusions

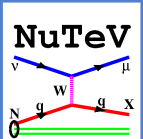
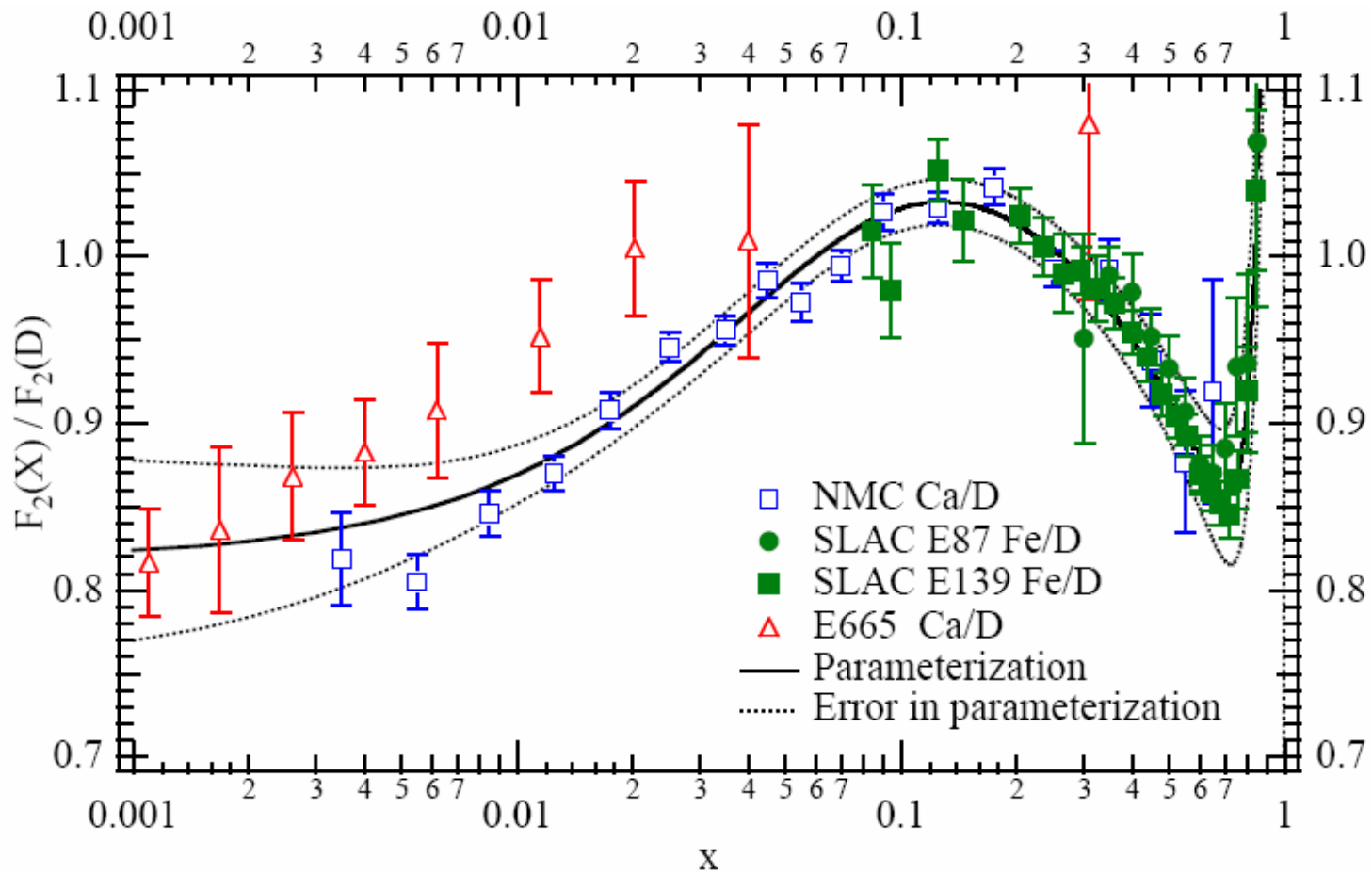
- NuTeV has extracted high precision $\nu - Fe$ and $\bar{\nu} - Fe$ differential cross section
- The **Sign Selected Quadrupole Train** is unique feature for NuTeV:
 - NuTeV measured differential cross section for inaccessible before high-y kinematic region.
- NuTeV measurement is in a good agreement with previous $\nu - Fe$ experiments (CCFR, CDHSW) for $x < 0.4$. For $x > 0.4$ NuTeV is systematically higher than CCFR:
5% at $x=0.45$, 10% at $x=0.55$, 20% at $x=0.65$
- **comparison with charge lepton data and theory shows that NuTeV favors smaller nuclear correction at high-x.**

Future:

- final checks
- QCD fits



Nuclear Correction



Buras-Gaemers Parameterization

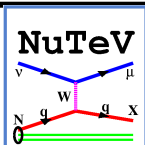
- **Valence quark distributions parameterization:**

$$x^{E_1} \times (1-x)^{E_2} + AV_2 \times x^{E_3} \times (1-x)^{E_4} + AV_3 \times x^{E_5} \times (1-x)^{E_6}$$

- **Sea quark distributions parameterization:**

$$AS_1 \times (1-x)^{ES_1} + AS_2 \times (1-x)^{ES_2}$$

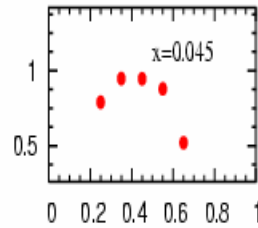
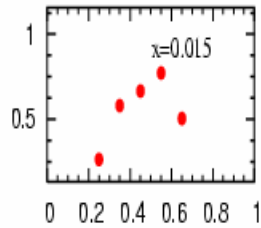
- **The exponents E_i and ES_i and the normalization coefficients AV_i and AS_i are fitted to NuTeV differential cross section data.**
- **13 parameters to fit. Fit includes standard QCD evolution, assumes $R=R_{\text{world}}$, charm mass, W-mass and sea constraint from dimuon data.**



Acceptance Correction

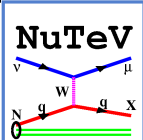
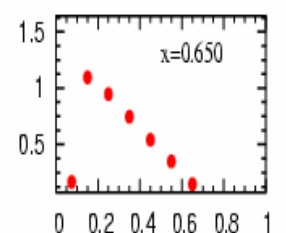
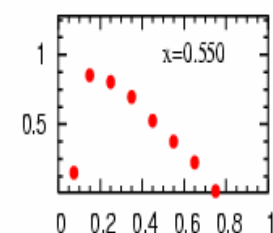
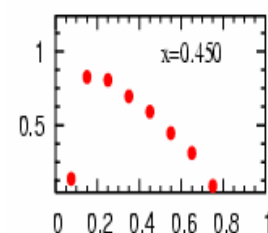
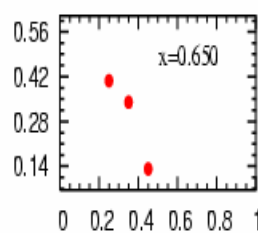
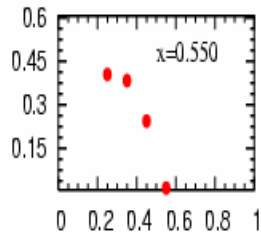
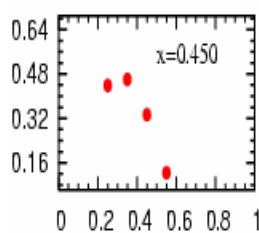
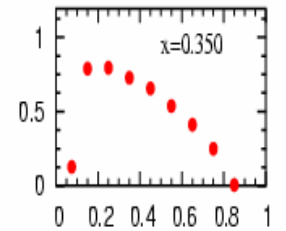
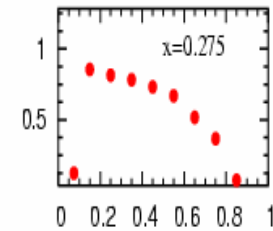
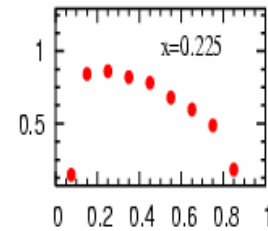
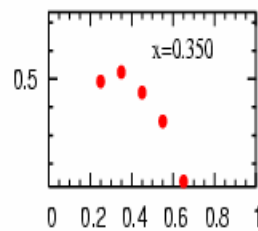
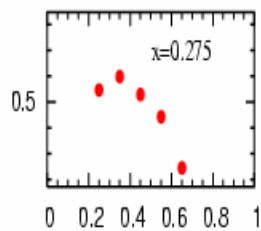
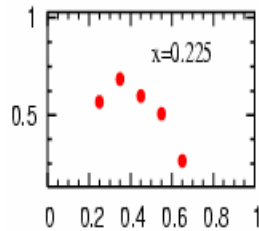
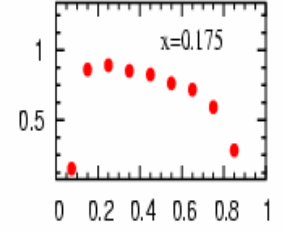
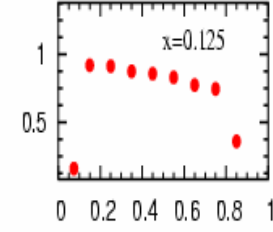
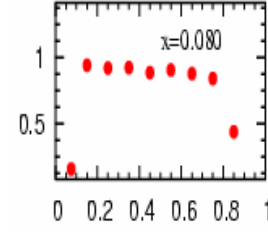
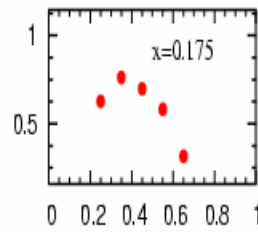
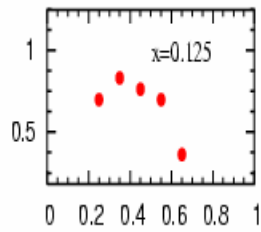
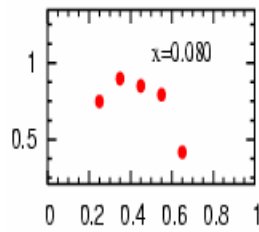
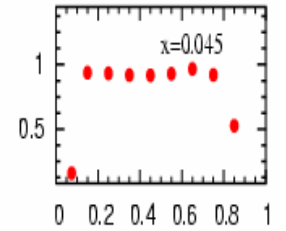
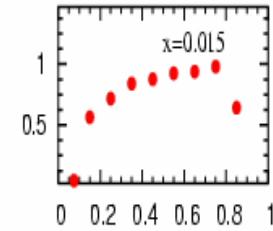
Energy = 45 GeV

Acceptance



Energy = 110 GeV

Acceptance



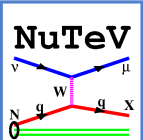
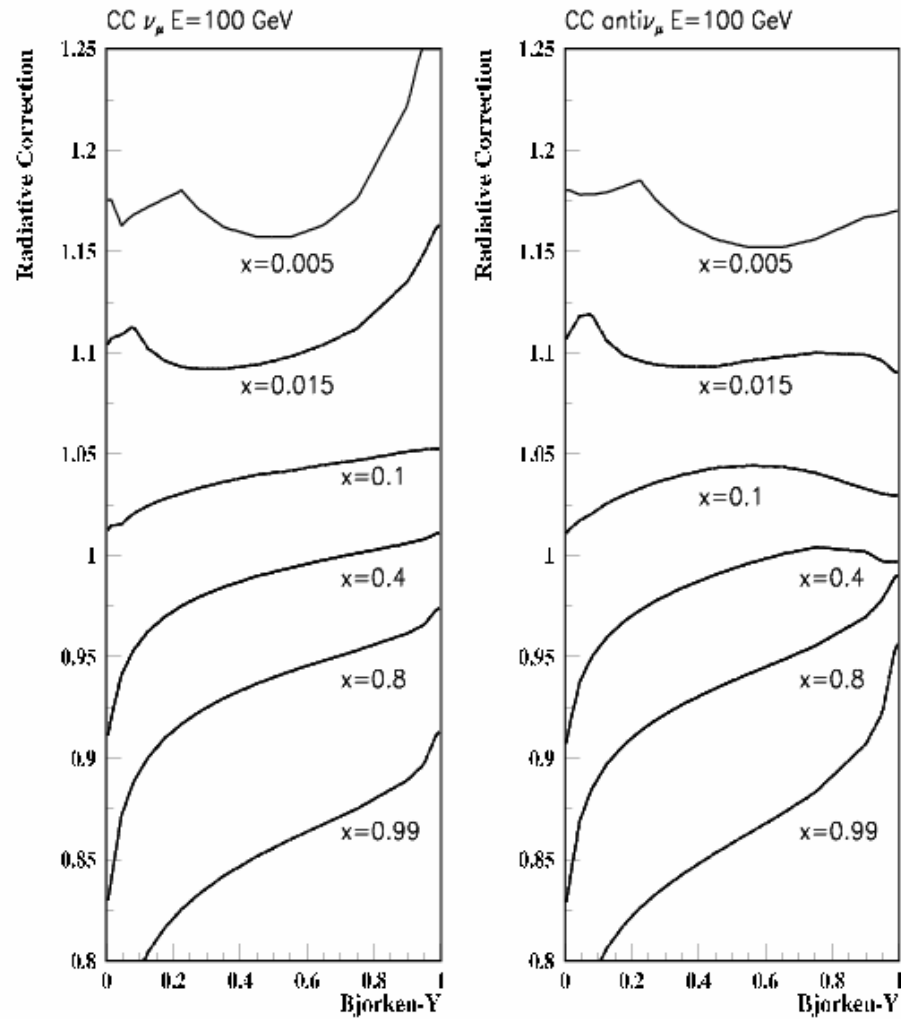
Martin Tzanov

DPF2004

August, 2004

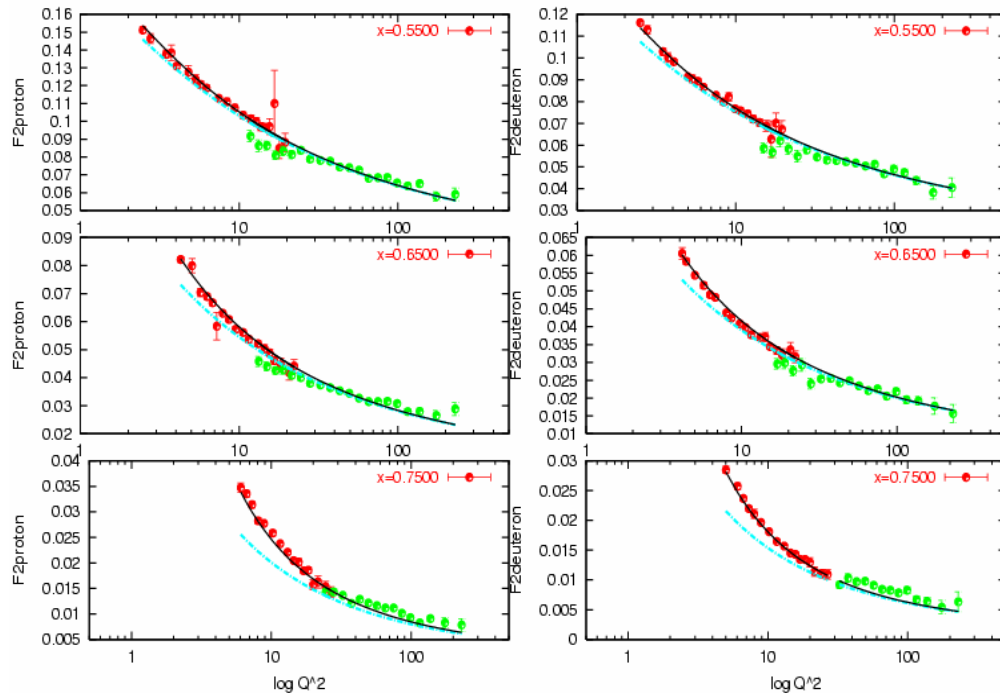


Radiative Correction



Higher Twist

- fit to ep, ed data (SLAC, NMC,BCDMS) to account for Target Mass and Higher-Twist effects in parton level cross section model
- important at high x and low Q²



$$x' = x \frac{Q^2 + B}{Q^2 + Ax}$$

$$F_2 \rightarrow \left(\frac{Q^2}{Q^2 + C} \right) F_2(x', Q^2)$$

A	0.57
B	0.22
C	0.06
χ^2/dof	792/312

