

# Search for $CP$ Violation in Hyperon Decays with the *HyperCP* Spectrometer at Fermilab

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## Why Search for $CP$ Violation in Hyperon Decays?

- After 40 years of intense experimental effort — and many beautiful experiments — we still know little about  $CP$  violation: the origin of  $CP$  violation remains unknown.
- Although  $CP$  is expected to be ubiquitous in weak interactions — albeit often vanishingly small — the experimental evidence is still meager.
- Although  $CP$  violation is accommodated quite nicely in the standard model, there is little hard evidence that it is the sole province of the standard model.
- Many beyond-the-standard-model theories can produce large new sources of  $CP$  violation, none of which have yet been seen.

**“We are willing to stake our reputation on the prediction that dedicated and comprehensive studies of  $CP$  violation will reveal the presence of New Physics.”**

*Bigi and Sanda, CP Violation*

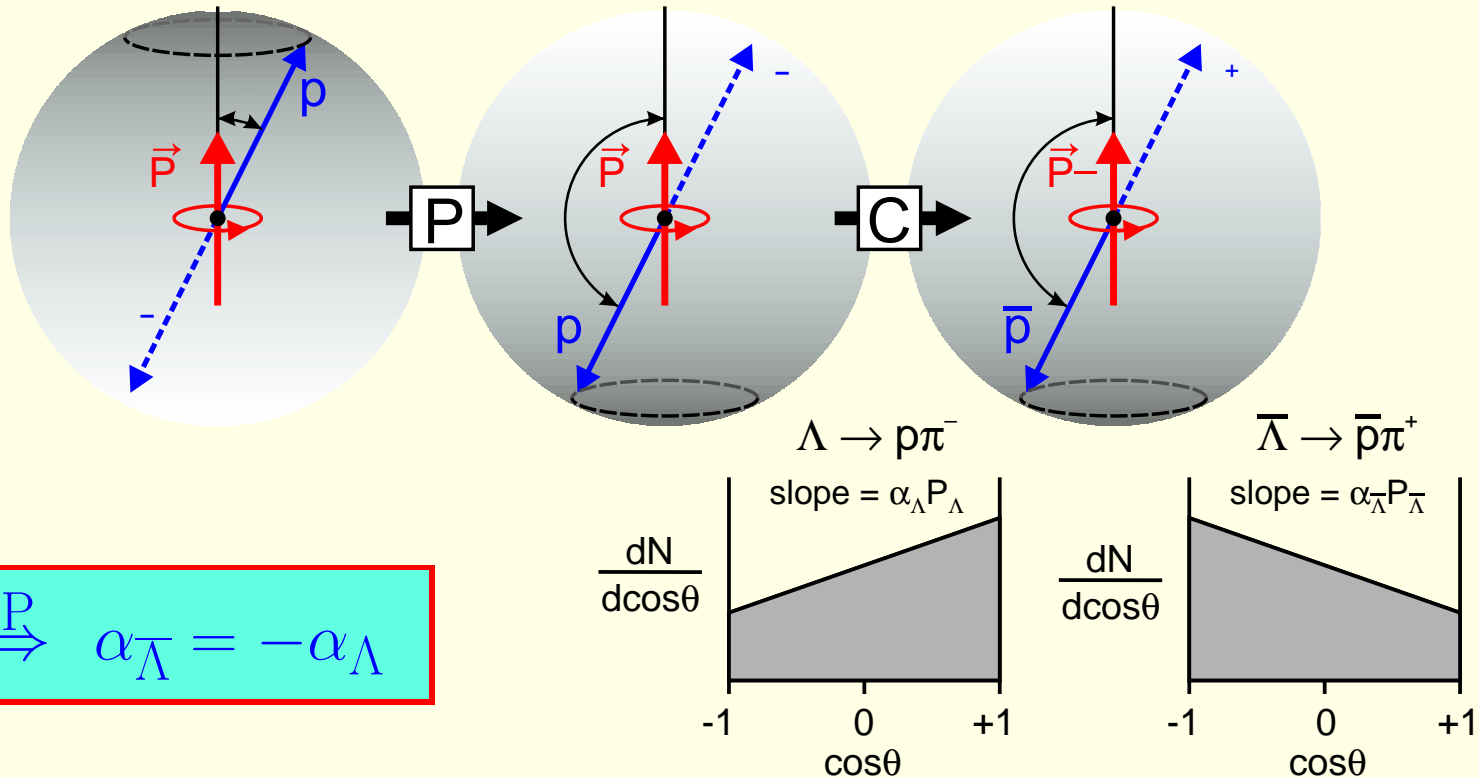
- Hyperons are sensitive to sources of  $CP$  violation that are not probed in other systems.
- These sources are experimentally accessible.
- The cost is small:
  - No new accelerators needed.
  - Apparatus is modest in scope and cost.

## How to Search for $CP$ Violation in $\Lambda$ Decays

Due to parity violation the proton likes to go in the direction of the  $\Lambda$  spin:

$$\Lambda \rightarrow p\pi^-: \quad \frac{dN(p)}{d\cos\theta} = \frac{N_0}{2}(1 + \alpha_\Lambda P_\Lambda \cos\theta) \quad \alpha = \frac{2\text{Re}(S^*P)}{|S|^2 + |P|^2} = 0.642$$

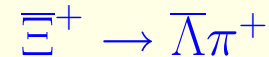
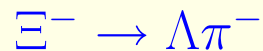
Under  $CP$  the antiproton likes to go in the direction **opposite** to the  $\bar{\Lambda}$  spin:



**Problem:** The  $\Lambda/\bar{\Lambda}$  polarizations have to be precisely known to extract  $\alpha_\Lambda/\bar{\alpha}_\Lambda$

## Producing Polarized $\Lambda/\bar{\Lambda}$ 's : unpolarized $\Xi$ Decays

In this technique, pioneered by *HyperCP*,  $\Lambda/\bar{\Lambda}$ 's of known polarization are produced from **unpolarized**  $\Xi^-/\bar{\Xi}^+$ 's:



If the  $\Xi$  is produced unpolarized — which can simply be done by targeting at 0 degrees — then the  $\Lambda$  is found in a helicity state, with a large polarization ( $\alpha_{\Xi} = -0.458$ ):

$$\vec{P}_{\Lambda} = \alpha_{\Xi}\hat{p}_{\Lambda}$$

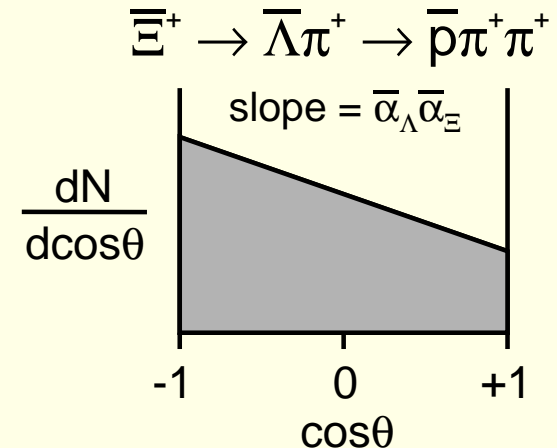
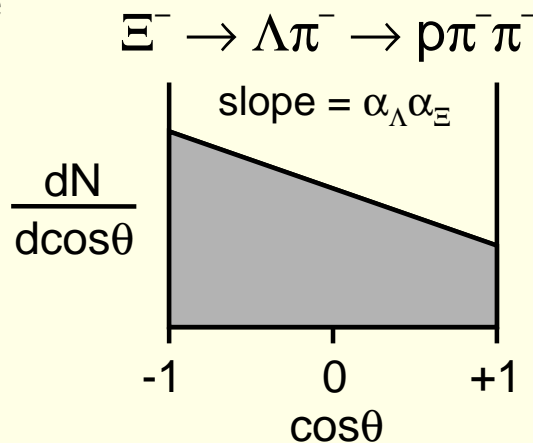
$$\vec{P}_{\bar{\Lambda}} = \bar{\alpha}_{\Xi}\hat{p}_{\bar{\Lambda}}$$

$$\frac{dN(p)}{d\cos\theta} = \frac{N_0}{2}(1 + \alpha_{\Lambda}\alpha_{\Xi}\cos\theta)$$

$$\frac{dN(\bar{p})}{d\cos\theta} = \frac{N_0}{2}(1 + \bar{\alpha}_{\Lambda}\bar{\alpha}_{\Xi}\cos\theta)$$

If *CP* is good, the slopes of the proton and antiproton  $\cos\theta$  distributions are identical, and:

$$\alpha_{\Xi}\alpha_{\Lambda} = \bar{\alpha}_{\Xi}\bar{\alpha}_{\Lambda}$$



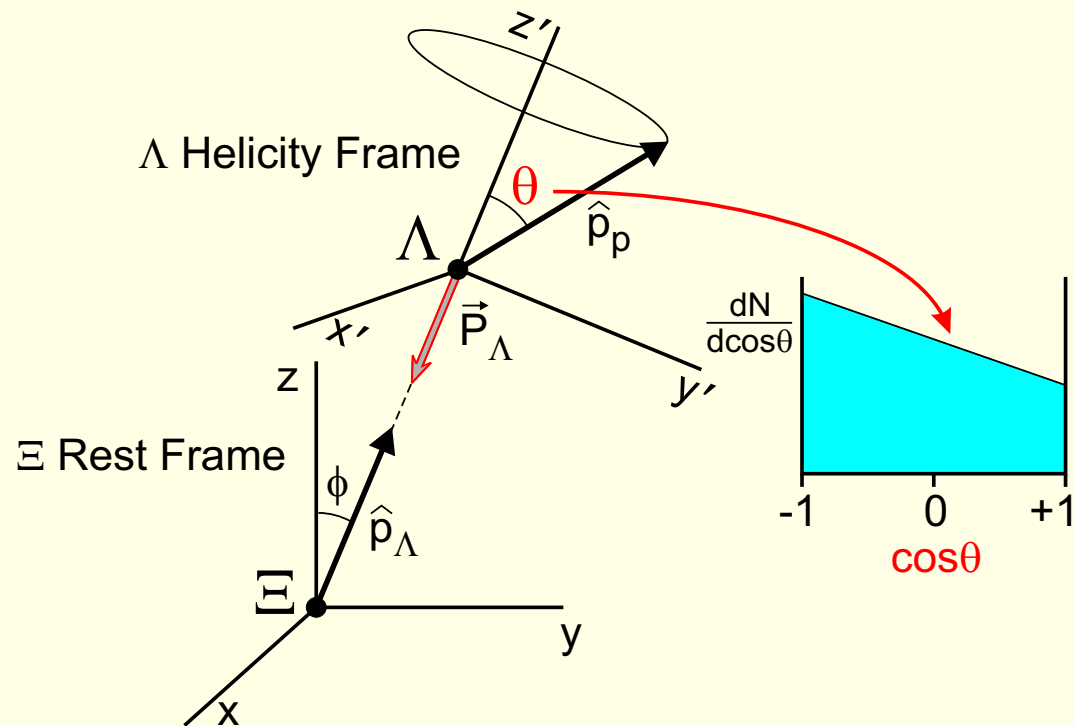
## HyperCP technique is sensitive to both $\Xi$ and $\Lambda$ $CP$ violation

$$\frac{\alpha_{\Xi}\alpha_{\Lambda} - \bar{\alpha}_{\Xi}\bar{\alpha}_{\Lambda}}{\alpha_{\Xi}\alpha_{\Lambda} + \bar{\alpha}_{\Xi}\bar{\alpha}_{\Lambda}} \approx A_{\Xi} + A_{\Lambda}$$

where:  $A_{\Xi} = \frac{\alpha_{\Xi} + \bar{\alpha}_{\Xi}}{\alpha_{\Xi} - \bar{\alpha}_{\Xi}}$  and  $A_{\Lambda} = \frac{\alpha_{\Lambda} + \bar{\alpha}_{\Lambda}}{\alpha_{\Lambda} - \bar{\alpha}_{\Lambda}}$

What HyperCP experimentally measures  $\Rightarrow$

Important: polar axis changes from event to event.



## Phenomenology of $CP$ Violation in $\Xi$ and $\Lambda$ Decay

- $CP$  violation in  $\Xi$  and  $\Lambda$  decays is manifestly **direct** with  $\Delta S = 1$ .
- Three ingredients are needed to get a non-zero asymmetry:
  1. At least two channels in the final state: the  $S$ - and  $P$ -wave amplitudes.
  2. The  $CP$  violating weak phases must be different in the two channels.
  3. There must be unequal final-state scattering phase shifts in the two channels.

$$\begin{aligned}
 A_\Lambda &= (\alpha_\Lambda + \alpha_{\bar{\Lambda}})/(\alpha_\Lambda - \alpha_{\bar{\Lambda}}) \cong -\tan(\delta_P - \delta_S) \sin(\phi_P - \phi_S), \\
 A_\Xi &= (\alpha_\Xi + \alpha_{\bar{\Xi}})/(\alpha_\Xi - \alpha_{\bar{\Xi}}) \cong -\tan(\underbrace{\delta_P - \delta_S}_{\text{strong phases}}) \sin(\underbrace{\phi_P - \phi_S}_{\text{weak phases}}).
 \end{aligned}$$

- Asymmetry greatly reduced by the small strong phase shifts.
  - The  $p\pi$  phase shifts have been measured to a precision of about one degree:

$$\Lambda \begin{cases} \delta_P = -1.1 \pm 1.0^\circ \\ \delta_S = 6.0 \pm 1.0^\circ \end{cases}$$

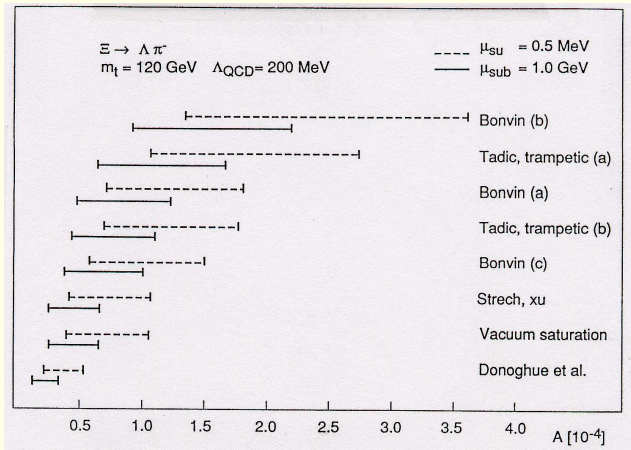
- The  $\Lambda\pi$  phase shifts can't be directly measured, theoretical predictions disagree:

$$\Xi^- \left\{ \begin{array}{l} \delta_P = -2.7^\circ \\ \delta_S = -18.7^\circ \end{array} \right\}_{1965} \quad \left. \begin{array}{l} = -1^\circ \\ = 0^\circ \end{array} \right\} \text{recent } \chi PT$$

**HyperCP has measured the  $\Lambda\pi$  phase shift:  $(4.6 \pm 2.3)^\circ$**

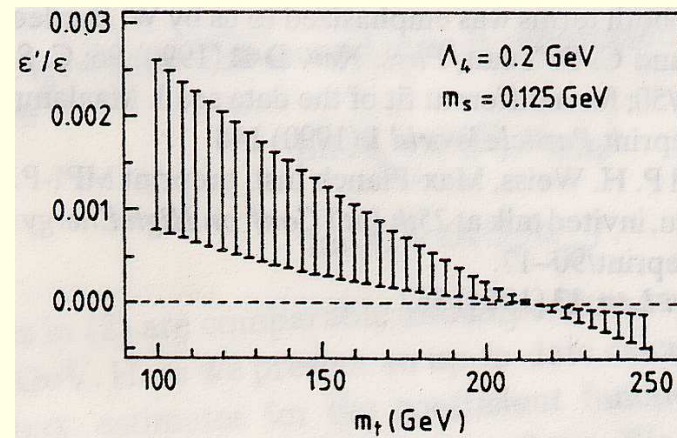
## Bad News: Standard Model Theory Predictions Small

- Much enthusiasm a decade ago as Standard Model predictions were relatively large.



Valencia (1991)

- At same time there was concern that accidental cancellation in the kaon system would lead to  $\epsilon'/\epsilon \approx 0$ .



Paschos (1991)

- Standard Model predictions have slowly fallen to:

$$-0.5 \times 10^{-4} < A_{\Xi\Lambda} < +0.5 \times 10^{-4}$$

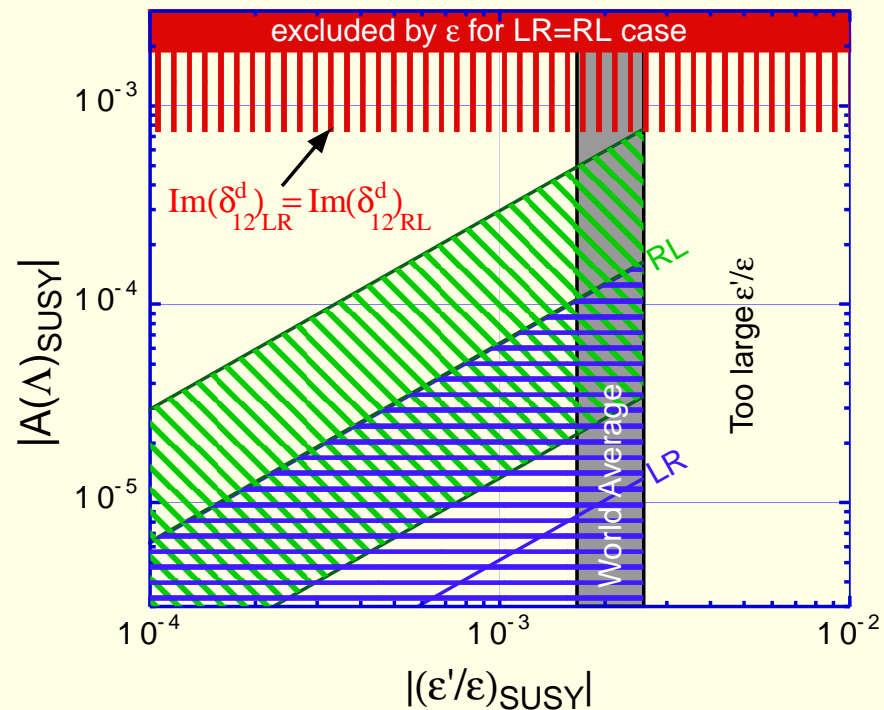
(Tandean & Valencia, 2003)

- The expected SM asymmetry is out of reach for any experiment, planned or otherwise.

**Important:** no unambiguous connection between:  $\delta_{CKM} \Leftrightarrow A_{\Xi}, A_{\Lambda}$

## Good News: Standard Model Theory Predictions Small

- Beyond-the-standard-model predictions larger, and not well constrained by kaon  $CP$  measurements: hyperon  $CP$  violation probes both parity conserving and parity violating amplitudes.
- Recent paper by Tandean (2004) shows that the upper bound on  $A_{\Xi\Lambda}$  from  $\epsilon'/\epsilon$  and  $\epsilon$  measurements is  $\sim 100 \times 10^{-4}$ .
- For example, some supersymmetric models that do not generate  $\epsilon'/\epsilon$  can lead to  $A_{\Lambda}$  of  $O(10^{-3})$ .
- Other BSM theories, such as Left-Right mixing models, (Chang, He, Pakvasa (1994)), also have enhanced asymmetries.



He et al., PRD 61 (2000) 071701(R).

Any  $CP$ -violation signal will almost certainly come from New Physics.

## What is the experimental situation?

- To date there are only upper limits on the asymmetries.
- $A_\Lambda$  has been measured to  $2 \times 10^{-2}$ :

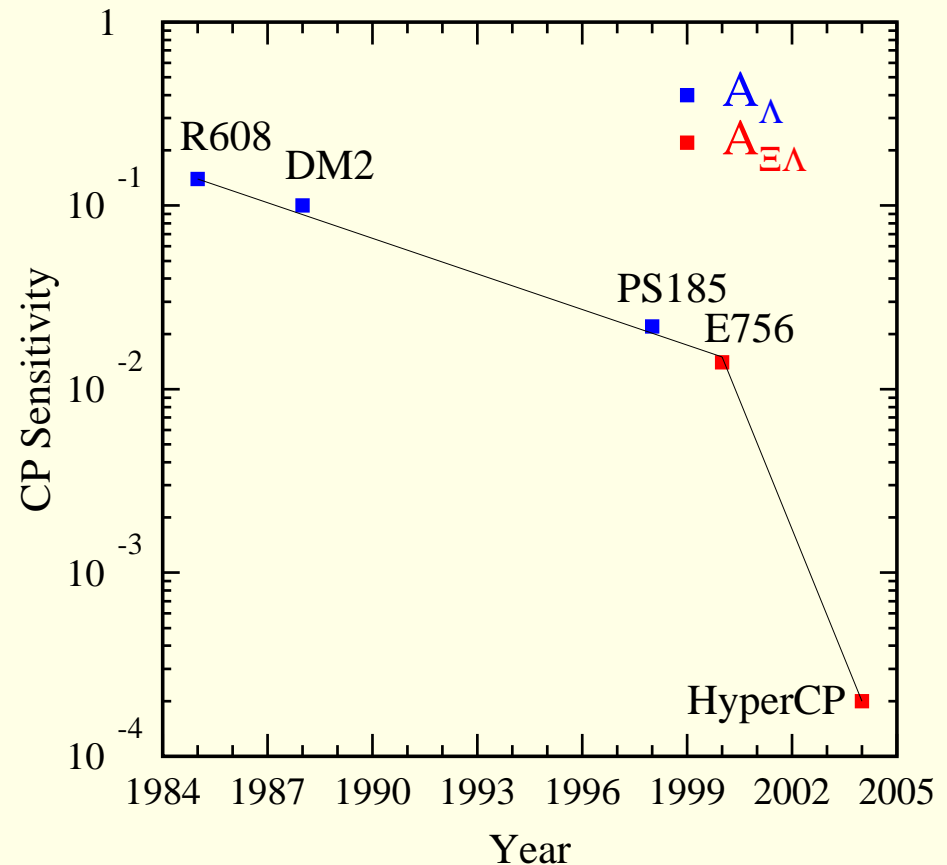
Exp	Mode	Method
R608	$A_\Lambda$	$p\bar{p} \rightarrow \Lambda X, p\bar{p} \rightarrow \bar{\Lambda} X$
DM2	$A_\Lambda$	$e^+e^- \rightarrow J/\psi \rightarrow \Lambda\bar{\Lambda}$
PS185	$A_\Lambda$	$p\bar{p} \rightarrow \Lambda\bar{\Lambda}$

- There is a recent measurement of  $A_{\Xi\Lambda}$ , based on the *HyperCP* technique:

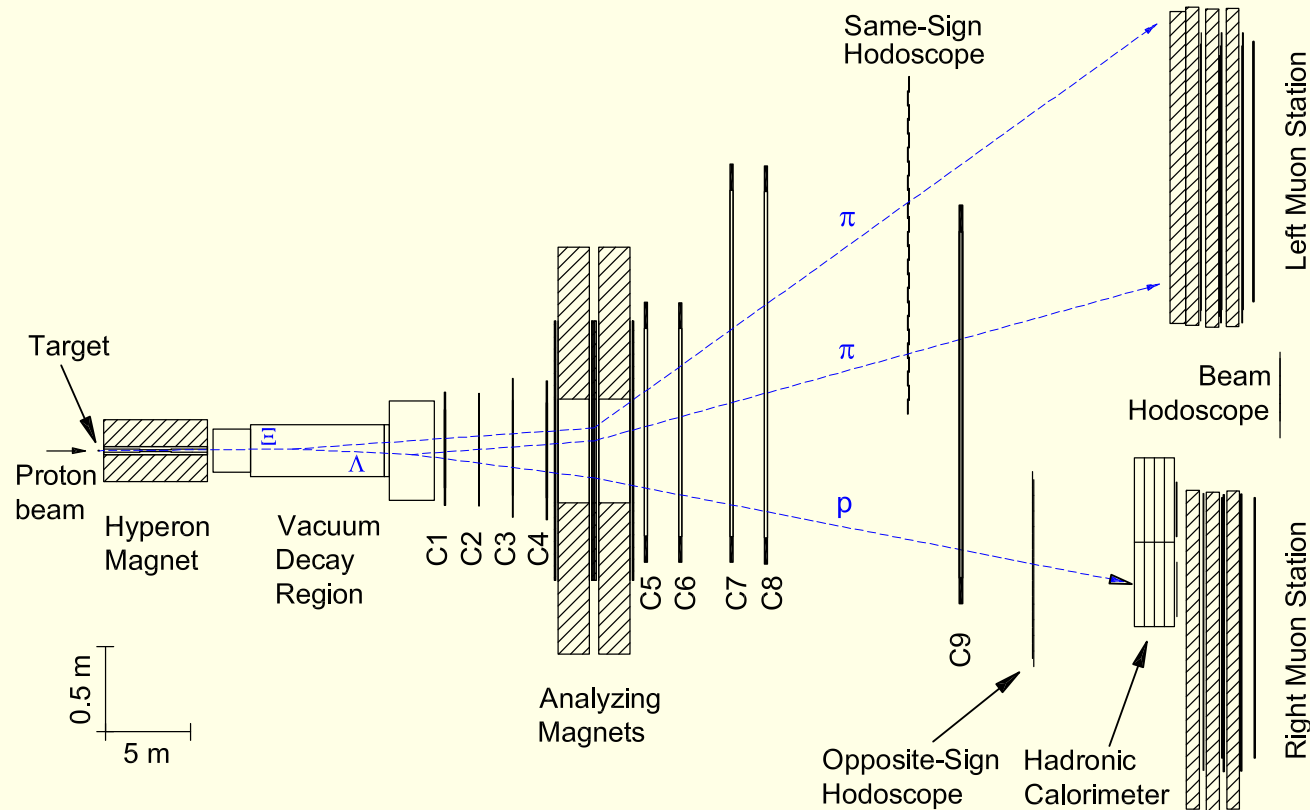
Exp	Mode	Method
E756	$A_{\Xi\Lambda}$	$pN \rightarrow \Xi^\pm X \rightarrow \Lambda\pi^\pm$

- This measurement of  $A_{\Xi\Lambda}$  can be used with measurements of  $A_\Lambda$  to infer a limit on  $A_\Xi$ .

- None of these measurements is in the regime of testing theory.
- HyperCP* is pushing two orders of magnitude beyond the best limit, to  $\sim 10^{-4}$ .**



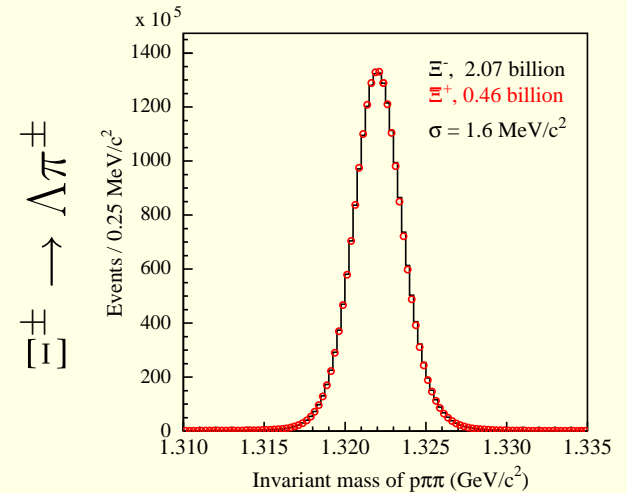
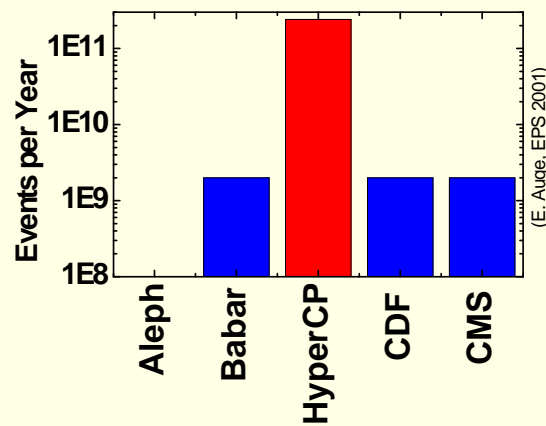
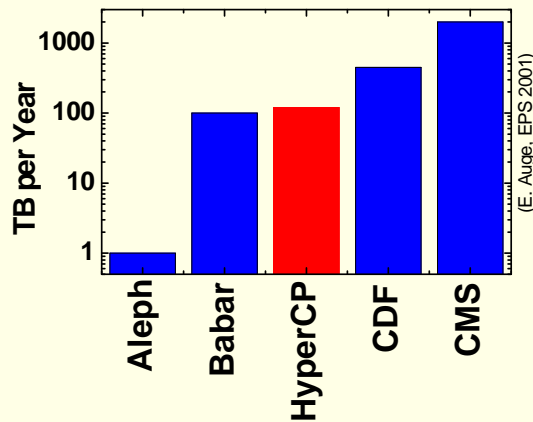
## The HyperCP Spectrometer



- Alternate + and - running.
- 800 GeV/c incident proton beam.
- 10–15 MHz, 167 GeV/c charged beam.
- High-rate, narrow-pitch wire chambers.
- Muon system for rare/forbidden hyperon and kaon decays.
- Very high-rate DAQ:
  - 50-80 KHz evts/spill-s to tape.
  - 27 MB/s on 27 Exabyte 8705 tape drives.
- Simple, low-bias trigger using hodoscopes and calorimeter.
  - SS( $\geq 1$  hit) · OS( $\geq 1$  hit) · Cal( $\geq 40$  G eV)

# HyperCP Yields

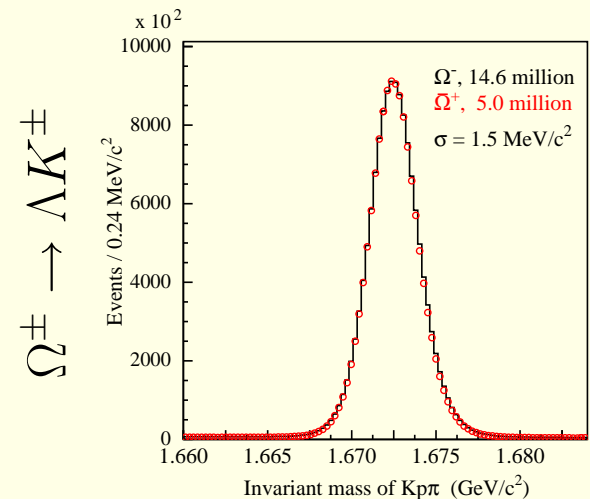
- In 12 months of data taking *HyperCP* recorded one the largest data samples ever by a particle physics experiment: 231 billion events, 29,401 tapes, and 119.5 TB data.



Entire WWW on 9/11/01 was **5 TB!**

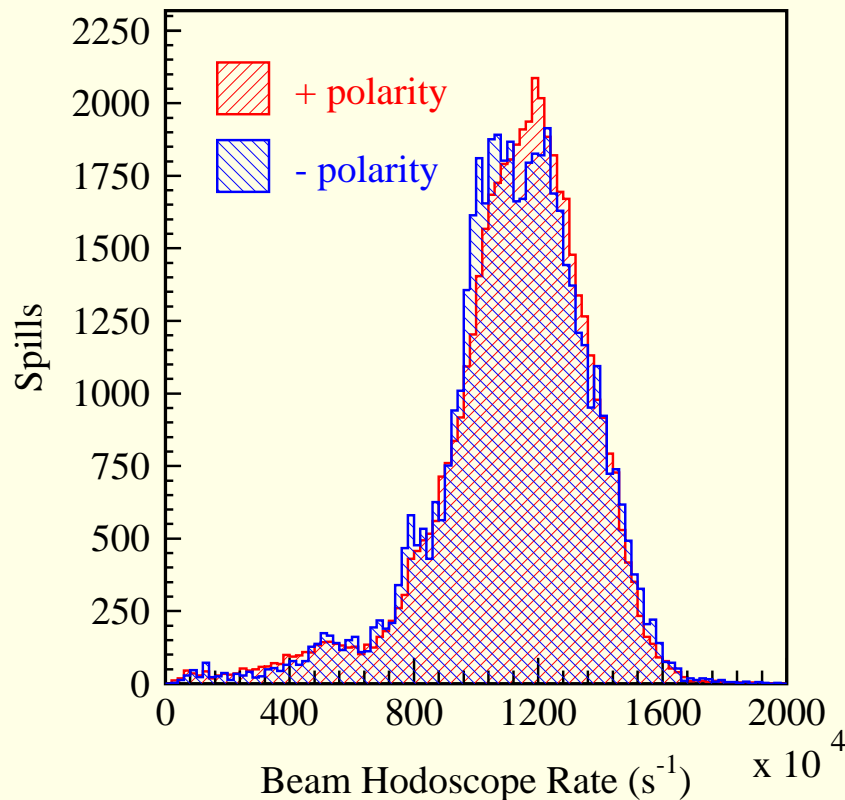
## Reconstructed Events

Type	Channeled beam polarity		Total
	+	-	
$\Xi \rightarrow \Lambda\pi$	$458 \times 10^6$	$2032 \times 10^6$	$2490 \times 10^6$
$K \rightarrow \pi\pi\pi$	$391 \times 10^6$	$164 \times 10^6$	$555 \times 10^6$
$\Omega \rightarrow \Lambda K$	$4.9 \times 10^6$	$14.1 \times 10^6$	$19.0 \times 10^6$

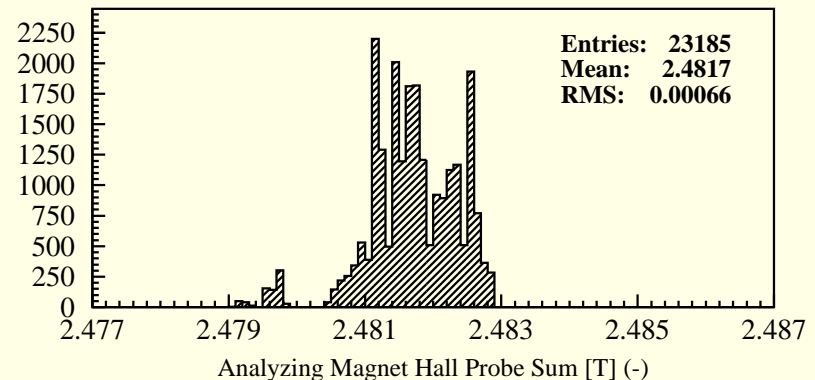
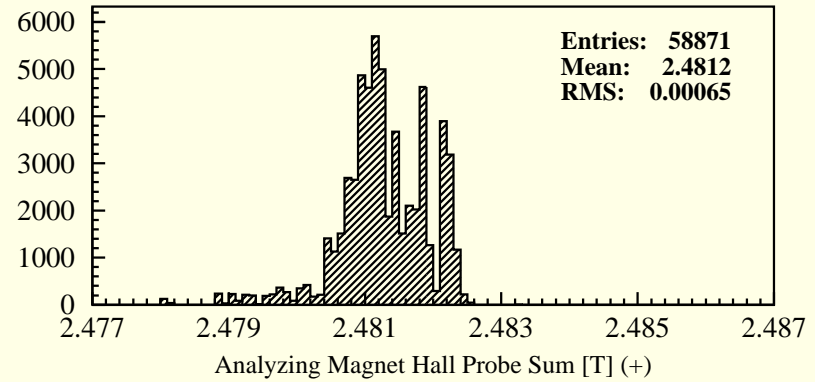


## Care Taken to Minimize Differences in + and - Running

- Targets changed to equalize secondary-beam rates.
  - + polarity: 2.0 cm Cu
  - polarity: 6.0 cm Cu
- When flipping polarity, field magnitude kept within  $\sim 2 \times 10^{-4}$ .
- This corresponds to a  $\sim 0.3$  mm deflection difference at 10 m for the lowest momentum ( $\sim 10$  GeV/c pions).



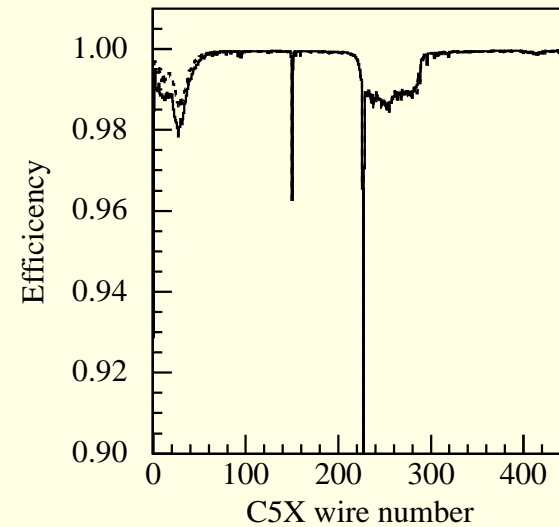
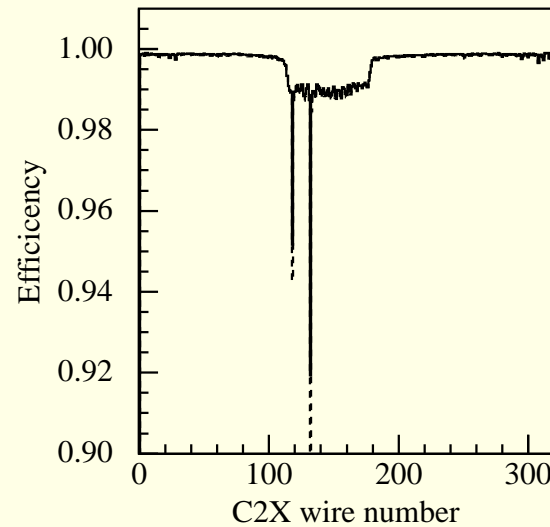
- About a 1% difference in rates.



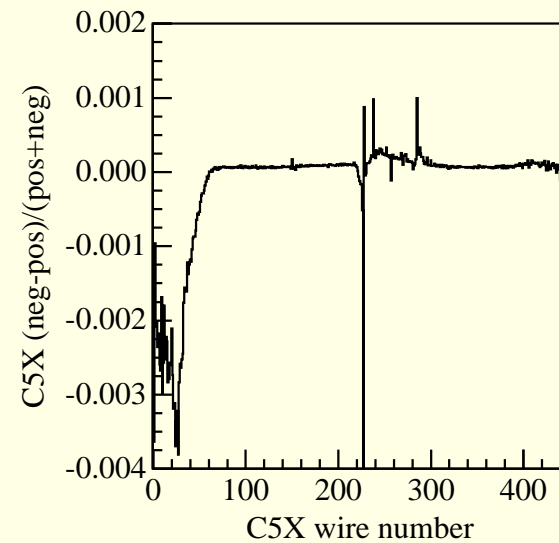
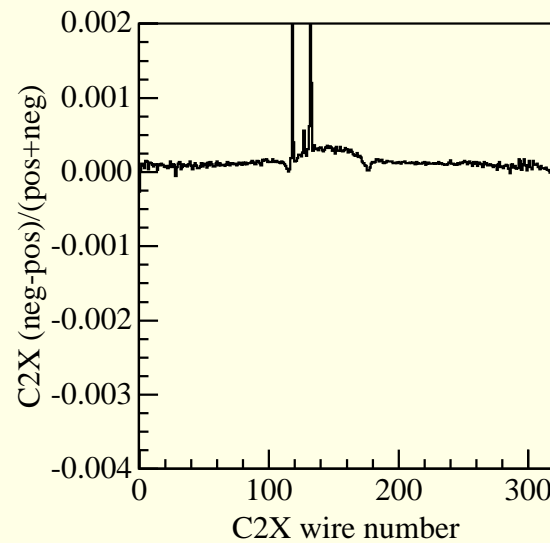
## Little Difference in PWC Efficiencies from + and - Running

- - data: solid line
- + data: dashed line
- 32 total planes  $\Rightarrow$  good redundancy

Absolute efficiency

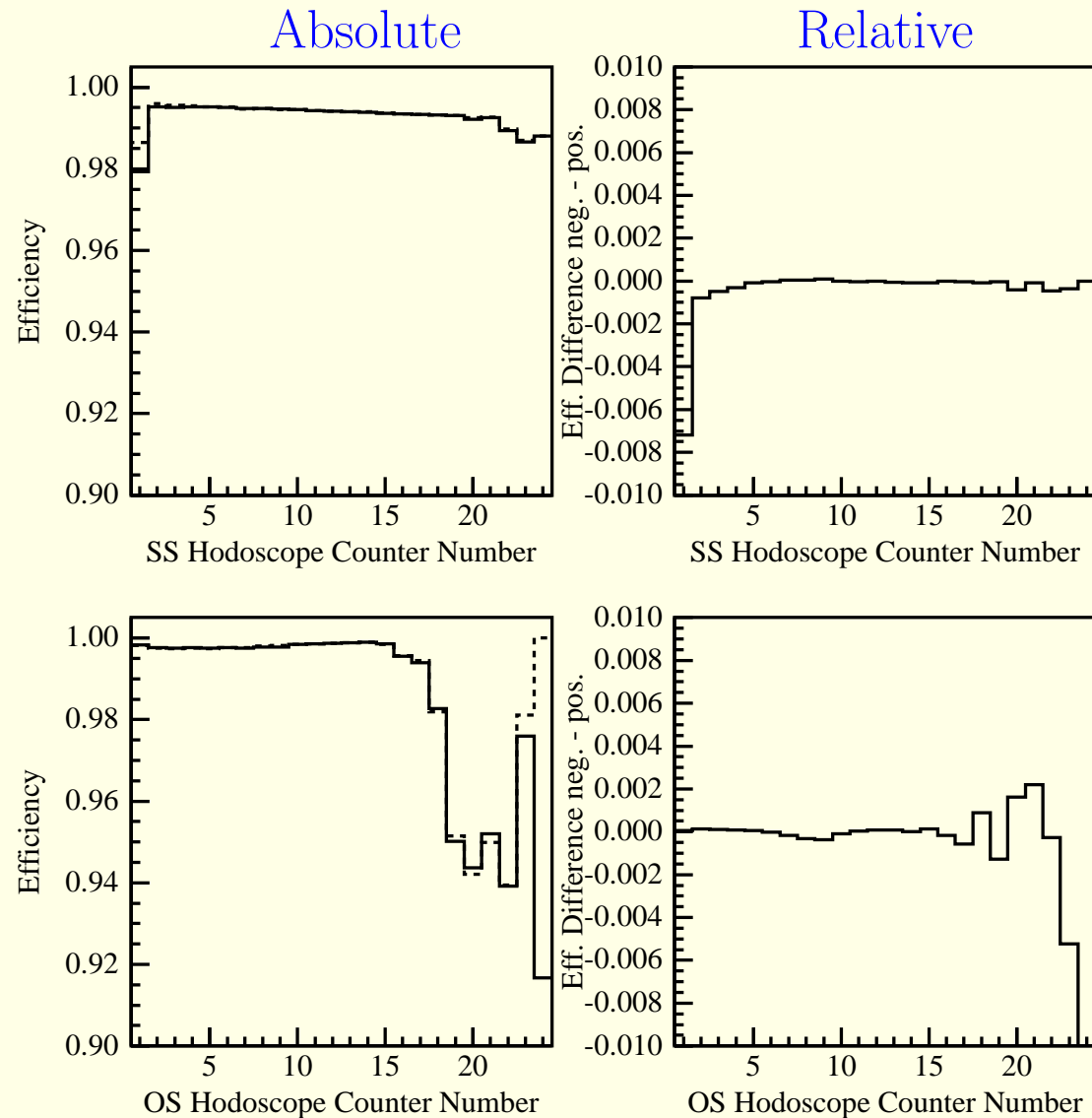


Difference (Neg. - Pos.)



## Little Difference in Hodo Efficiencies from + and - Running

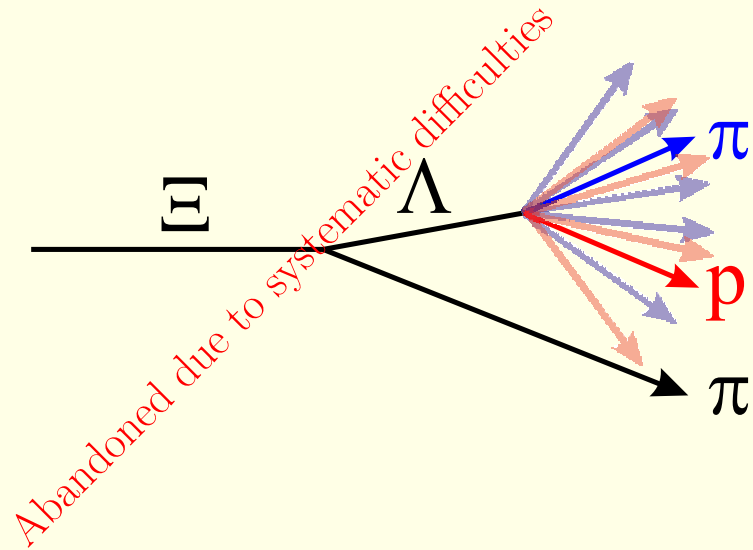
- - data: solid line.
- + data: dashed line.
- Differences where it matters  $< 0.1\%$ .
- Redundant counters make real inefficiencies vanishingly small.
- Two rows on OS side.
- Two particles on SS side.



## Two Different $CP$ Analyses Attempted

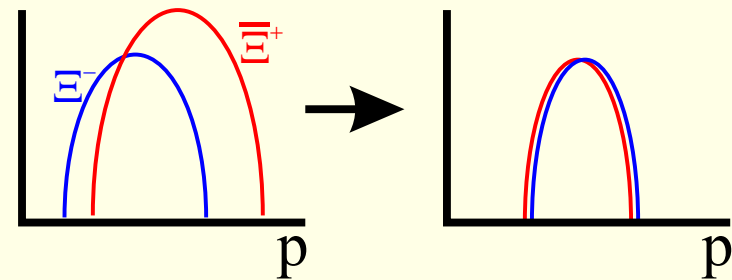
### Hybrid Monte Carlo Method:

- Compare corrected  $\cos \theta$  distributions.
- Take a real  $\Xi \rightarrow \Lambda \pi$ ,  $\Lambda \rightarrow p \pi$  event, discard proton and pion, generate 10 new unpolarized  $\Lambda$  decays.
- **Advantage:** Absolute measurement of  $\alpha_{\Lambda} \alpha_{\Xi}$ .
- **Disadvantage:** Monte Carlo must be very, very good, and fast:  $\sim 20$  billion events needed.



### Weighting Method:

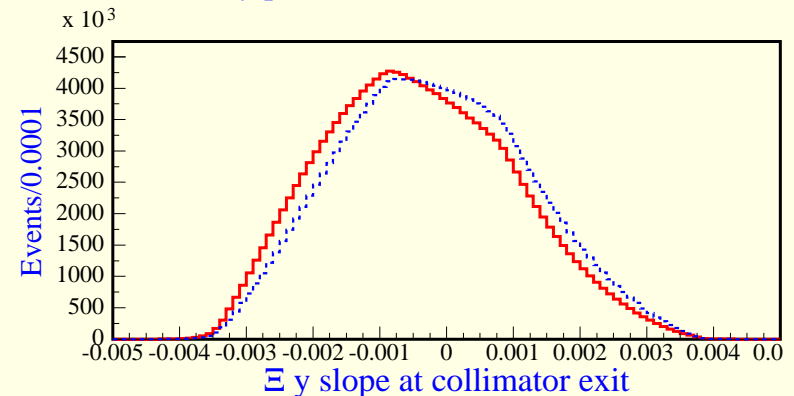
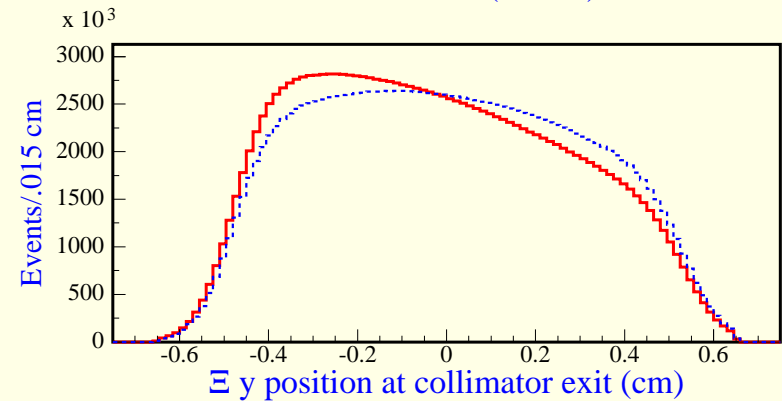
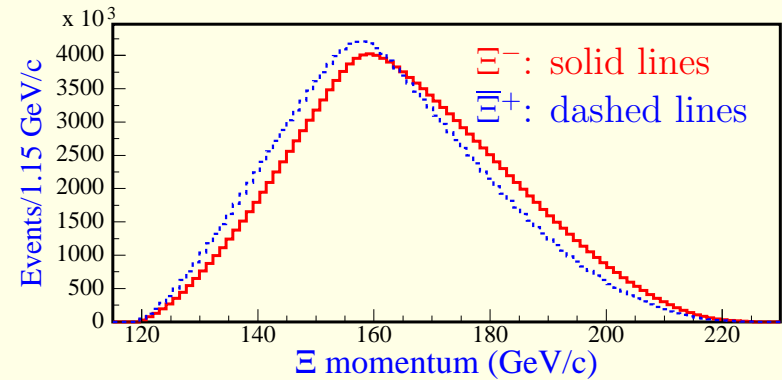
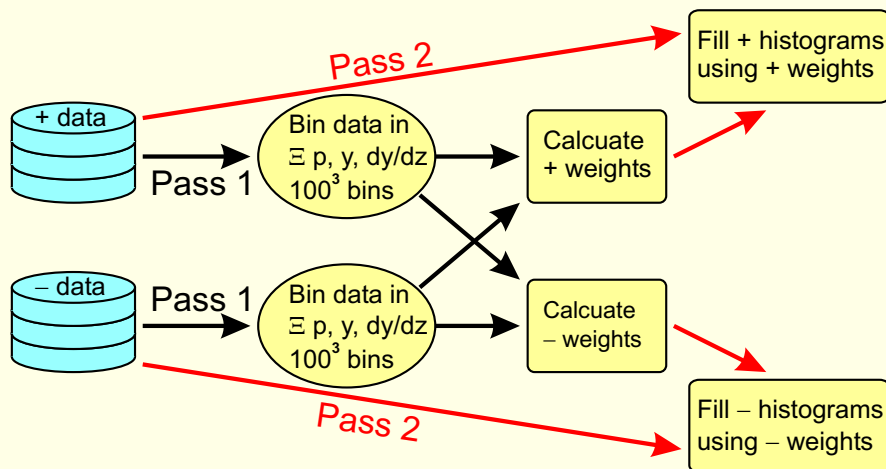
- Compare uncorrected  $\cos \theta$  distributions.
- Force the  $\Xi^-$  and  $\Xi^+$  events to have similar momentum and spatial distributions by appropriate weighting.
- **Advantage:** No Monte Carlo needed to measure apparatus acceptance, smaller statistical error.
- **Disadvantage:** inflexible, event-size dependent analysis.



Large data set,  $\sim 1$  billion events, in both cases makes the analysis difficult.

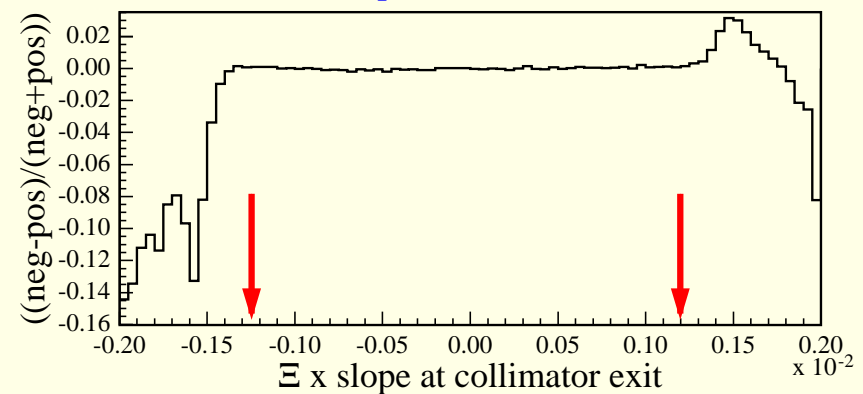
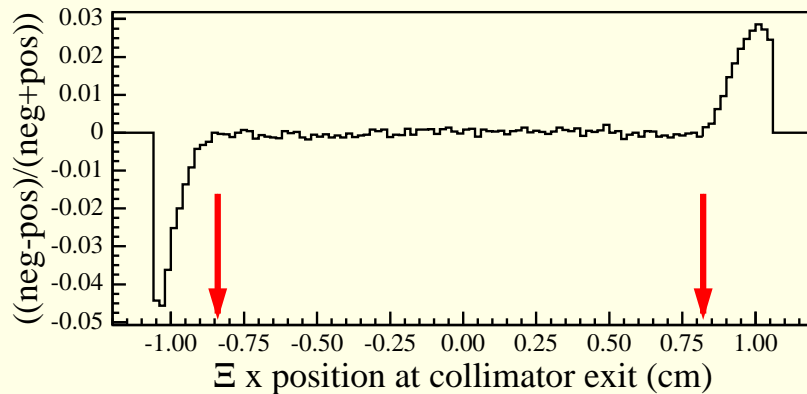
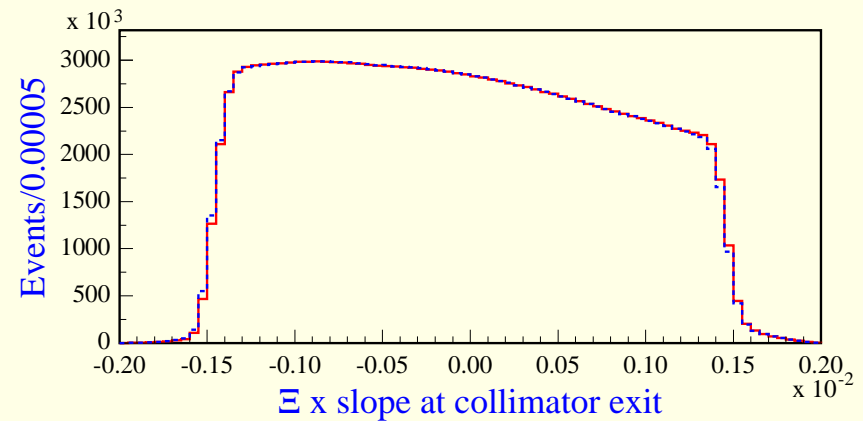
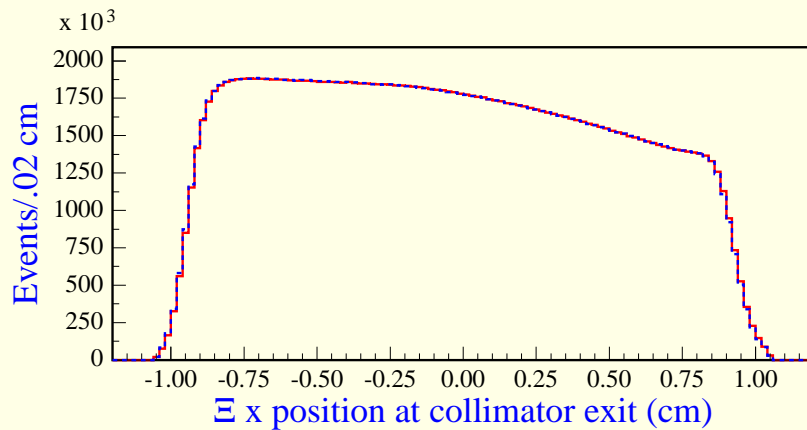
## Weighting Technique

- **Problem:** Geometrical acceptance identical for  $\Xi^-$  and  $\Xi^+$  decay products only if parent  $\Xi^-$  and  $\Xi^+$  have same momentum and inhabit the same phase space exiting the collimator.
- They are not the same due to different production dynamics.
- **Solution:** Weight the  $\Xi^-$  and  $\Xi^+$  events to force the two distributions to be identical.
- Momentum-dependent parameters of  $\Xi$  at collimator exit matched.
- $100 \times 100 \times 100 = 1 \times 10^6$  bins used.



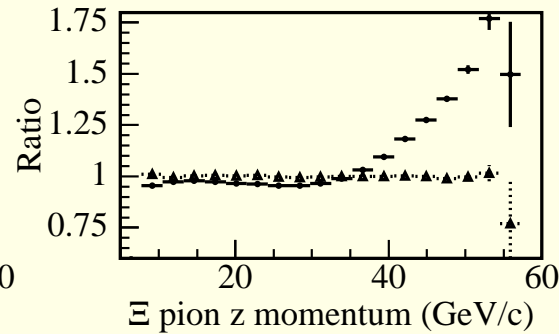
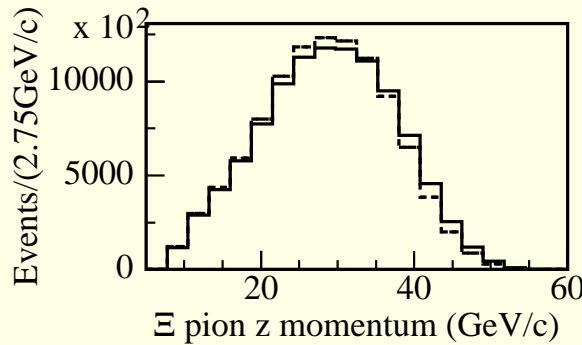
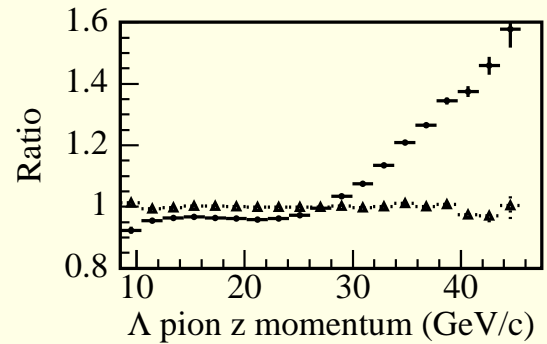
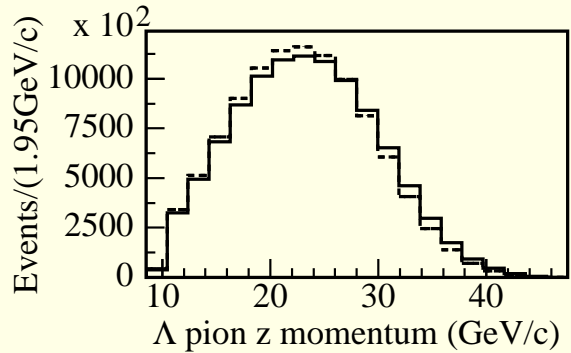
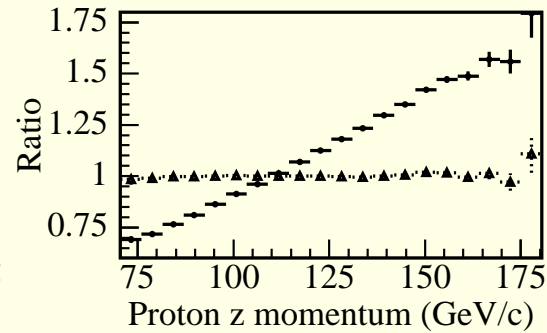
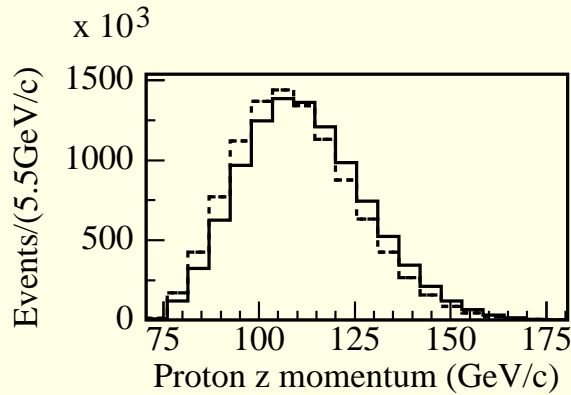
## $\Xi^-$ and $\Xi^+$ $x$ Slopes and Positions not Weighted

- Not momentum dependent  $\Rightarrow$  distributions almost identical
- Cut out regions where they are not.
- $\Xi^-$ : Solid lines
- $\Xi^+$ : Dashed lines



# Proton, $\Lambda$ -pion, $\Xi$ -pion Momenta Before/After Weighting

— Solid lines  
+ dashed lines



Ratio

o: before weighting  
 $\Delta$  after weighting

## Extracting the $CP$ Asymmetry

- Determine weighted proton and weighted antiproton  $\cos \theta$  distributions.

$$\frac{dN_-}{d \cos \theta_-} = A_- \frac{N_-}{2} (1 + \alpha_{\Xi} \alpha_{\Lambda} \cos \theta_-)$$

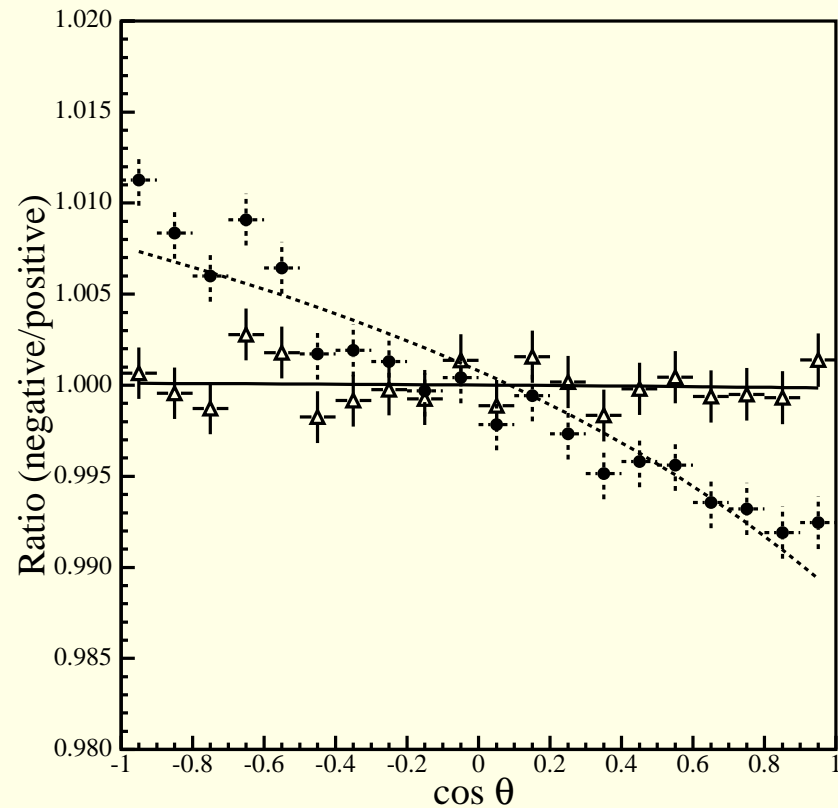
$$\frac{dN_+}{d \cos \theta_+} = A_+ \frac{N_+}{2} (1 + \bar{\alpha}_{\Xi} \bar{\alpha}_{\Lambda} \cos \theta_+)$$

- Assume the acceptances  $A_-$  and  $A_+$  have the same  $\cos \theta$  dependence.
- Take the ratio of proton and antiproton  $\cos \theta$  distributions: a nonzero slope is evidence of  $CP$  violation.
- Fit ratios to:

$$R(\theta, \delta) = C \frac{1 + \alpha_{\Xi} \alpha_{\Lambda} \cos \theta}{1 + (\alpha_{\Xi} \alpha_{\Lambda} - \delta) \cos \theta}$$

to extract asymmetry  $\delta$ :

$$\begin{aligned} \delta &= \alpha_{\Xi} \alpha_{\Lambda} - \bar{\alpha}_{\Xi} \bar{\alpha}_{\Lambda} \\ A_{\Xi\Lambda} &= \frac{\delta}{\alpha_{\Xi} \alpha_{\Lambda} + \bar{\alpha}_{\Xi} \bar{\alpha}_{\Lambda}} = \frac{\delta}{2\alpha_{\Xi} \alpha_{\Lambda}} \\ &= 1.71 \delta \end{aligned}$$



Proton  $\cos \theta$  ratio before (●) and after (Δ) weighting.

## Monte Carlo Tests

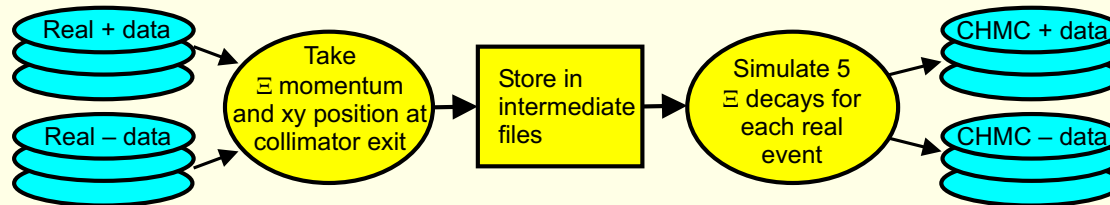
**Important!** Monte Carlo only used to:

- Verify code and algorithm.
- Study a few systematics.

**Final result has no Monte Carlo dependence!**

**Problem:** How do you generate  $\sim 1$  billion MC events?

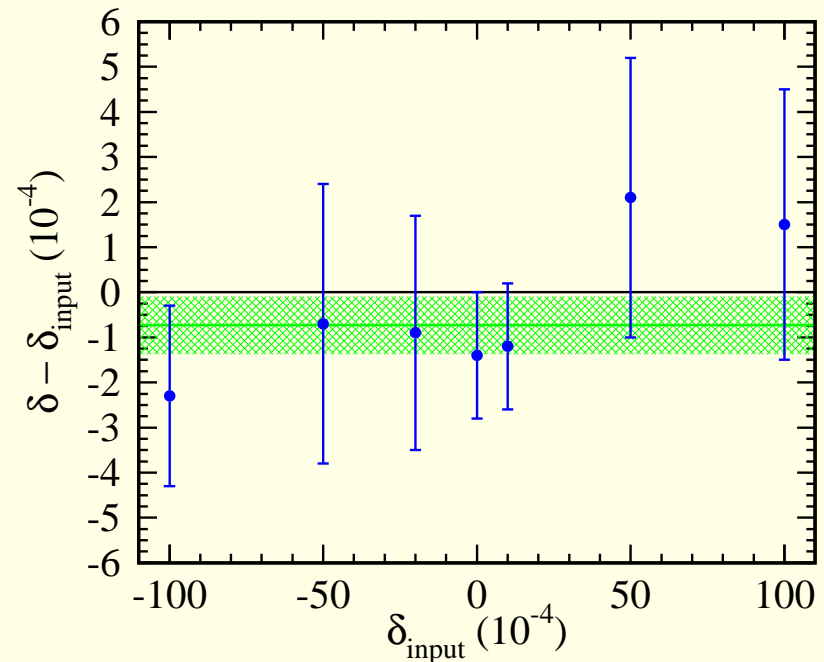
**Solution:**



We get the input asymmetry back  $\Rightarrow$

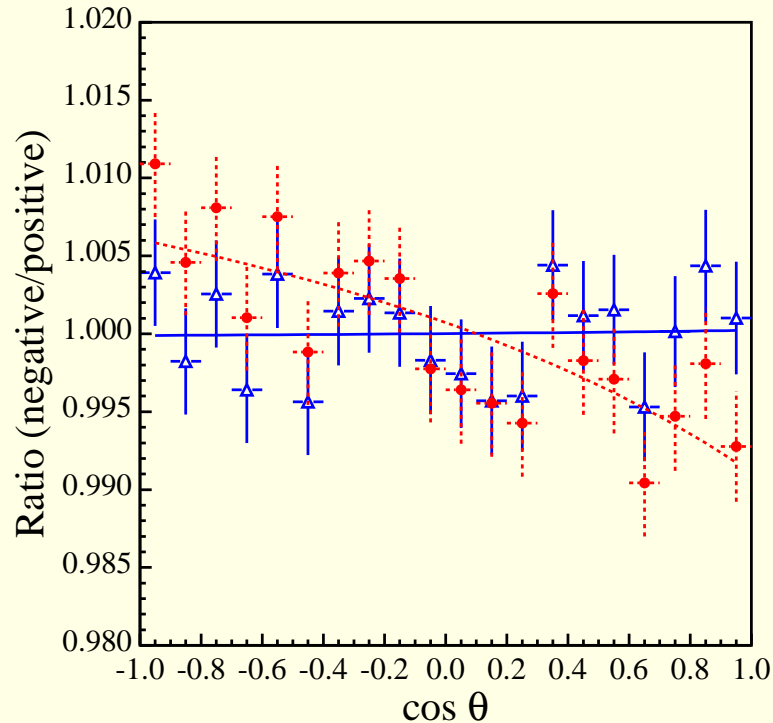
$$\delta = (-0.73 \pm 0.64) \times 10^{-4}$$

$$A_{\Xi\Lambda} = (1.24 \pm 1.09) \times 10^{-4}$$



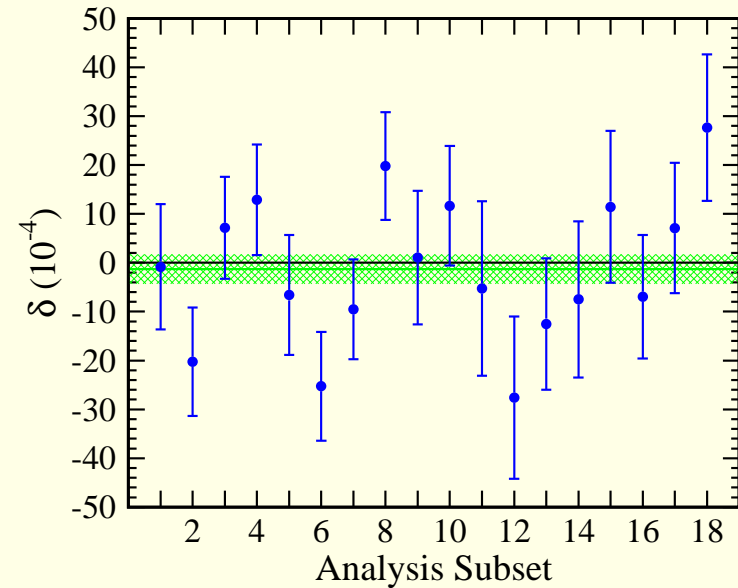
## The $CP$ Asymmetry $A_{\Xi\Lambda}$ from Weighting Method

- Data broken up into 18 sets, each with positive and negative events.
- No acceptance corrections.
- No efficiency corrections.
- No background subtraction.



Proton  $\cos\theta$  ratio before ( $\bullet$ ) after ( $\triangle$ ) weighting, from Analysis Set 1

$$\delta = \alpha_{\Xi}\alpha_{\Lambda} - \bar{\alpha}_{\Xi}\bar{\alpha}_{\Lambda}$$

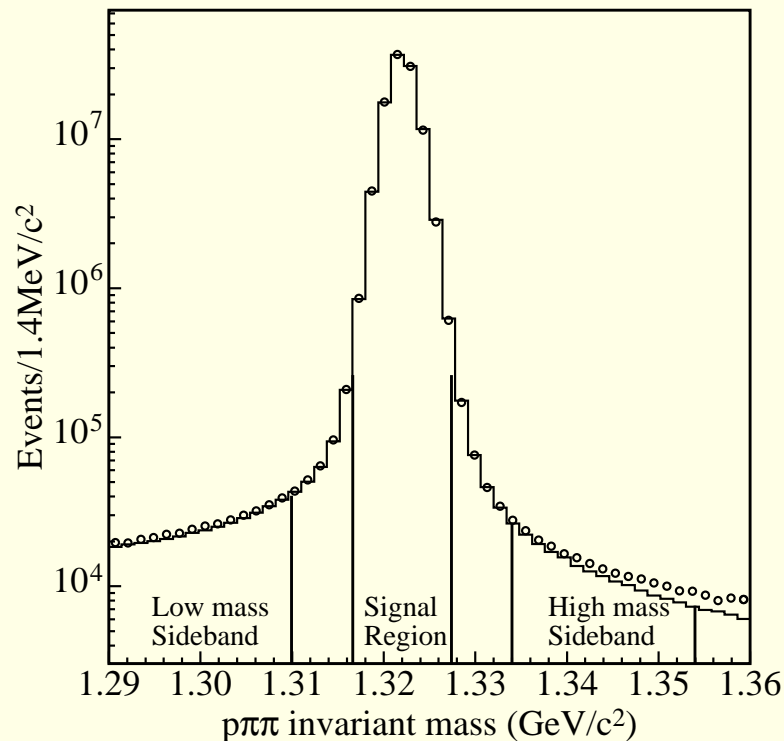


Weighted average of all 18 data sets:

$$\begin{aligned} \delta &= (-1.3 \pm 3.0) \times 10^{-4} \\ A_{\Xi\Lambda} &= (2.2 \pm 5.1) \times 10^{-4} \\ \chi^2 &= 24 \end{aligned}$$

## Background Subtraction Has Little Effect

- Triple Gaussian fit with fourth-order polynomial for background.
- Background fraction:
  - $\Xi^-$ : 0.43% (lines)
  - $\Xi^+$ : 0.41% (circles)

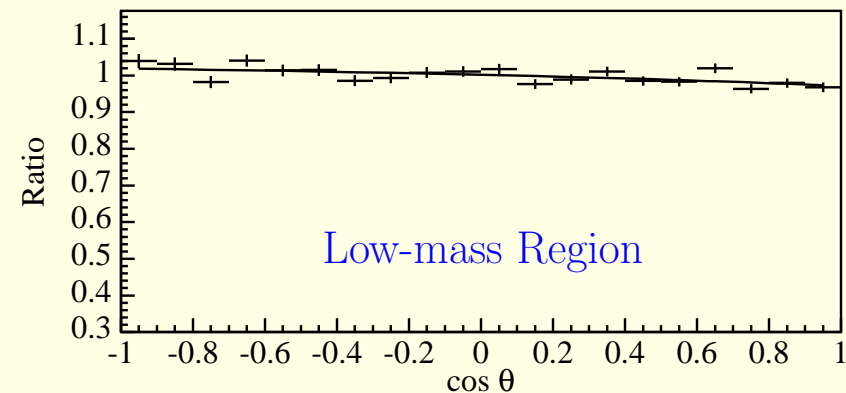
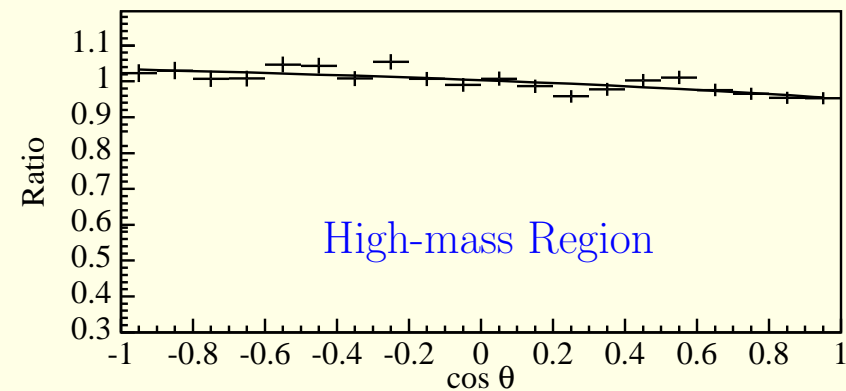


Low mass:  $\delta = (-2.2 \pm 0.5) \times 10^{-2}$

High mass:  $\delta = (-3.8 \pm 0.7) \times 10^{-2}$

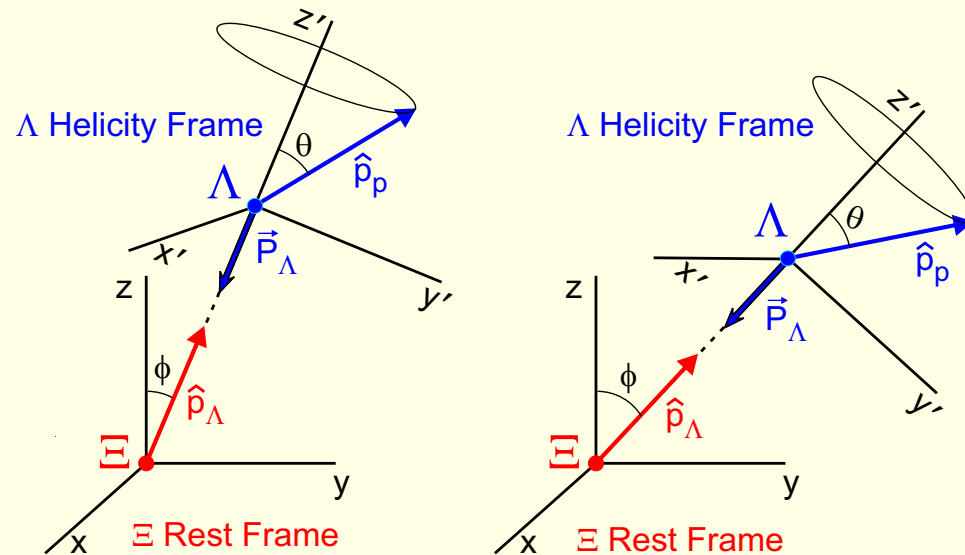
- Weighted background asymmetry:

$$A_{\Xi\Lambda}(bs) = (0.0 \pm 5.1) \times 10^{-4}$$

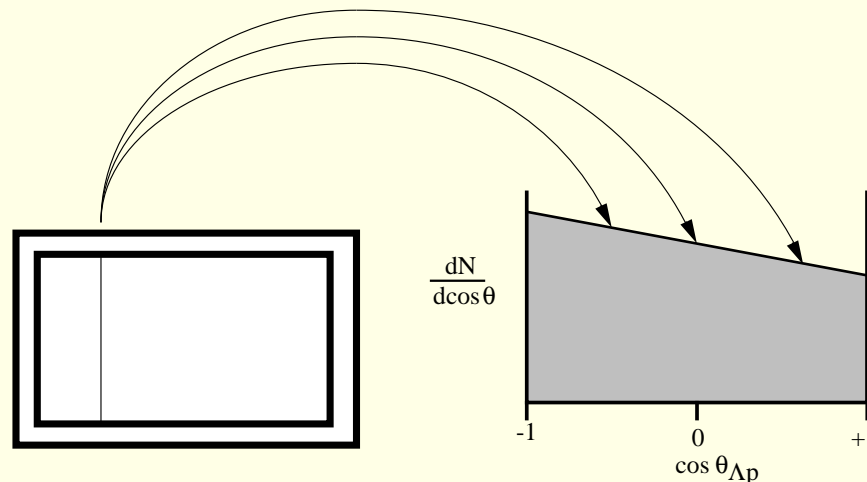


## Helicity Frame Analysis Naturally Minimizes Biases

- The **helicity frame** axes changes from event to event since we always define the polar axis to be the direction of the  $\Lambda$  momentum in the  $\Xi$  rest frame.



- Acceptance differences localized in a particular part of the apparatus do **not** map into a particular part of the proton (antiproton)  $\cos \theta$  distribution.



**Important!** Overall acceptance differences do not cause any biases.

## Weighted Analysis Bias Error Summary

Systematic	Method	$\delta A_{\Xi\Lambda}(10^{-4})$
Analyzing Magnets field uncertainties	Data	2.4
Calorimeter inefficiency uncertainty	Data	2.1
Validation of analysis code	CHMC	1.9
Collimator exit $x$ slope cut	Data	1.4
Collimator exit $x$ position cut	Data	1.2
PWC inefficiency uncertainty	CHMC	1.0
Hodoscope inefficiency uncertainty	Data	0.3
Particle/antiparticle interaction differences	MC	0.9
Momentum weights bin size	Data	0.4
Background subtraction uncertainty	Data	0.3
Error on $\alpha\alpha_{PDG}$	Data	0.03
Polarization	MC	negligible
Earth's magnetic field	CHMC	negligible
Total systematic error		4.4

## Results from $CP$ Violation Search

### Weighting Technique:

- $\sim 10\%$  total data sample
- selected from end of 1999 run
- 118.6 million  $\Xi^-$
- 41.9 million  $\Xi^+$
- no acceptance or efficiency corrections

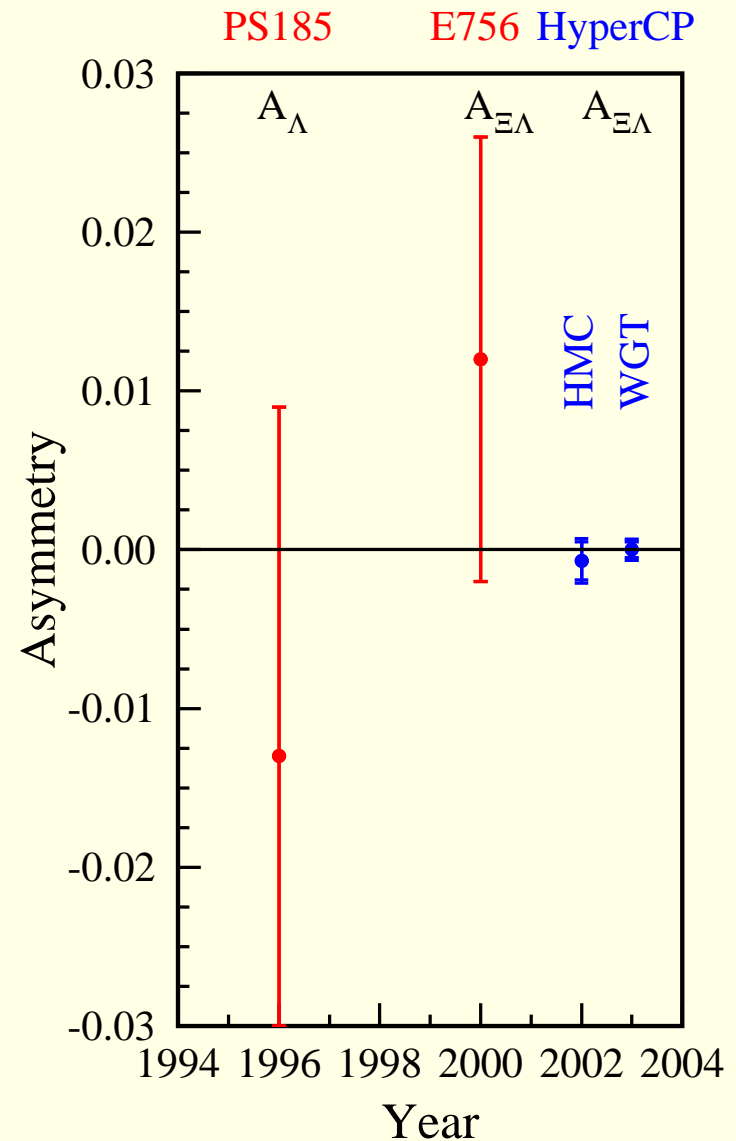
$$A_{\Xi\Lambda} = [0.0 \pm 5.1(\text{stat}) \pm 4.4(\text{syst})] \times 10^{-4}$$

### Check with HMC Technique:

- $\sim 5\%$  of the total data sample
- prescaled selection of 1997 and 1999
- 15 million  $\Xi^-$
- 30 million  $\Xi^+$

$$A_{\Xi\Lambda} = [-7 \pm 12(\text{stat}) \pm 6.2(\text{syst})] \times 10^{-4}$$

$\Rightarrow 20\times$  improvement on previous result.



## Conclusions and Outlook

- Hyperon  $CP$  violation measurements probing limits not constrained by Kaon, B, or EDM measurements.

“...we can then conclude that the available preliminary measurement by HyperCP has already begun to probe the parity-even contributions better than  $\epsilon$  does.”

Tandean (2004)

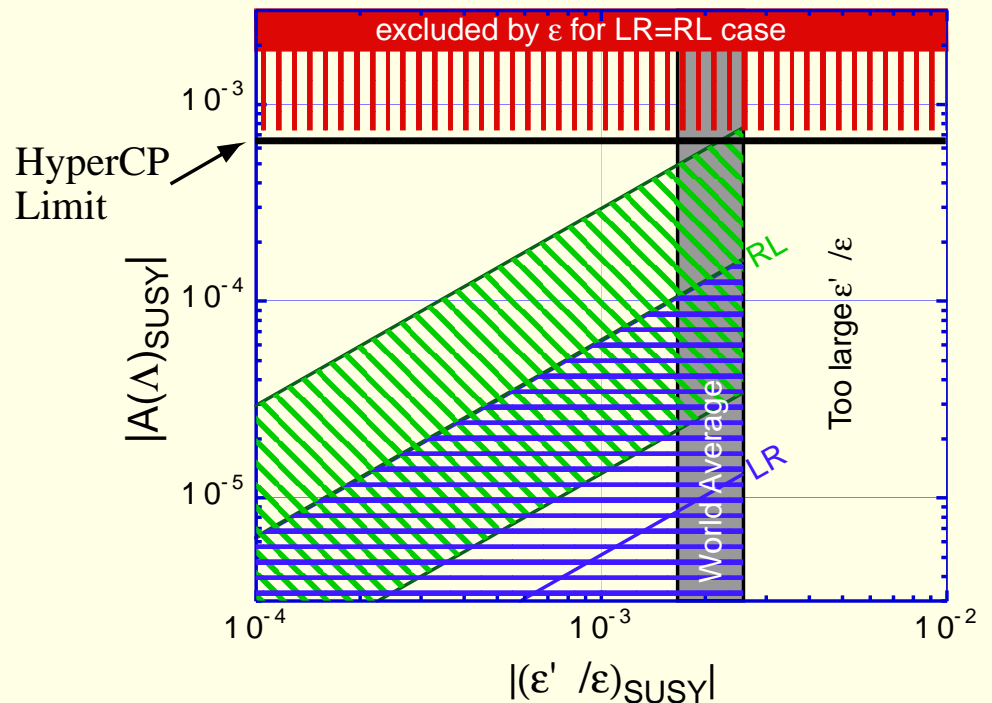
- $HyperCP$ , in particular, the first dedicated hyperon  $CP$  violation experiment, has pushed into the region where SUSY models allow an effect.
- $HyperCP$  finds no evidence of  $CP$  violation in  $\Xi^\pm$  and  $\Lambda$  decays:

$$\delta A_{\Xi\Lambda} = (0.0 \pm 5.1 \pm 4.4) \times 10^{-4}$$

- Shortly we should push our statistical limit to:

$$\delta A_{\Xi\Lambda} \approx 2 \times 10^{-4}$$

two orders of magnitude better than the present limit.



## Backup Slides

## Measurement of the $\Lambda$ - $\pi$ Phase Shift

- This is done by analyzing the  $\Lambda$  decay distribution from 144 million **polarized**  $\Xi^-$ 's.
- $\Lambda$  has three components of polarization:

$$\vec{P}_\Lambda = \frac{(\alpha_\Xi + \vec{P}_\Xi \cdot \hat{p}_\Lambda) \hat{p}_\Lambda + \beta_\Xi (\vec{P}_\Xi \times \hat{p}_\Lambda) + \gamma_\Xi (\hat{p}_\Lambda \times (\vec{P}_\Xi \times \hat{p}_\Lambda))}{(1 + \alpha_\Xi \vec{P}_\Xi \cdot \hat{p}_\Lambda)}$$

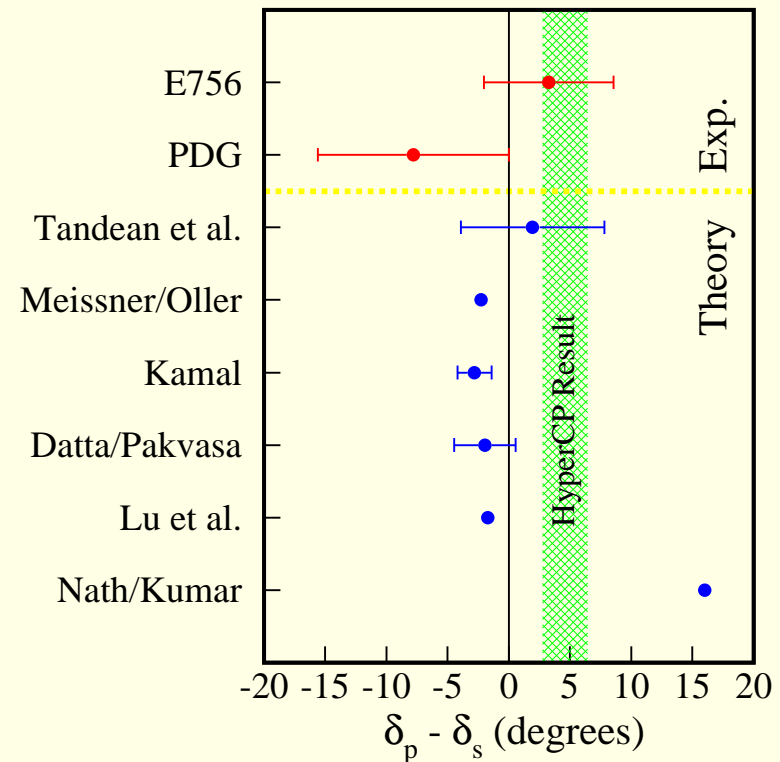
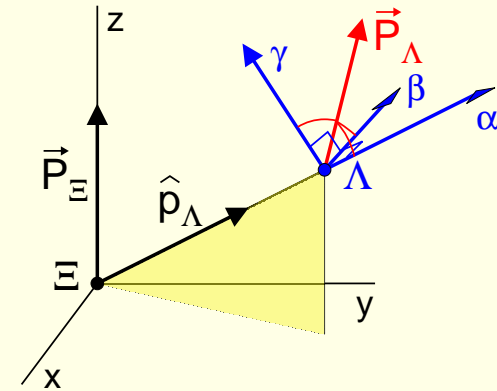
$$\beta_\Xi = -0.037 \pm 0.011(\text{stat}) \pm 0.010(\text{syst})$$

$$\gamma_\Xi = 0.888 \pm 0.0004(\text{stat}) \pm 0.006(\text{syst})$$

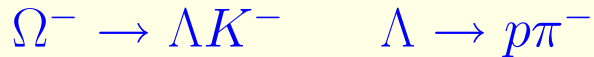
- Using the known value of  $\alpha_\Xi$ :

$$\delta_P - \delta_S = \tan^{-1} \left( \frac{\beta_\Xi}{\alpha_\Xi} \right) = (4.6 \pm 1.4 \pm 1.2)^\circ$$

- First non-zero measurement of phaseshift.
- This is about the same magnitude as the  $p$ - $\pi$  phase shift:
  - $\Rightarrow$   $CP$  equally likely in  $\Xi$  and  $\Lambda$  decays.
  - $\Rightarrow$   $CP$  predictions underestimated,
  - $\Rightarrow$   $\chi$ PT calculations off.



## Search for Parity Violation in $\Omega^- \rightarrow \Lambda K^-$ Decays



- Although spin-3/2,  $\Omega^- \rightarrow \Lambda K^-$  decay goes much like the other hyperon two-body decays:

$$\frac{dP}{d \cos \theta} = \frac{1}{2}(1 + \alpha_\Omega P_\Omega \cos \theta)$$

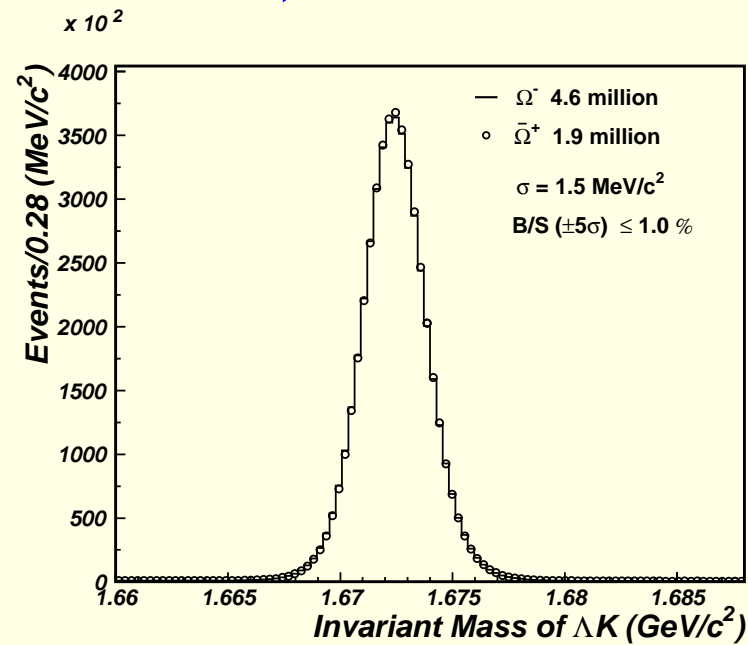
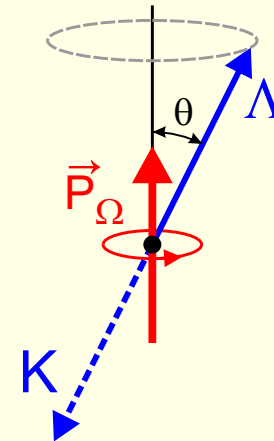
- Here:

$$\alpha_\Omega = \frac{2\text{Re}(P^* D)}{|P|^2 + |D|^2}$$

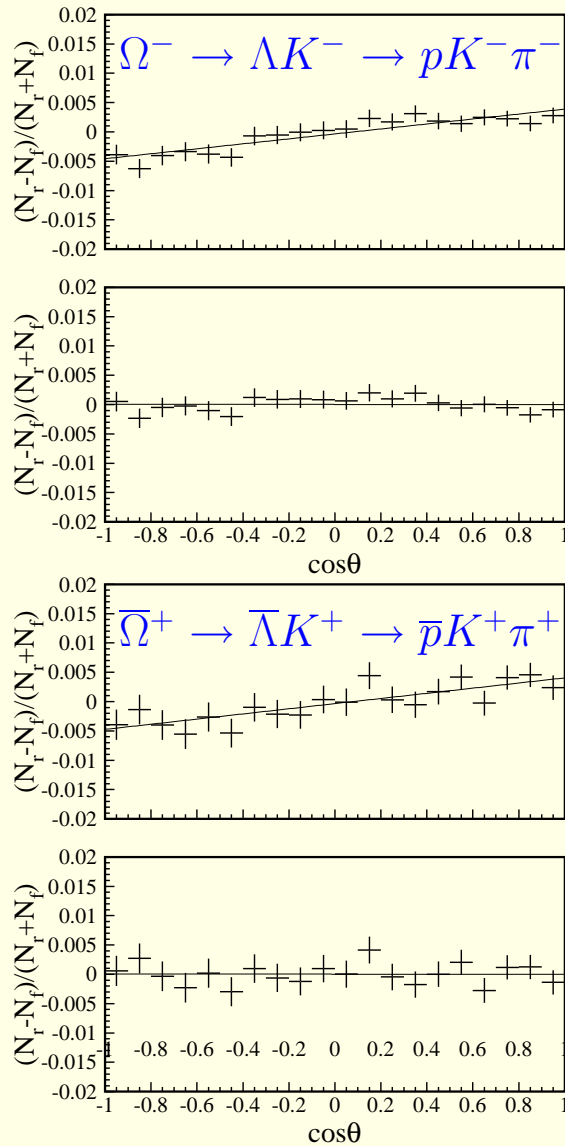
- A non-zero  $\alpha_\Omega$  indicates parity violation.
- All other hyperons have non-zero  $\alpha$  parameters; only the  $\Omega^-$  has resisted efforts to find an asymmetrical decay distribution.
- HyperCP* is measuring  $\alpha_\Omega$  using unpolarized  $\Omega^-$ 's through the polarization given to the daughter  $\Lambda$ , which is  $\alpha_\Omega$ :

$$\frac{dP}{d \cos \theta} = \frac{1}{2}(1 + \alpha_\Omega \alpha_\Lambda \cos \theta)$$

- Large data sample, little background.

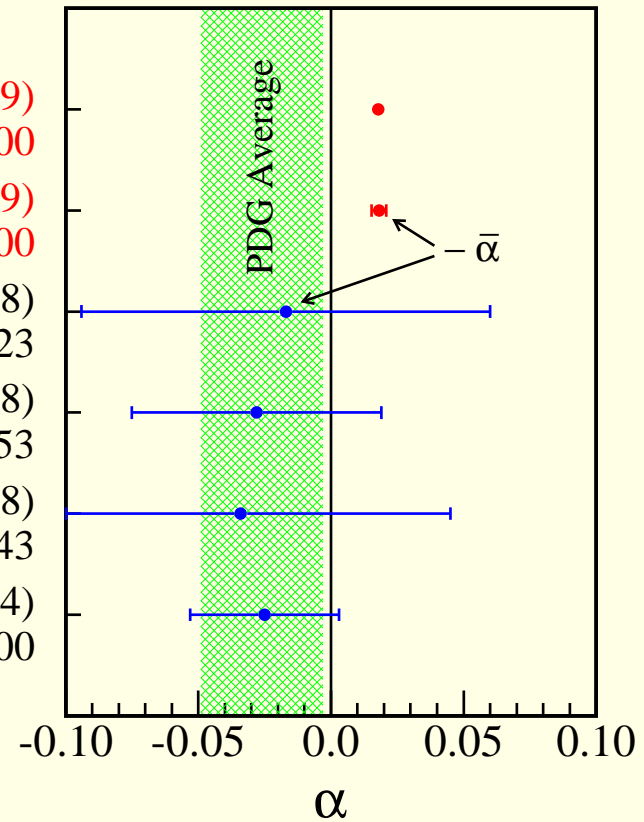


# Preliminary Measurement of $\alpha_\Omega$ and $\bar{\alpha}_\Omega$ in $\Omega^- \rightarrow \Lambda K^-$ Decays



1999 :  $\alpha_\Omega = [ 1.78 \pm 0.19(\text{stat}) \pm 0.10(\text{syst}) ] \times 10^{-2}$   
 1999 :  $\bar{\alpha}_\Omega = [ -1.81 \pm 0.28(\text{stat}) ] \times 10^{-2}$

HyperCP (1999)  
 4,500,000  
 HyperCP (1999)  
 1,900,000  
 FNAL-756 (1998)  
 1823  
 FNAL-756 (1998)  
 6953  
 FNAL-620 (1988)  
 1743  
 CERN (1984)  
 12,000



- First evidence of parity violation in  $\Omega^-$  decays.
- Can search for  $CP$  violation in  $\Omega^-/\bar{\Omega}^+$  decays.