

# Neutrinoless double beta decay: from rates to Majorana masses

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# New era of neutrino physics

## 1. Atmospheric neutrino oscillations

*(in particular zenith angle dependence of the muon neutrino flux)*

## 2. Solar neutrino deficit

*(in particular the difference in neutrino fluxes deduced from the charged and neutral current reaction rates)*

## 3. KamLAND

*(reactor  $\bar{\nu}_e$  disappearance, oscillations seen with the terrestrial sources)*

## 4. LSND oscillation observations

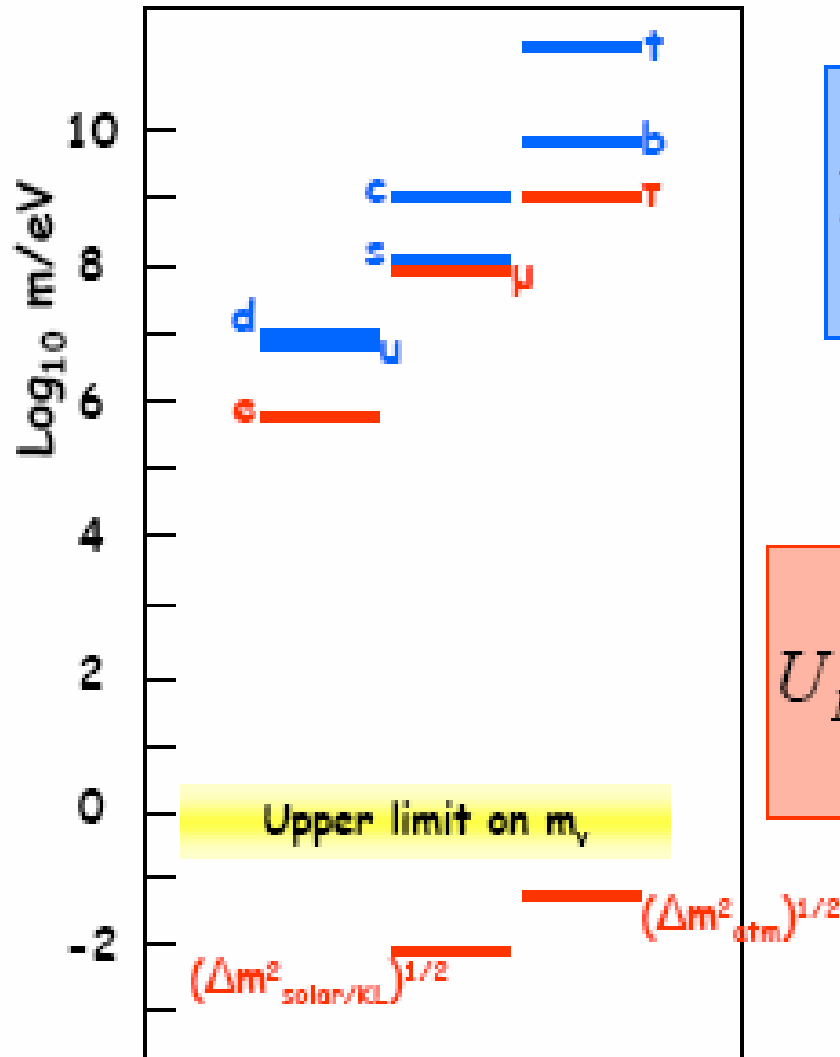
*(unconfirmed, but soon to be checked; if correct all bets are off)*

# Consequences:

- At least some neutrinos are massive.  
*(lower limits are 50 and 10 meV)*
- Mixing exists  
*(two mixing angles are large, one is small)*
- But we do not know the absolute mass scale
- We do not know the behavior under charge conjugation

$0\nu\beta\beta$  might help with the answers

So the neutrino sector is really strange !

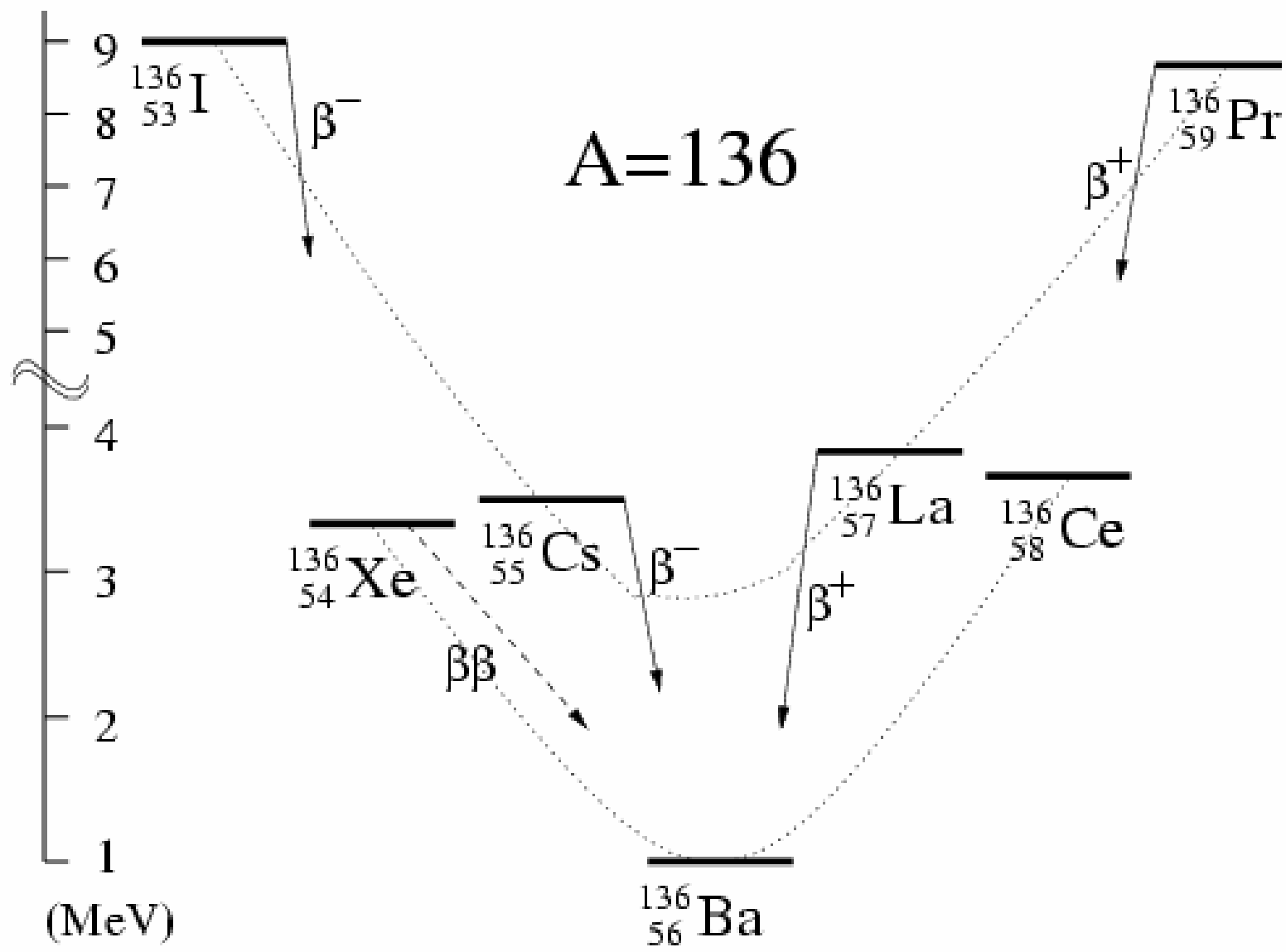


$$V_{\text{CKM}} \cong \begin{pmatrix} 0.97 & 0.22 & 0.00 \\ 0.22 & 1.00 & 0.04 \\ 0.01 & 0.04 & 1.00 \end{pmatrix}$$

Adapted from PDG 2002

$$U_{\text{PMNS}} \cong \begin{pmatrix} 0.84 & 0.54 & 0.1 \\ -0.44 & 0.56 & 0.71 \\ 0.32 & -0.63 & 0.71 \end{pmatrix}$$

(entries evaluated for  $U_{e3} = 0.1$ , near the middle of allowed range)



# Candidate Nuclei for Double Beta Decay

Candidate

Q  
(MeV)

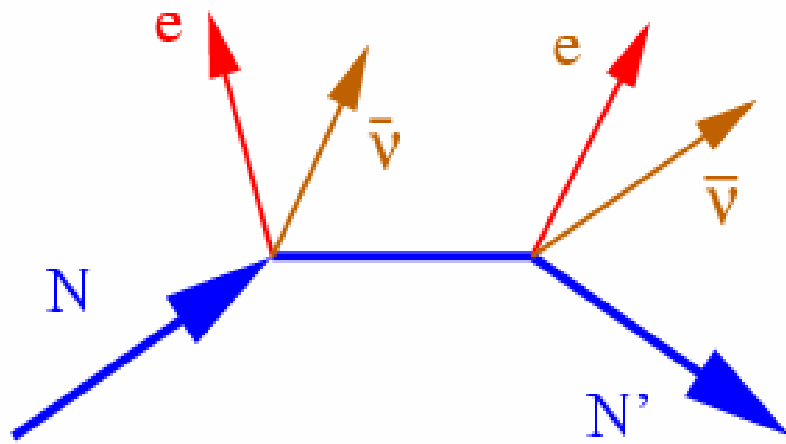
Abund.  
(%)

$^{48}\text{Ca} \rightarrow ^{48}\text{Ti}$	4.271	0.187
$^{76}\text{Ge} \rightarrow ^{76}\text{Se}$	2.040	7.8
$^{82}\text{Se} \rightarrow ^{82}\text{Kr}$	2.995	9.2
$^{96}\text{Zr} \rightarrow ^{96}\text{Mo}$	3.350	2.8
$^{100}\text{Mo} \rightarrow ^{100}\text{Ru}$	3.034	9.6
$^{110}\text{Pd} \rightarrow ^{110}\text{Cd}$	2.013	11.8
$^{116}\text{Cd} \rightarrow ^{116}\text{Sn}$	2.802	7.5
$^{124}\text{Sn} \rightarrow ^{124}\text{Te}$	2.228	5.64
$^{130}\text{Te} \rightarrow ^{130}\text{Xe}$	2.533	34.5
$^{136}\text{Xe} \rightarrow ^{136}\text{Ba}$	2.479	8.9
$^{150}\text{Nd} \rightarrow ^{150}\text{Sm}$	3.367	5.6

Issues include : Q value, abundance, ease of purification (chemical and isotopic), radioactivity (incl. cosmogenesis), & experimental ease of use.

# Most sensitive neutrino mass measurements can be obtained from double-beta decay

$2\nu \beta\beta$  decay: a standard  
process in nuclear physics



$0\nu \beta\beta$  decay: a hypothetical  
process

→  $m_\nu \neq 0$   
→  $\bar{\nu} = \nu$  since helicity  
has to "flip"

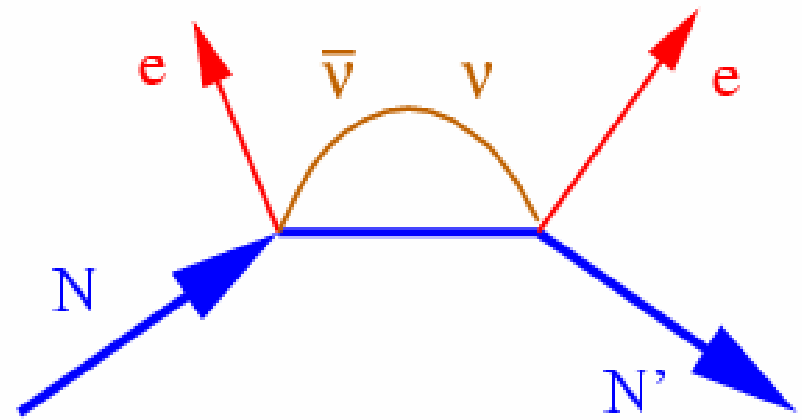


FIG. 1. How to separate the  $2\nu$  and  $0\nu$  decays? Dotted line is the sum electron spectrum for the  $2\nu$  decay, full line is for the  $0\nu$  decay. 5% resolution is assumed. The ratio of rates is  $10^2$ , and  $10^6$  (goal) in the insert.

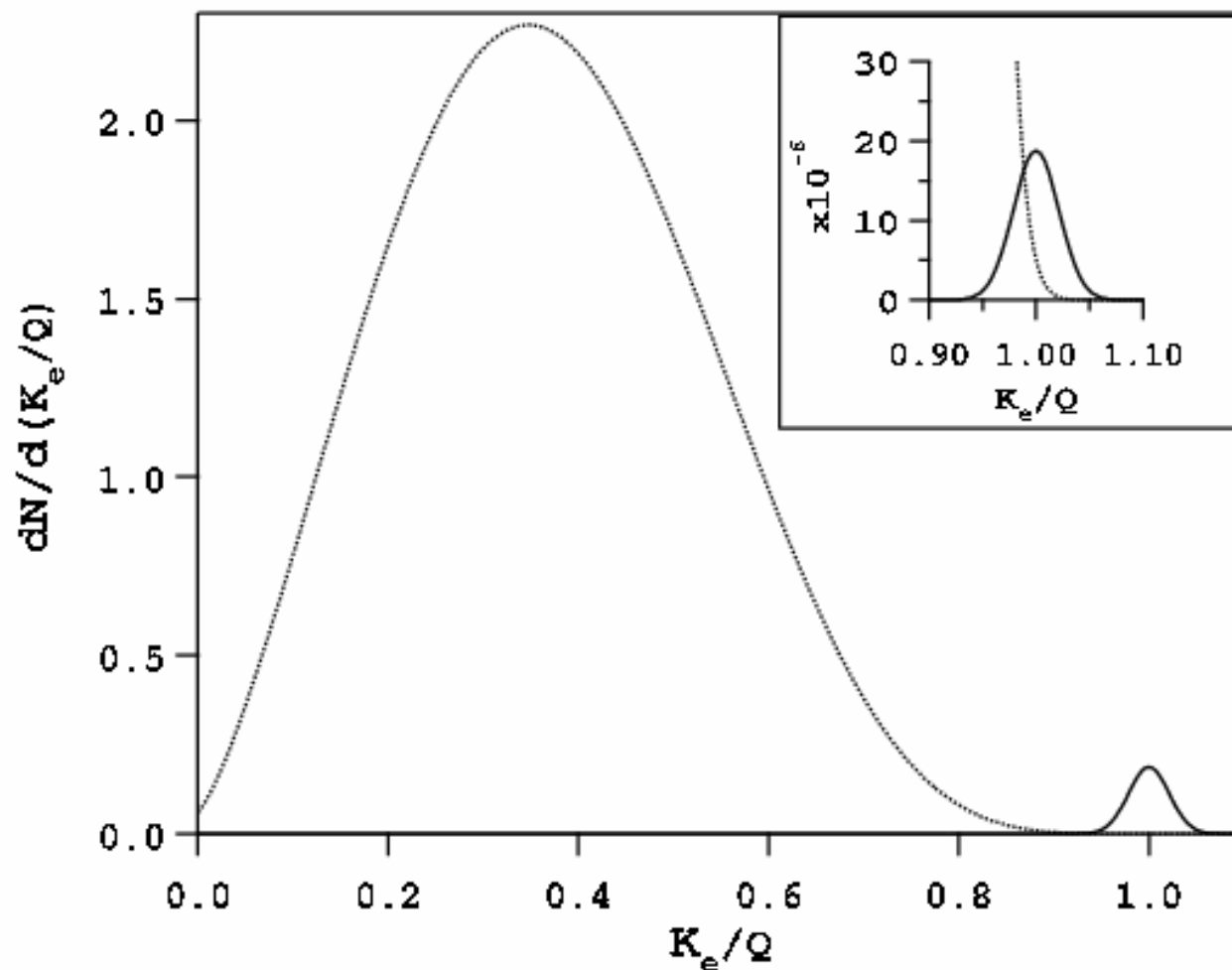


Table 1: Summary of experimentally measured  $2\nu\beta\beta$  half-lives and matrix elements ( $^{136}\text{Xe}$  is an important exception where a limit is quoted).

Isotope	$T_{1/2}^{2\nu}$ (y)	References	$M_{GT}^{2\nu}$ (MeV $^{-1}$ )
$^{48}\text{Ca}$	$(4.2 \pm 1.2) \times 10^{19}$	BAL96,BRU00	0.05
$^{76}\text{Ge}$	$(1.3 \pm 0.1) \times 10^{21}$	KLA01a,AVI91,AAL96	0.15
$^{82}\text{Se}$	$(9.2 \pm 1.0) \times 10^{19}$	ELL92,ARN98	0.10
$^{96}\text{Zr}^\dagger$	$(1.4_{-0.5}^{+3.5}) \times 10^{19}$	ARN99,KAW93,Wieser01	0.12
$^{100}\text{Mo}$	$(8.0 \pm 0.6) \times 10^{18}$	DAS95,EJI91a,EJI91c, DES97,ALS97,ASH01	0.22
$^{116}\text{Cd}$	$(3.2 \pm 0.3) \times 10^{19}$	ARN96,DAN00,EJI95	0.12
$^{128}\text{Te}^{(1)}$	$(7.2 \pm 0.3) \times 10^{24}$	BER93,CRU93	0.025
$^{130}\text{Te}^{(2)}$	$(2.7 \pm 0.1) \times 10^{21}$	BER93	0.017
$^{136}\text{Xe}$	$> 8.1 \times 10^{20}$ (90% CL)	GAV00	$< 0.03$
$^{150}\text{Nd}^\dagger$	$7.0_{-0.3}^{+11.8} \times 10^{18}$	DES97,ART95	0.07
$^{238}\text{U}^{(3)}$	$(2.0 \pm 0.6) \times 10^{21}$	TUR91	0.05

<sup>(1)</sup>deduced from the geochemically determined half-life ratio  $^{128}\text{Te}/^{130}\text{Te}$

<sup>(2)</sup>geochemical result includes all decay modes; other geochemical determinations only marginally agree

<sup>(3)</sup>radiochemical result, again for all decay modes

$$M_{GT}^{2\nu} = \sum_m \frac{\langle f | \vec{\sigma} \tau^+ | m \rangle \cdot \langle m | \vec{\sigma} \tau^+ | i \rangle}{E_m - (M_i + M_f)/2}$$

The  $2\nu\beta\beta$  matrix element consists of two amplitudes:  
 $\langle m || \sigma\tau^+ || i \rangle$  corresponds to the  $\beta^-$  strength of initial nucleus  
 $\langle f || \sigma\tau^+ || m \rangle$  corresponds to the  $\beta^+$  strength of final nucleus  
 $E_m$  is the energy of the intermediate  $1^+$  state

Note that the sum converges faster than expected, often already the first  $1^+$  state gives a dominant contribution.

The  $\beta^-$  and  $\beta^+$  strengths are constrained by the Ikeda sum rule  
 $S(\beta^-) - S(\beta^+) = 3(N-Z)$

Note that for the  $2\nu\beta\beta$  we need also signs, not just magnitudes of the GT matrix elements.

**Various processes could contribute to the  $0\nu$  decay amplitude:**

**(light Majorana neutrinos, heavy neutrinos, right-handed current interactions, various SUSY particles, etc.) All of them imply “Physics beyond the Standard Model”.**

**Since we know that light neutrinos exist, we primarily consider the process mediated by them.**

**If heavy virtual particles are exchanged, short range physics issues are involved, for the treatment in EFT see Prezeau, Ramsey-Musolf and P.V., Phys.Rev.D68,035501(2003)**

**For relation to the Lepton Flavor Violation processes see Cirigliano, Kurylov, Ramsey-Musolf and P.V. hep-ph/0406199**

In  $0\nu\beta\beta$  decay with the exchange of light Majorana neutrinos,

$$\langle m_{\beta\beta} \rangle = \sum_i |U_{ei}|^2 m_i e^{i\alpha(i)},$$

where  $\alpha(i)$  are the CP phases of  $U_{ei}$ .

These phases are relevant only for Majorana neutrinos; they do not affect flavor oscillations.

Independently of these phases

$$\text{Max}[2|U_{ei}|^2 m_i] - \sum |U_{ei}|^2 \ll \langle m_{\beta\beta} \rangle \ll \sum |U_{ei}|^2 m_i .$$

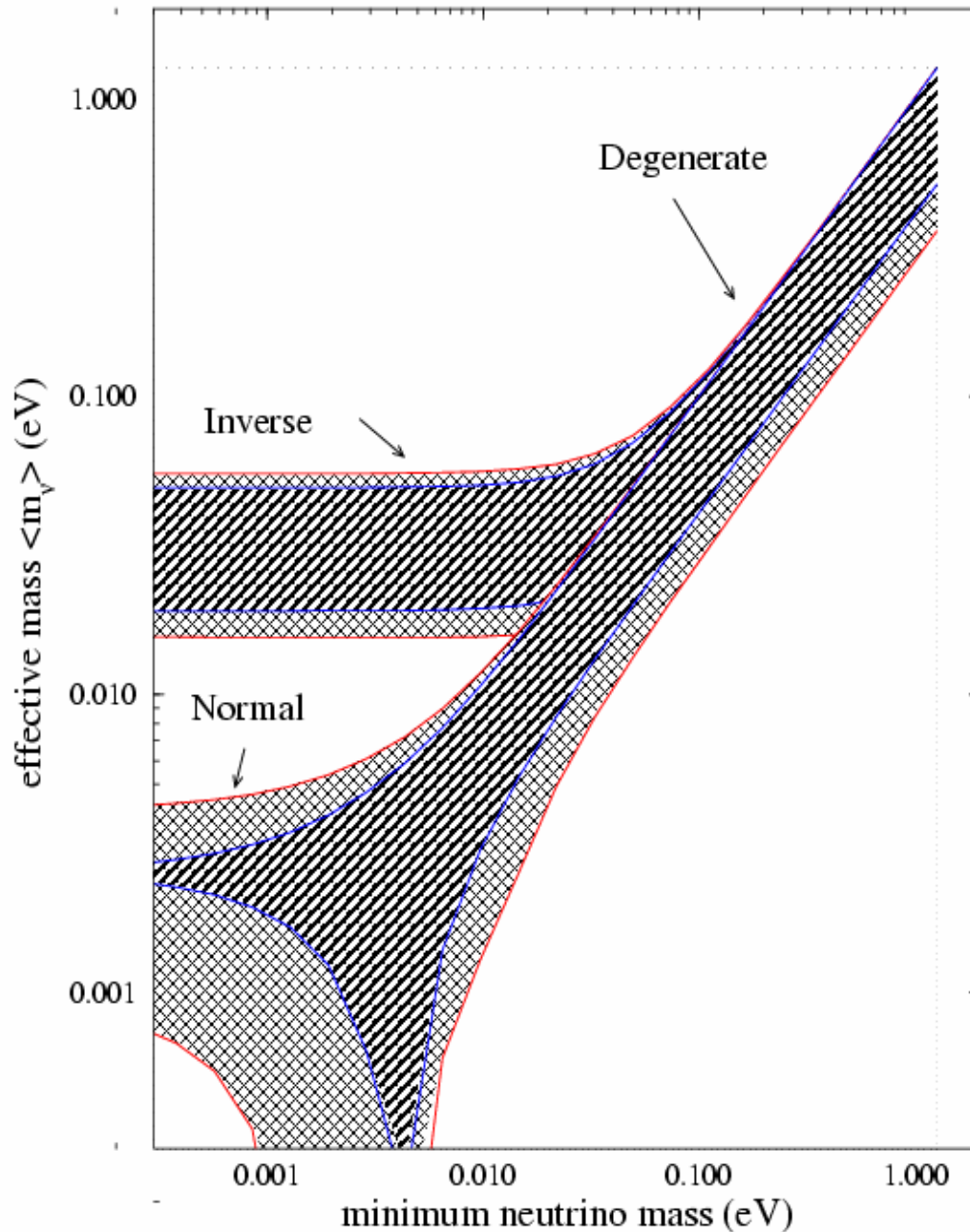
These upper and lower limits depend only on the oscillation parameters and on one mass  $m_{min}$ .

# What do we know?

Table 1

Neutrino Oscillation Parameters Determined From Various Experiments (2003)

Parameter	Value $\pm 1\sigma$	Reference	Comment
$\Delta m_{12}^2$	$7.1_{-0.6}^{+1.2} \times 10^{-5} \text{ eV}^2$	[71]	$8.2_{-0.5}^{+0.6} \times 10^{-5} (2004)$
$\theta_{12}$	$32.5_{-2.3}^{+2.4}$	[71]	For $\theta_{13} = 0$
$\Delta m_{32}^2$	$2.0_{-0.4}^{+0.6} \times 10^{-3} \text{ eV}^2$	[61]	
$\sin^2 2\theta_{23}$	$> 0.94$	[61]	For $\theta_{13} = 0$
$\sin^2 2\theta_{13}$	$< 0.11$	[63]	For $\Delta m_{atm}^2 = 2 \times 10^{-3} \text{ eV}^2$



**This is based on the parameters of the LMA solution. The cross-hatched region is for  $1\sigma$  errors. Arrows indicate that by determining  $\langle m_{\beta\beta} \rangle$ , even crudely we can fix the neutrino mass pattern.**

# Possible neutrino mass patterns

Normal

Inverse

Degenerate

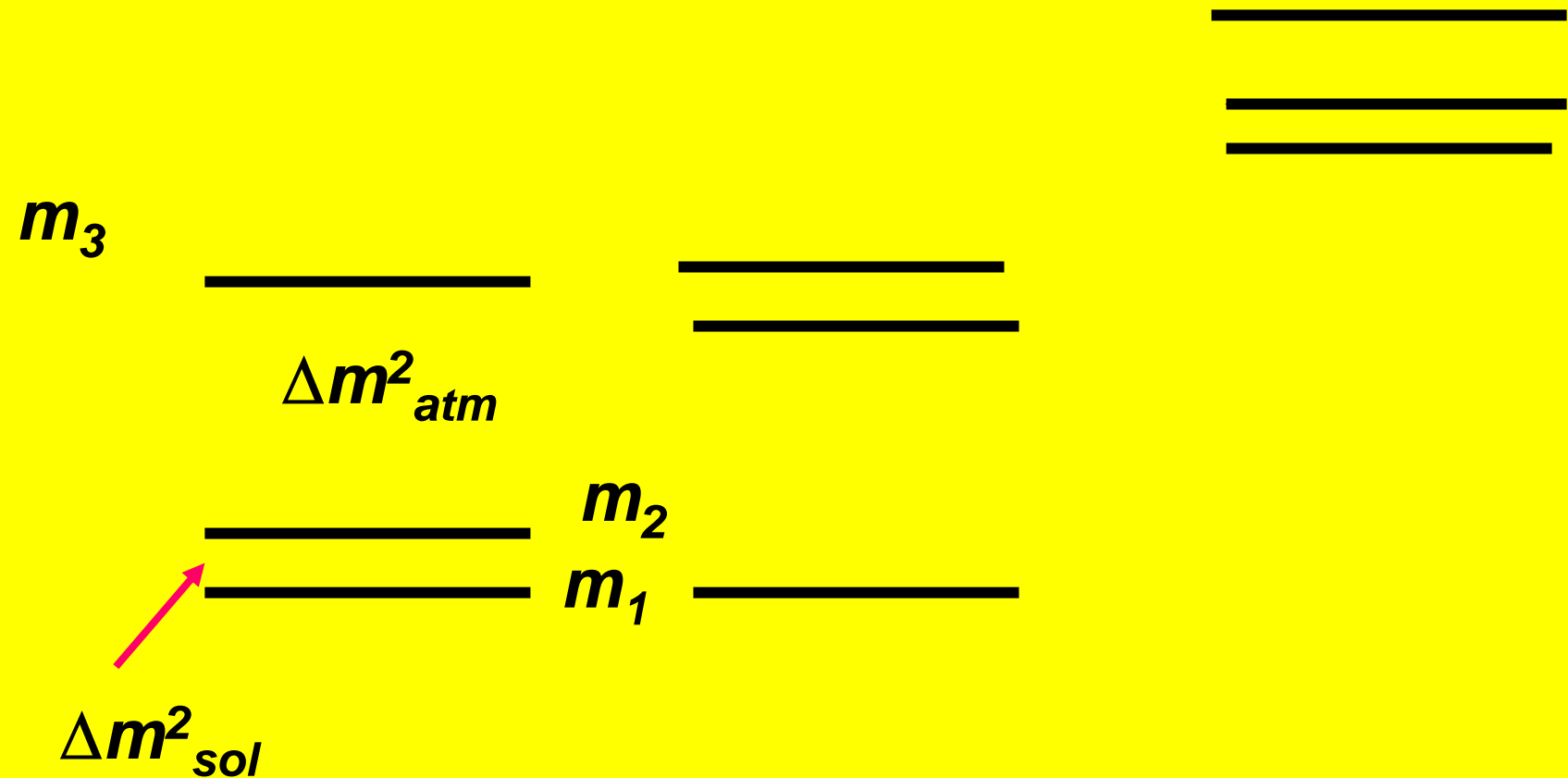
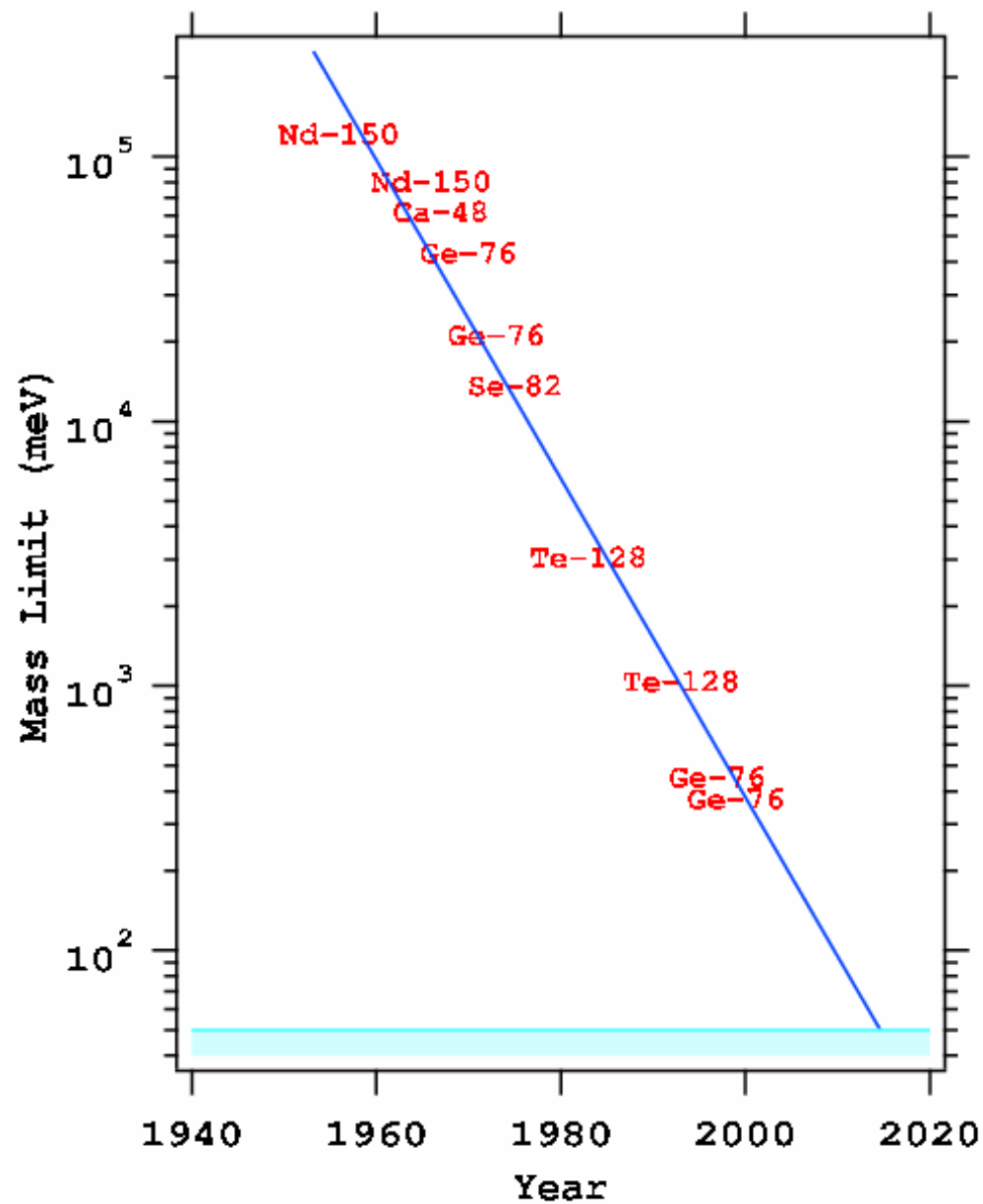


FIG. 4. Moore's plot for the  $0\nu\beta\beta$  decay. The shaded blue region at bottom indicates the desired sensitivity range.



The decay rate is related to  $\langle m_{\beta\beta} \rangle$  by

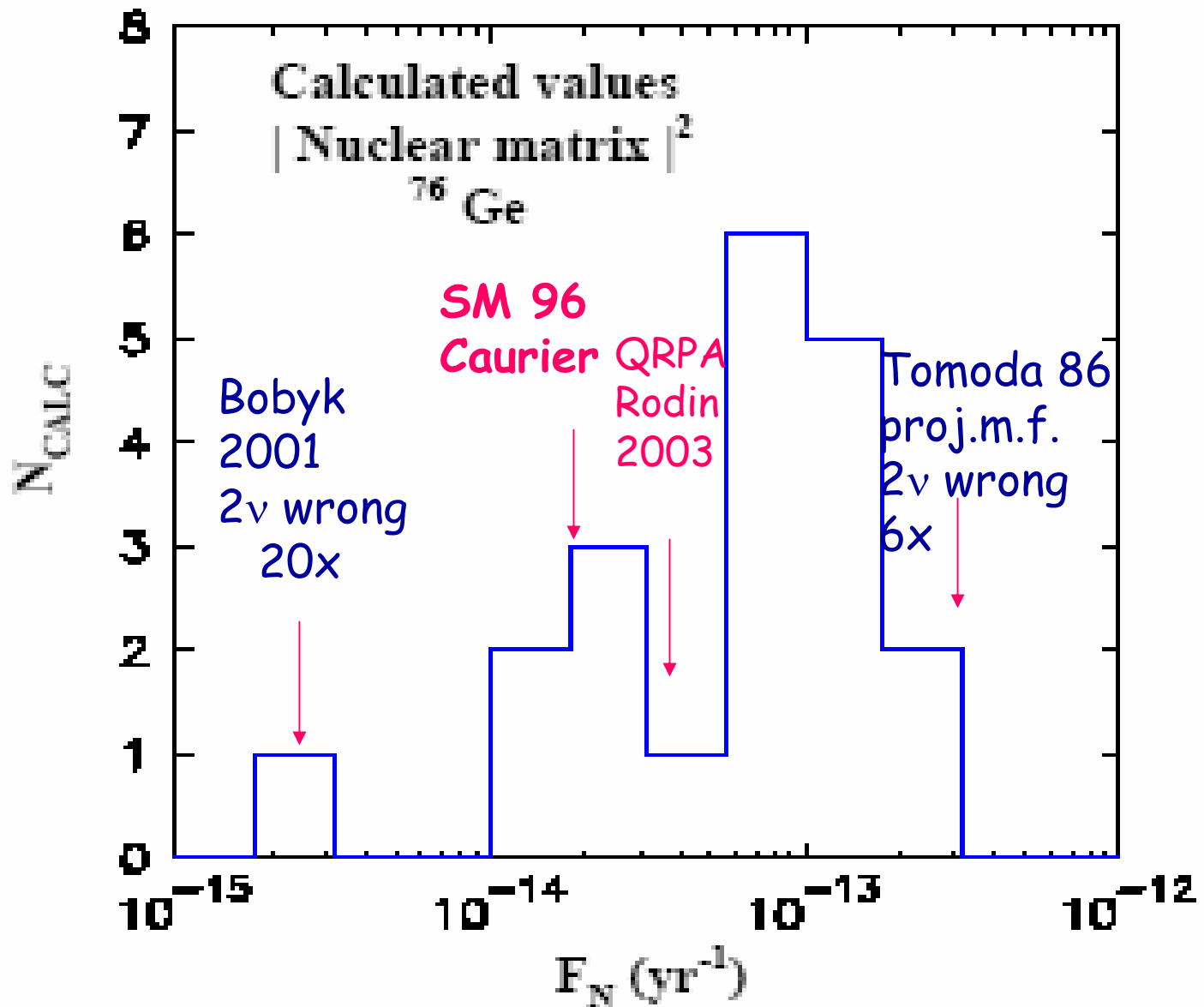
$$1/T_{1/2} = G^{0\nu}(E_0, Z) |M|^2 \langle m_{\beta\beta} \rangle^2,$$

Where  $G^{0\nu}(E_0, Z)$  is calculable phase space factor, and  $|M|^2$  is the square of the nuclear matrix element, calculable with difficulties.

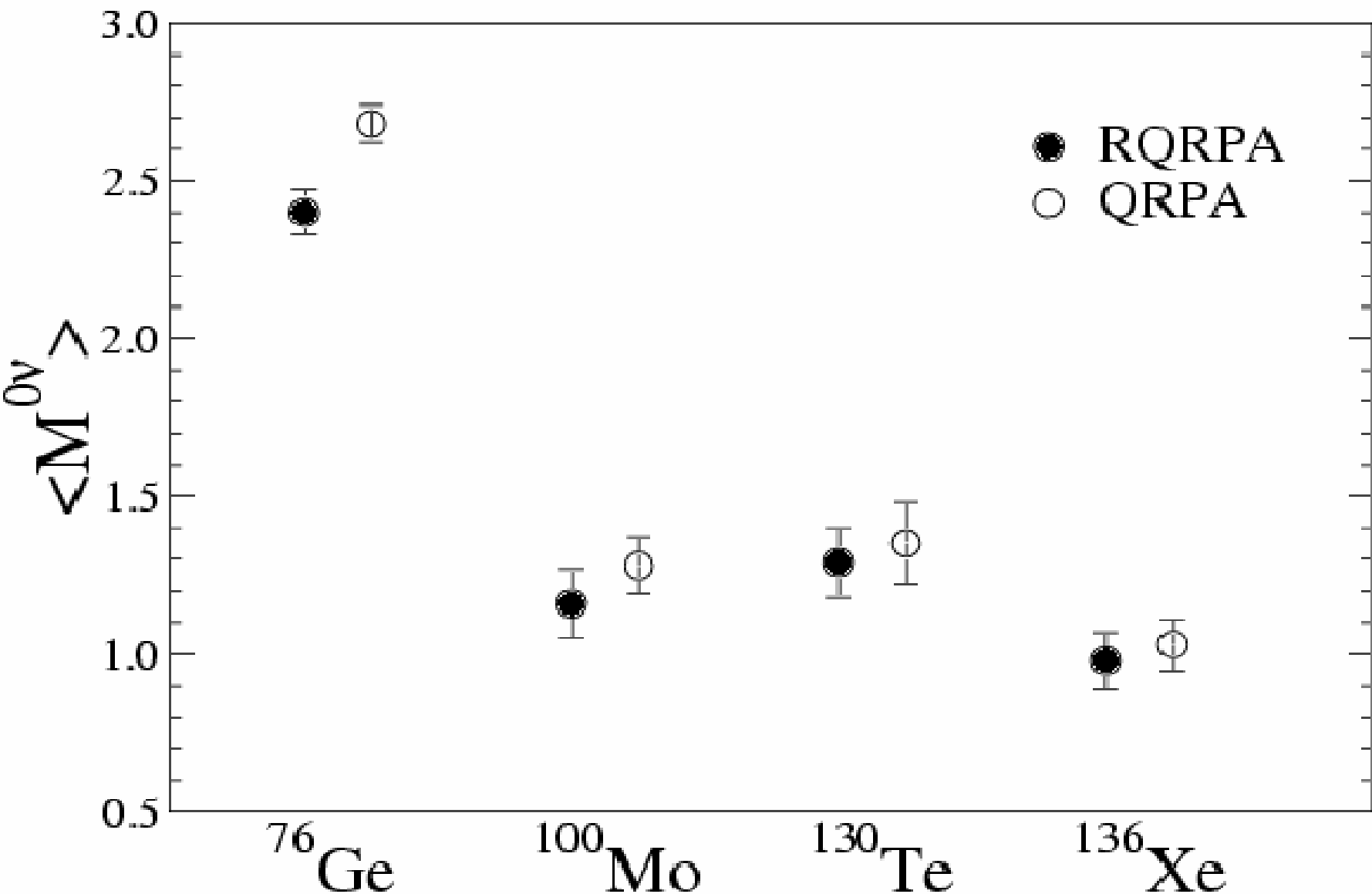
The two basic and complementary methods of evaluating  $M$  are the nuclear shell model (SM) and the Quasiparticle Random Phase Approximation (QRPA) and its various modifications.

One often uses  $C_{mm} = G^{0\nu}(E_0, Z) |M|^2 m_e^2$  to characterize the matrix elements

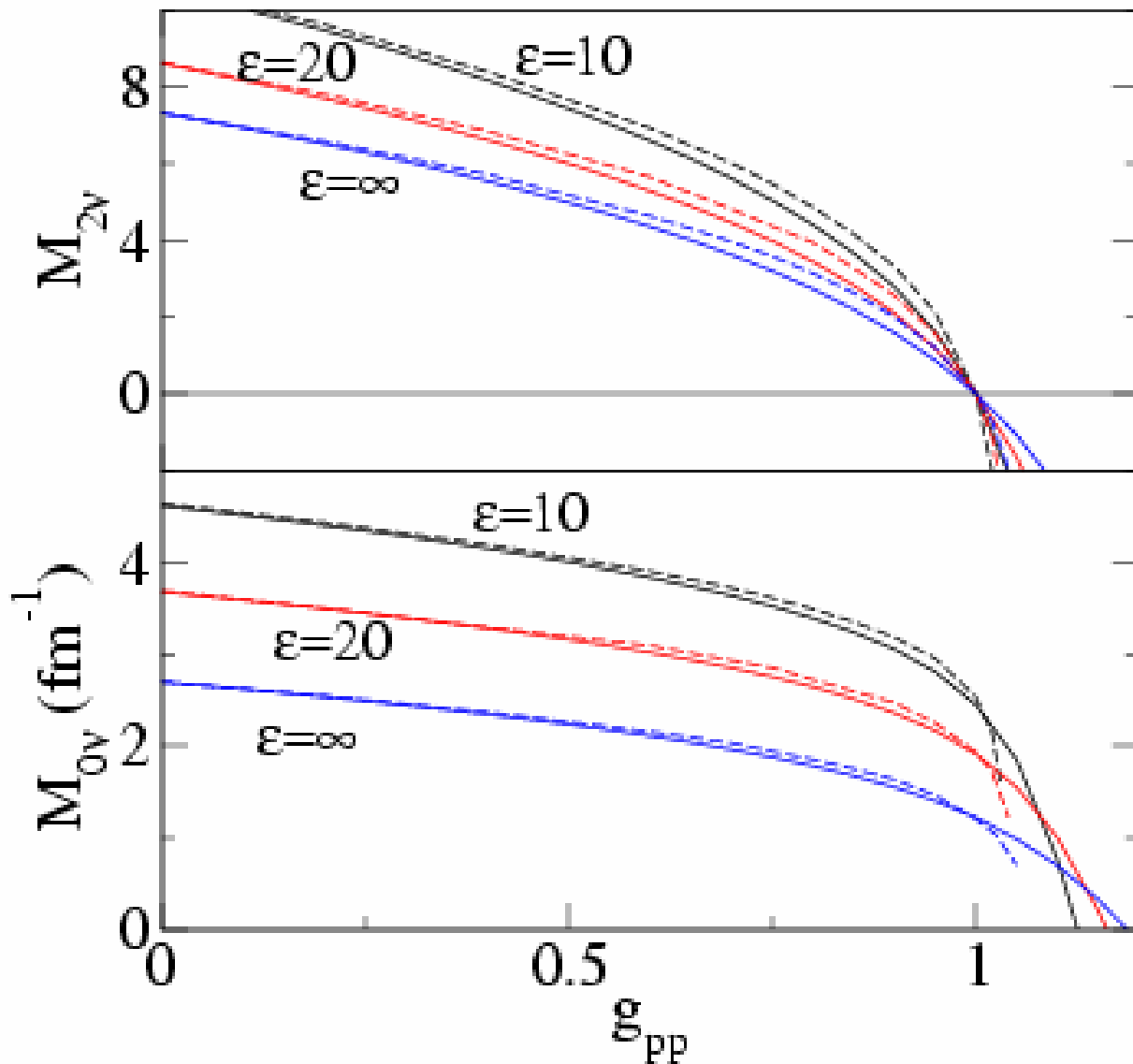
# Histogram from Bahcall et al. hep-ph/0403167



Average matrix elements and their variance  
see Rodin et al. Phys.Rev. C68,044302(2003)



From Engel & Vogel, PRC69,034304(2004)  
Comparison of exact (full) and QRPA (dashed) m.e.



Two shells,  $\epsilon$  apart,  
schematic force,  
but similar to the  
treatment of real  
nuclei.

# Comparison of the shell model and RQRPA matrix elements:

Shell model: Caurier et al., PRL 77(1996)1954 and  
Caurier et al., Nucl. Phys. A654(1999)973c

QRPA: Rodin et al., Phys. Rev. C68(2003)044302 and  
private comm. (for  $^{82}\text{Se}$ )

Nucleus	$C_{mm}(\text{s.m.})$	$C_{mm}(\text{RQRPA})$ (in $10^{-14} \text{ y}^{-1}$ )
$^{76}\text{Ge}$	1.9	3.6
$^{82}\text{Se}$	10.9	11.3
$^{130}\text{Te}$	4.5	7.4
$^{136}\text{Xe}$	2.2	3.4

Note:  $1/T_{1/2} = C_{mm} * (m_\nu/m_e)^2$

# Conclusions

- ❑ There is a good possibility that  $0\nu\beta\beta$  decay corresponding to  $\langle m_{\beta\beta} \rangle \sim 10 \text{ meV}$  really exists.
- ❑ It is of extreme interest to observe it.
- ❑ This can be achieved only with more challenging and costly experiments than those currently running.
- ❑ There is a wealth of proposals; the best of them clearly deserve support.
- ❑ Hand in hand with the development of novel detection schemes, the community should also concentrate on attempts to improve the evaluation of the nuclear matrix elements,

# APS Joint Study on the Future of Neutrino Physics (2004)

(preliminary version, subject to change):

Will recommend that a program of sensitive searches for neutrinoless double beta decay be initiated as soon as possible. Most likely a phased approach, with several scalable 100-200 kg experiments, will be recommended.

This is based on the understanding that the answer to the question whether neutrinos are their own antiparticles is a crucial ingredient in the development of a "New Standard Model"

# Search for Neutrinoless Double Beta Decay with Enriched $^{76}\text{Ge}$ in Gran Sasso 1990-2003

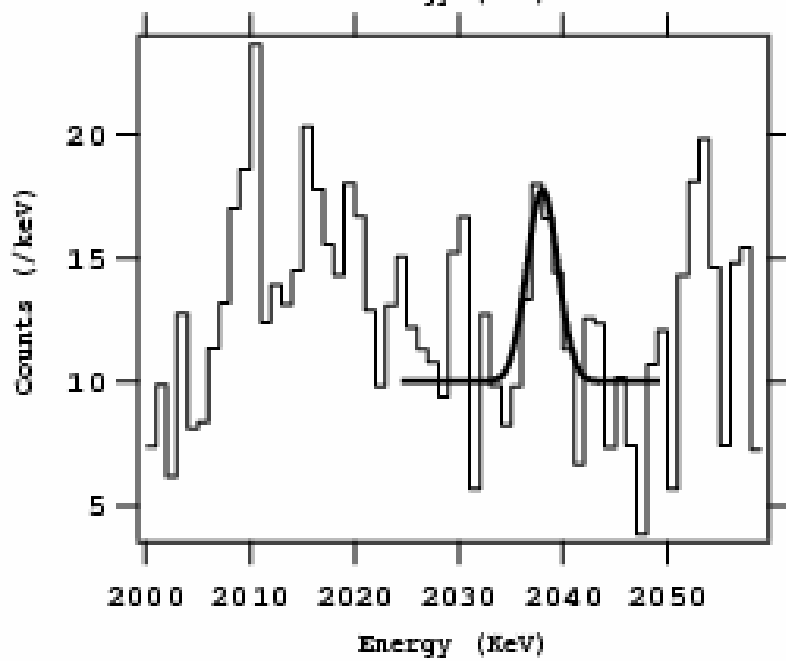
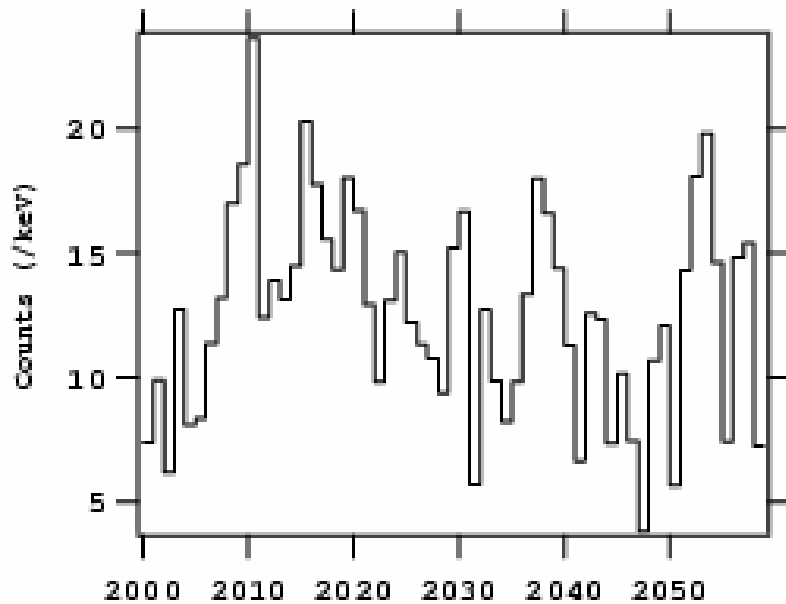
H.V. Klapdor-Kleingrothaus <sup>1</sup>,  
I.V. Krivosheina <sup>2</sup>, A. Dietz, and O. Chkvorets

*Max-Planck-Institut für Kernphysik, PO 10 39 80, D-69029 Heidelberg, Germany*

Phys. Lett. **B586**, 198 (2004); hep-ph/0404088

All detectors, 71.7 kg y,  $28.75 \pm 6.86$  events  
at 2039 keV, the Q-value for  $0\nu\beta\beta$  decay

This is interpreted as  $T_{1/2} = 1.15_{-0.5}^{+2.99} \times 10^{25}$  y  
or  $\langle m_{\beta\beta} \rangle = 0.44_{-0.2}^{+0.14}$  eV for the matrix  
elements chosen by the authors.



*From Klapdor-Kleingrothaus,  
Dietz, Krivosheina, Chkvorets,  
hep-hp/0403018.  
71.7 kg-y exposure,  
spectrum without and with  
the '0νββ' line.*

*Additional lines at 2010, 2017,  
2022, 2053 are assigned to the  
<sup>214</sup>Bi decay, line at 2030,  
previously not seen, is of an  
unknown origin.*

*Unidentified previously fitted  
lines at ~2066 and 2075 keV  
are not shown now.*

- To confirm or reject the KDHK claim of the  $0\nu\beta\beta$  discovery convincingly one needs 5-10 kmole-years exposure with low background.
- This is in accord with plans of most groups. So, no change in strategy is required. It might take  $\sim 3-5$  years to accomplish.
- One needs 100-200 kg of source material, and lower background than KDHK to be able to reject or confirm the claim at  $\sim 5\sigma$ .

Quote from the executive summary  
(preliminary version, subject to change)  
the APS Joint Study on the Future of  
Neutrino Physics (2004)

“We recommend, as a high priority, that a phased program of sensitive searches for neutrinoless double beta decay be initiated as soon as possible.”

It continues by noting that “An understanding of nature through a New Standard Model cannot be completed without the answer to this question”