

**Production of the X (3872)
in B meson Decay
by the Coalescence of Charm Mesons**

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Production of the X (3872) in B meson Decay
by the Coalescence of Charm Mesons

Outline

1. What is the X(3872)?
2. Low-Energy Universality
3. Decay amplitude for $B^+ \rightarrow XK^+$
4. Invariant mass distribution at threshold
5. Estimate for $\mathcal{B}[B^+ \rightarrow XK^+]$
6. Summary

What is the $X(3872)$?

Discovery by Belle

$$B^+ \rightarrow X(3872)K^+$$

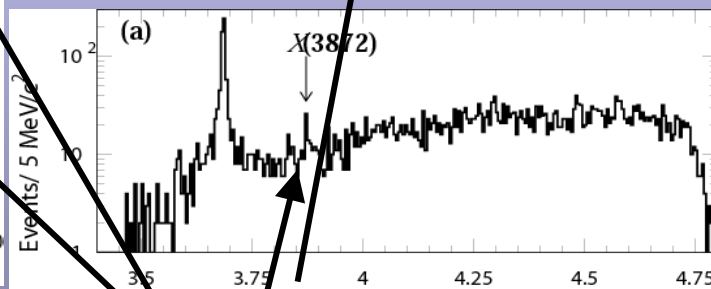
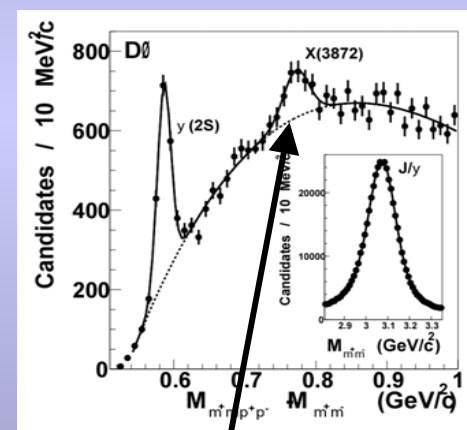
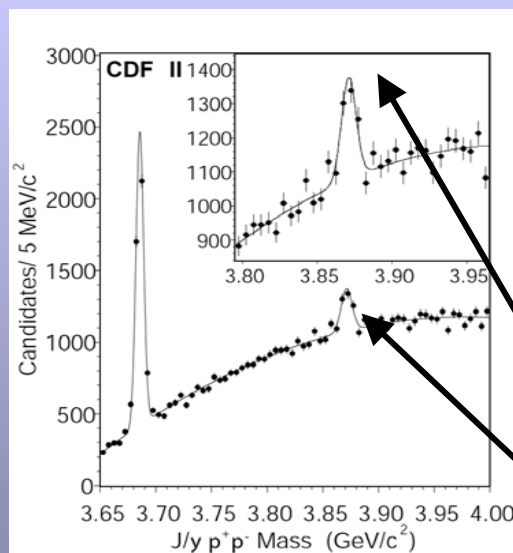
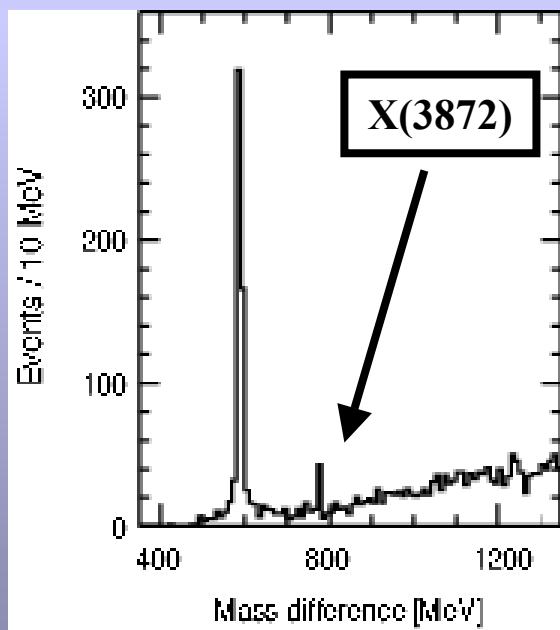
$$X \rightarrow J/\psi\pi^+\pi^-$$

$$3872.0 \pm 0.6 \pm 0.5 \text{ MeV}$$

$$\Gamma < 2.3 \text{ MeV}$$

Confirmed by CDF II, D0 and BABAR

$$M(J/\psi\pi^+\pi^-) = 3871.4 \pm 0.7 \pm 0.4 \text{ MeV}$$



$X(3872)$

Proposed Interpretations

1. Charmonium

D-wave: $J^{PC} = 2^{-\pm}$

(Eichten, Lane, Quigg)
(Barns, Godfrey)

P-wave: $J^{PC} = 1^{+\pm}$

Hybrid ($c\bar{c}g$): mass 4100 ± 200 MeV (ground-state)

(Juge, Kuti, Morningstar)

2. Loosely-bound $\bar{D}D^*$ Molecule

$$m_D + m_{D^*} = 3871.2 \pm 0.7 \text{ MeV}$$

$$m_X = 3871.9 \pm 0.6 \text{ MeV (Average)}$$

Bound by one pion exchange

$$J^{PC} = 1^{++} \text{ (S-wave)}$$

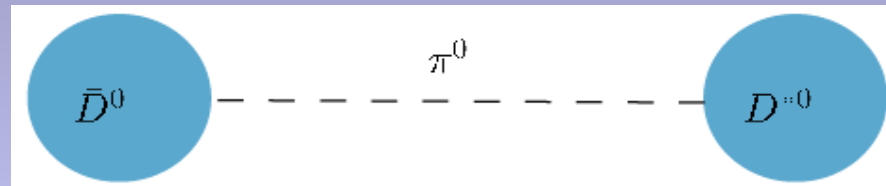
$$J^{PC} = 0^{-+} \text{ (P-wave)}$$

(Tornqvist)

If molecule, it is important to know the **production rate**

Low-Energy Universality

If X is S-wave $\bar{D}^0 D^{*0} / D^0 \bar{D}^{*0}$ Molecule:



Natural scale for the binding energy $\simeq m_\pi^2 / 2\mu = 10 \text{ MeV}$

1. Extremely small binding energy ($< 0.4 \text{ MeV}$)

$$E_b = m_D + m_{D^*} - m_X \simeq \frac{1}{2\mu a^2}$$

$> 7 \text{ fm}$

2. Some properties depend only on a (s-wave scattering length),

insensitive to shorter scales of QCD \longrightarrow low-energy universality

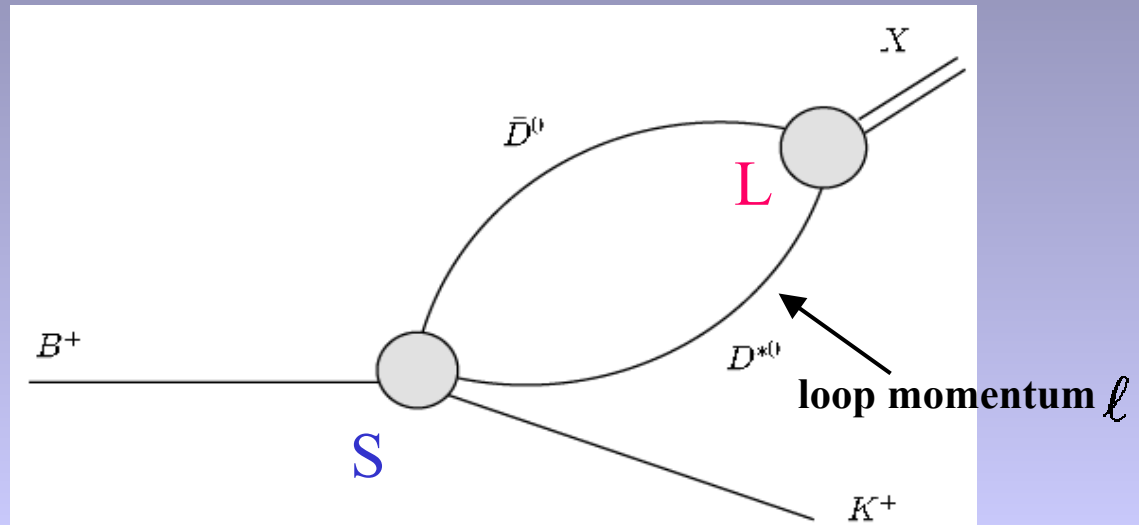
3. Universal wave function

$$|\mathbf{k}| \ll m_\pi \quad \Psi(\mathbf{k}) = \left(\frac{8\pi}{a} \right)^{1/2} \frac{1}{\mathbf{k}^2 + 1/a^2}$$

Production of X from the B^+ Decay

Discovery mode:

$$B^+ \longrightarrow XK^+$$



Short-distance Process: $B^+ \longrightarrow \bar{D}^0 D^{*0} K^+$

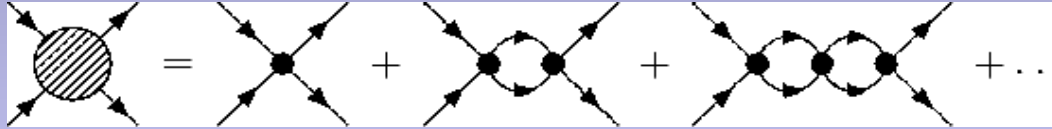
Long-distance Process: $\bar{D}^0 D^{*0} \longrightarrow X$

Amplitude $\mathcal{A}_1[B^+ \rightarrow XK^+]$ is

$$-i \int \frac{d^4 \ell}{(2\pi)^4} \mathcal{A}[B^+ \rightarrow D^0 \bar{D}^{*0} K^+] \mathcal{A}[D^0 \bar{D}^{*0} \rightarrow X] \\ \times \frac{1}{(q + \ell)^2 - m_D^2 + i\epsilon} \frac{1}{(q' - \ell)^2 - m_{D^*}^2 + i\epsilon}$$

Coalescence Process $\mathcal{A}[\bar{D}^0 D^{*0} \rightarrow X]$

Large S-wave Scattering Length a \longrightarrow Resonant Scattering
 $k_{\text{cm}} \ll m_\pi$



$$\bar{D}^0 D^{*0} \rightarrow X \rightarrow \bar{D}^0 D^{*0}$$

Universal elastic scattering amplitude:

$$\mathcal{A}[\bar{D}^0 D^{*0} \rightarrow \bar{D}^0 D^{*0}] \simeq \frac{8\pi m_D m_{D^*}}{\mu (-1/a - ik_{\text{cm}})}$$

Universal coalescence amplitude:

$$\mathcal{A}[\bar{D}^0 D^{*0} \rightarrow X] = \left(\frac{16\pi Z m_X m_D m_{D^*}}{\mu^2 a} \right)^{1/2} \epsilon_X^* \cdot \epsilon_{D^*}$$

What about the loop integral over ℓ ?

$$\int \frac{d^4\ell}{(2\pi)^4} \frac{1}{(q+\ell)^2 - m_D^2 + i\epsilon} \frac{1}{(q'-\ell)^2 - m_{D^*}^2 + i\epsilon}$$

$$\simeq \int^\Lambda \frac{d^3\ell}{(2\pi)^3} \frac{i\mu a^{1/2}}{4\sqrt{2}\pi^{1/2}m_D m_{D^*}} \Psi(\ell)$$

$$= \frac{i\mu\Lambda}{4\pi^2 m_D m_{D^*}}$$

Momentum-space wave function

→ Linearly depend on the ultraviolet cutoff Λ

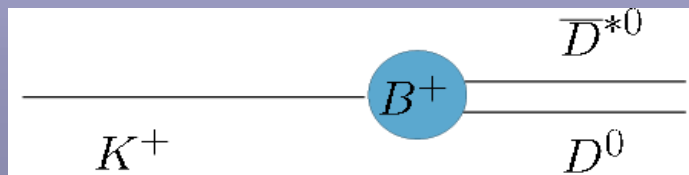
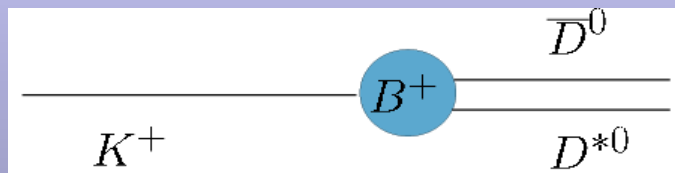
($\Lambda \simeq m_\pi$) : chosen for an order-of-magnitude estimate

What about the **short**-distance process?

$$\mathcal{A}[B^+ \longrightarrow \bar{D}^0 D^{*0} K^+] \quad (\text{or } \mathcal{A}[B^+ \longrightarrow D^0 \bar{D}^{*0} K^+])$$

In order to produce X(3872) by coalescence, D mesons should have small relative momentum and are almost collinear

At the corner of phase space where \bar{D}^0 and D^{*0} are **collinear**:



$$= c_1 P \cdot \epsilon^*$$

constant

momentum of B^+

$$= c_2 P \cdot \epsilon^*$$

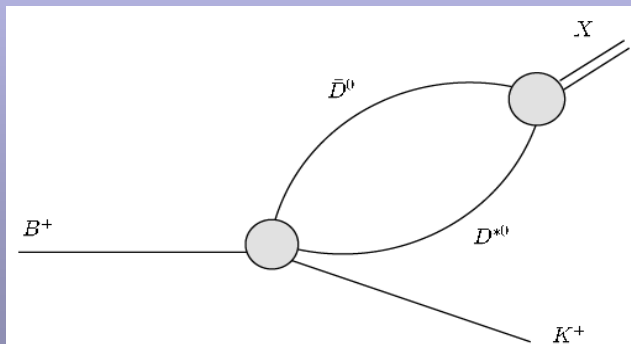
Putting all together...

$$\Gamma[B^+ \rightarrow XK^+] = \frac{\lambda^{3/2}(m_B, m_X, m_K)}{64\pi^4 m_B^3 m_X^2 \mu a} |c_1 \pm c_2|^2 \Lambda^2$$

+ : C = even

- : C = odd

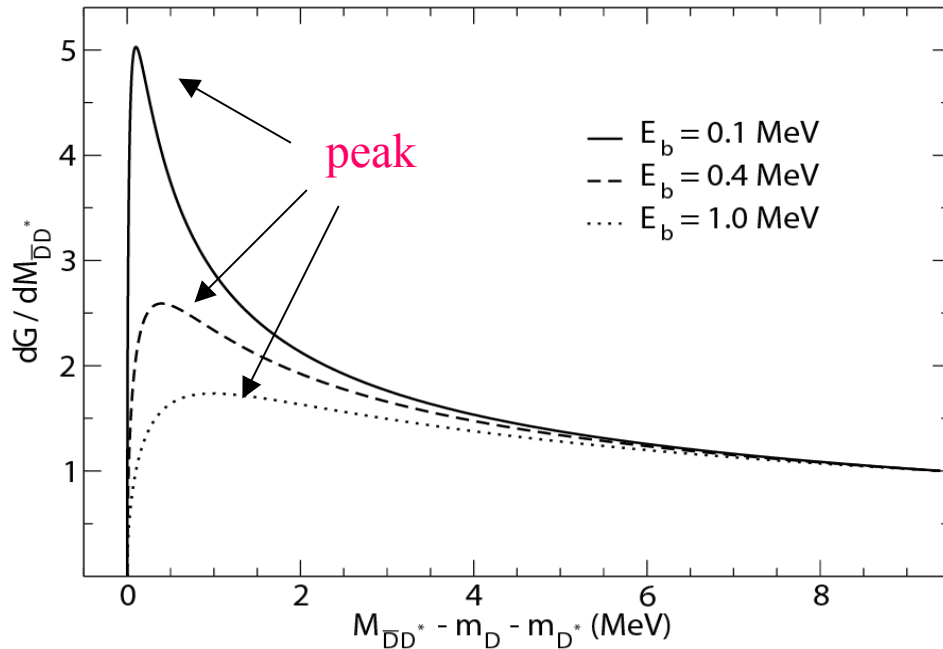
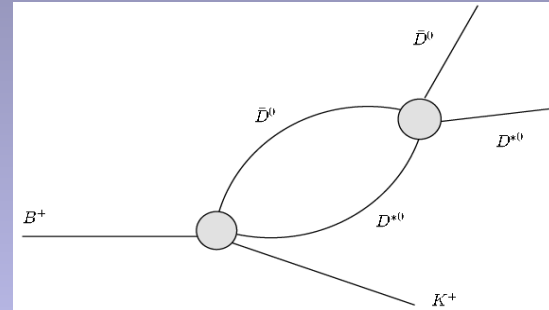
$$\frac{1}{\sqrt{2}} (|D^0 \bar{D}^{*0}\rangle \pm |\bar{D}^0 D^{*0}\rangle)$$



$$\pm \begin{pmatrix} D^0 & \bar{D}^0 \\ \bar{D}^{*0} & D^{*0} \end{pmatrix}$$

$\bar{D}^0 D^{*0}$ invariant mass distribution at the threshold

$B^+ \rightarrow \bar{D}^0 D^{*0} K^+$ near the threshold



$$\frac{d\Gamma}{dM_{\bar{D}D^*}} [B^+ \rightarrow \bar{D}^0 D^{*0} K^+] = \Gamma[B^+ \rightarrow X K^+] \frac{\mu a^3 k_{\text{cm}}}{\pi(1 + a^2 k_{\text{cm}}^2)}$$

Observation of a **peak** would confirm the X(3872) as a loosely-bound molecule

An Estimate for c_1 and c_2 from Babar data

$$\mathcal{A}[B^+ \longrightarrow \bar{D}^0 D^{*0} K^+] = c_1 P \cdot \epsilon^*$$

Babar has measured the branching fractions of $B \rightarrow \bar{D}^{(*)} D^{(*)} K$

1. Use factorization assumption and Isospin relations (Bauer et.al.)

For example,

$$\begin{aligned} \mathcal{A}[B^+ \rightarrow \bar{D}^0 D^{*0} K^+] &\rightarrow \lambda_1 \langle \bar{D}^0 | \bar{c} \gamma^\mu (1 - \gamma_5) b | B^+ \rangle \langle D^{*0} K^+ | \bar{s} \gamma^\mu (1 - \gamma_5) c | 0 \rangle \\ &\quad + \lambda_2 \langle K^+ | \bar{s} \gamma^\mu (1 - \gamma_5) b | B^+ \rangle \langle \bar{D}^0 D^{*0} | \bar{c} \gamma^\mu (1 - \gamma_5) c | 0 \rangle \end{aligned}$$

$$\mathcal{A}[B^+ \rightarrow \bar{D}^0 D^+ K^0] \rightarrow \lambda_1 \langle \bar{D}^0 | \bar{c} \gamma^\mu (1 - \gamma_5) b | B^+ \rangle \langle D^+ K^+ | \bar{s} \gamma^\mu (1 - \gamma_5) c | 0 \rangle$$

2. Need 4 form factors from Heavy-quark spin symmetry

For example,

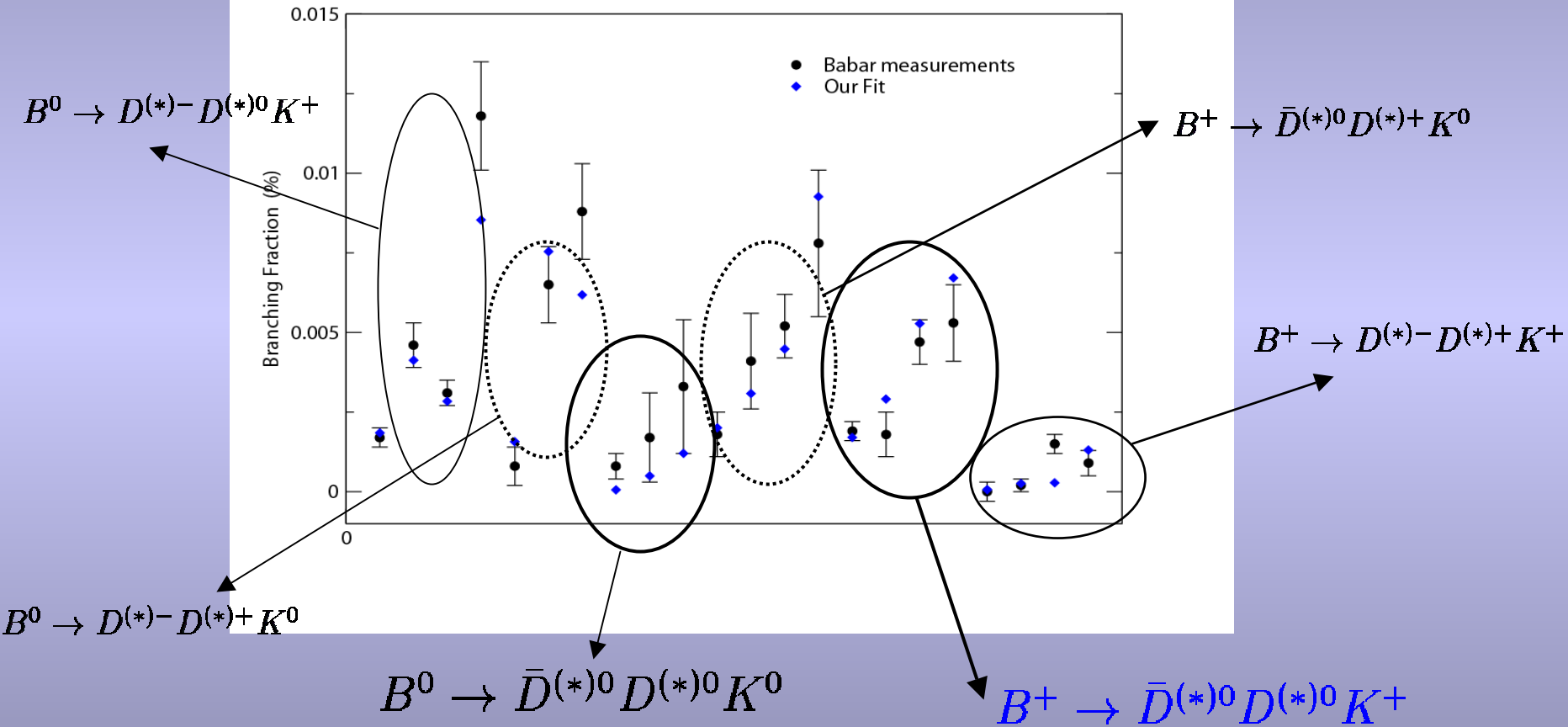
$$\mathcal{A}[B^+ \rightarrow \bar{D}^0 D^{*0} K^+] = -i G_1 \epsilon^* \cdot (V + v)$$

$$-i (G_2 / m_B) \epsilon_\nu^* [v_* \cdot k (V + v)^\nu - v_* \cdot (V + v) k^\nu$$

$$-i \epsilon^{\nu\mu\alpha\beta} (V + v)_\mu v_{*\alpha} k_\beta]$$

3. Fit $G_{i=1,2,3,4}$ for the best description of the data, assuming they are constants (7 parameters to fit)

Branching fractions of 3-body double charm decay:



One can estimate c_1 and c_2 from this analysis:

$$c_1 = c_2 = -i[m_B G_1 - (m_B + m_D + m_{D^*}) G_2] / m_B^2$$

If $\mathcal{C} = +$ (constructive interference),

$$\mathcal{B}[B^+ \rightarrow XK^+] \approx 3.9 \times 10^{-5} \frac{\Lambda^2}{m_\pi^2} \left(\frac{E_b}{0.4 \text{ MeV}} \right)^{1/2}$$

If we take $\Lambda \simeq m_\pi$ and $E_b = 0.4 \text{ MeV}$ and if $\mathcal{C} = +$, ($J^{PC} = 1^{++}$)

$$\mathcal{B}[B^+ \rightarrow XK^+] \simeq 3.9 \times 10^{-5} \quad (\text{our estimate})$$

(Belle collaboration)

$$\mathcal{B}[B^+ \rightarrow XK^+] \times \mathcal{B}[X \rightarrow J/\psi \pi^+ \pi^-] = (1.3 \pm 0.3) \times 10^{-5}$$

Our estimate is compatible, if $J^{PC} = 1^{++}$

and $X \rightarrow J/\psi \pi^+ \pi^-$ is one of the major decay modes of X(3872)

Summary

Applied the **low-energy universality** analysis, assuming the X(3872) is S-wave $\bar{D}^0 D^{*0} / D^0 \bar{D}^{*0}$ molecule.

Invariant mass distribution $d\Gamma/dM_{DD^*} [X \rightarrow \bar{D}^0 D^{*0}]$
should have a peak at the threshold

Applied the **heavy-quark spin symmetry** to estimate unknown constants c_1 and c_2 from Babar measurements

An order of magnitude estimate is compatible

if X has $J^{PC} = 1^{++}$ and

$X \rightarrow J/\psi \pi^+ \pi^-$ is one of the major decay modes